

Quarkonium detection and physics with ECCE

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27/10/21



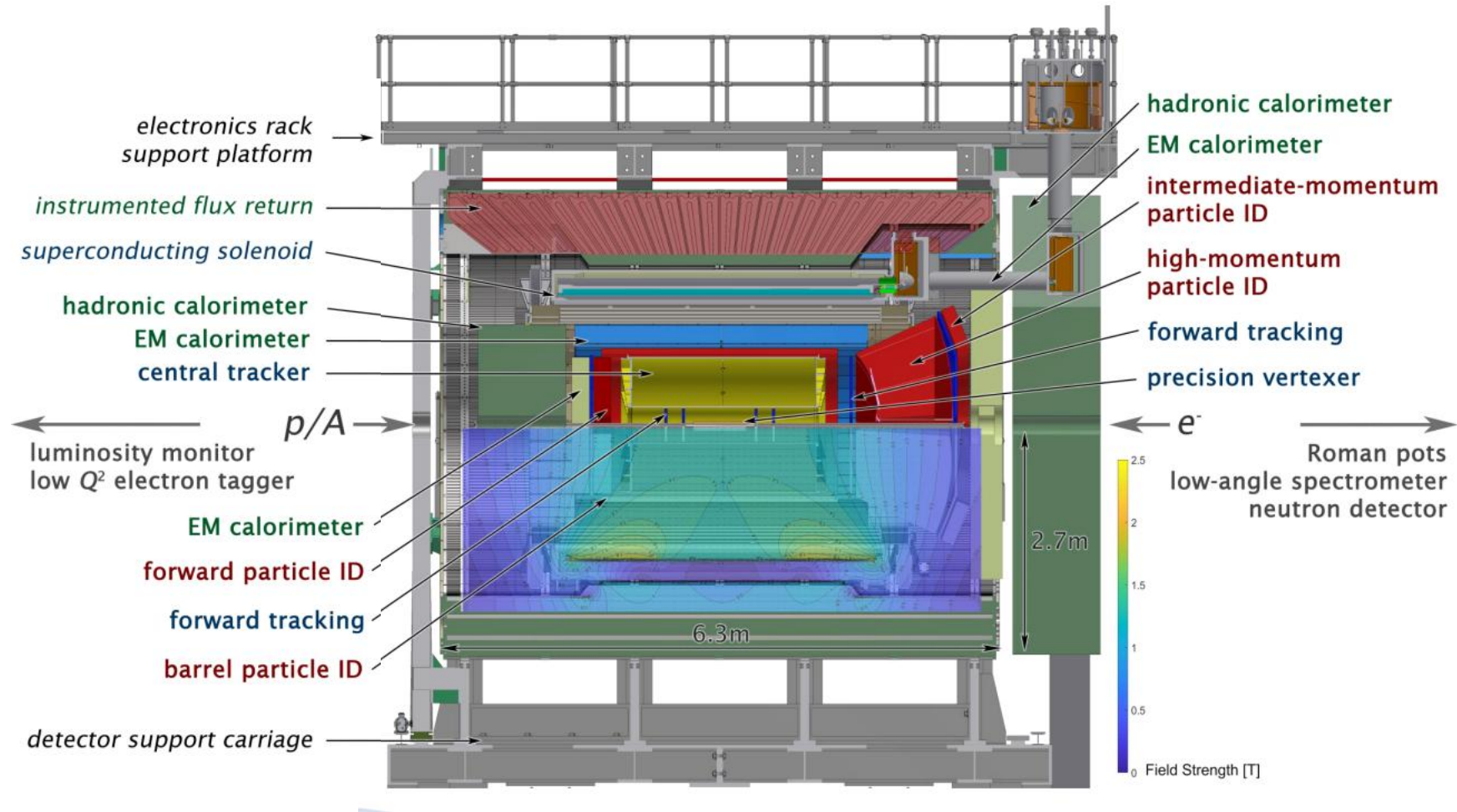
Outline

- Physics Opportunities for Heavy Quarkonia at the EIC
- Detector Configuration of ECCE
- Tracking and electron identification capability
- The simulation of Quarkonia reconstruction at ECCE
- The projection results for J/ψ photoproduction
- Summary and future plan

Physics Opportunities

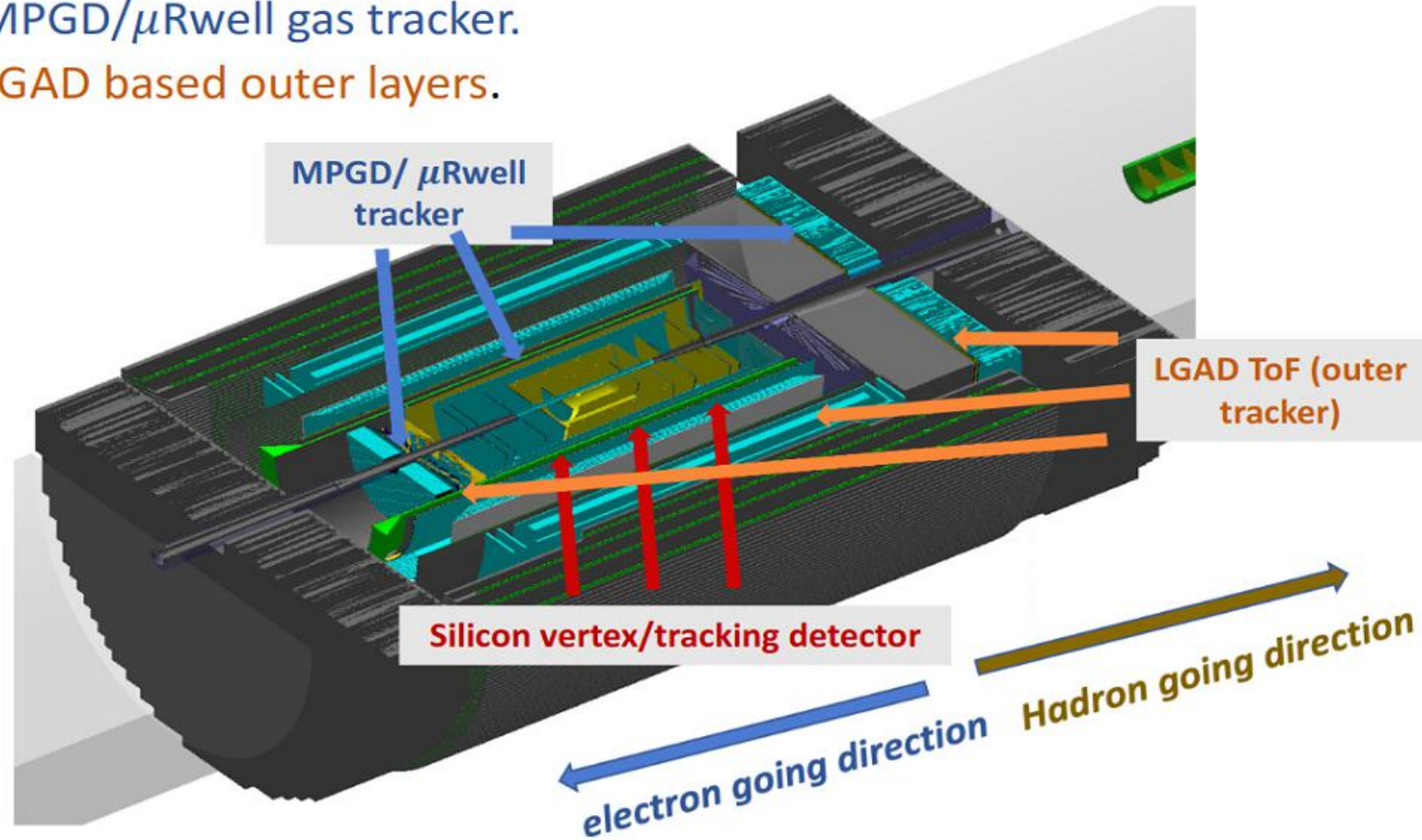
- Production mechanism for quarkonia
 - ✓ Constrain the NRQCD matrix elements
 - ✓ Study the hadronization in nucleus
- 3D tomography of gluon distribution
 - ✓ gluon nPDF (z direction)
 - ✓ Transverse distribution of gluon (x-y direction)
- Near threshold photoproduction
 - ✓ 2g, 3g exchange
 - ✓ The proton mass decomposition

Detector Configuration

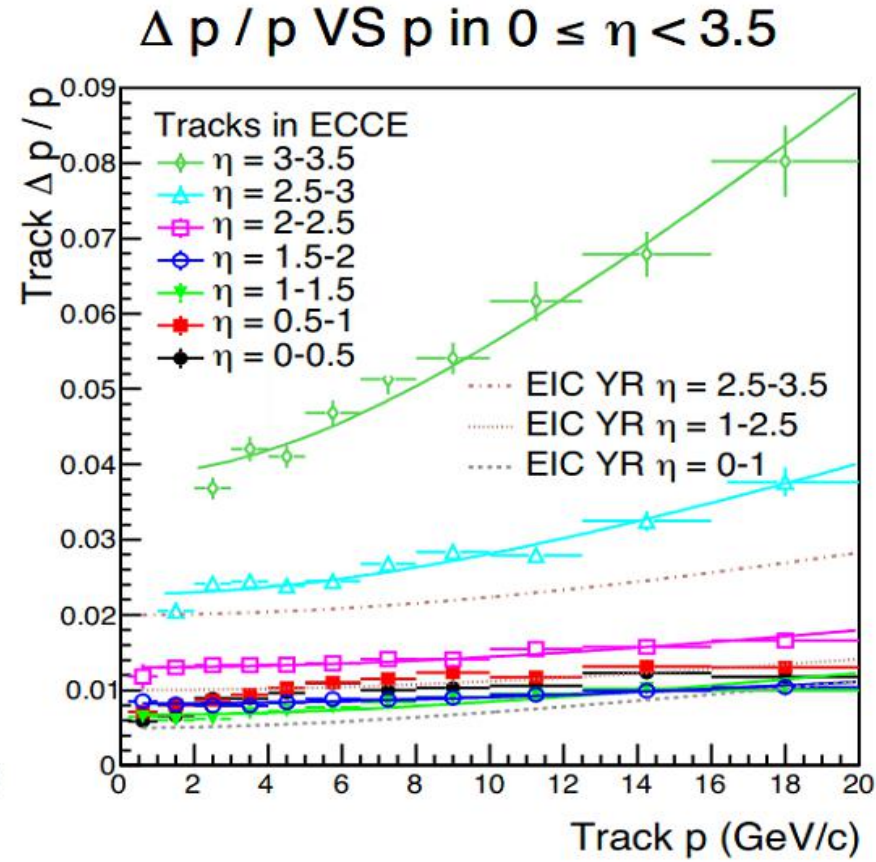
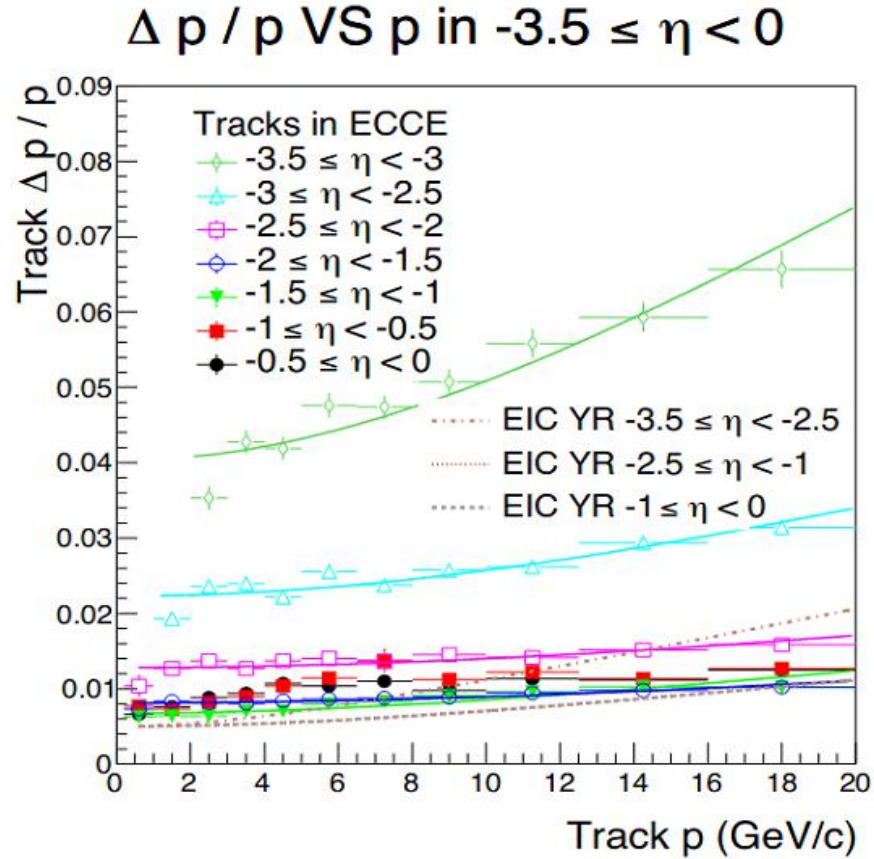


Detector Configuration (June Concept)——Tracking

- 1st simulation camping ECCE tracking detector consists of
 - MAPS based silicon vertex/tracking layers/planes.
 - MPGD/ μ Rwell gas tracker.
 - LGAD based outer layers.



Tracking performance at ECCE

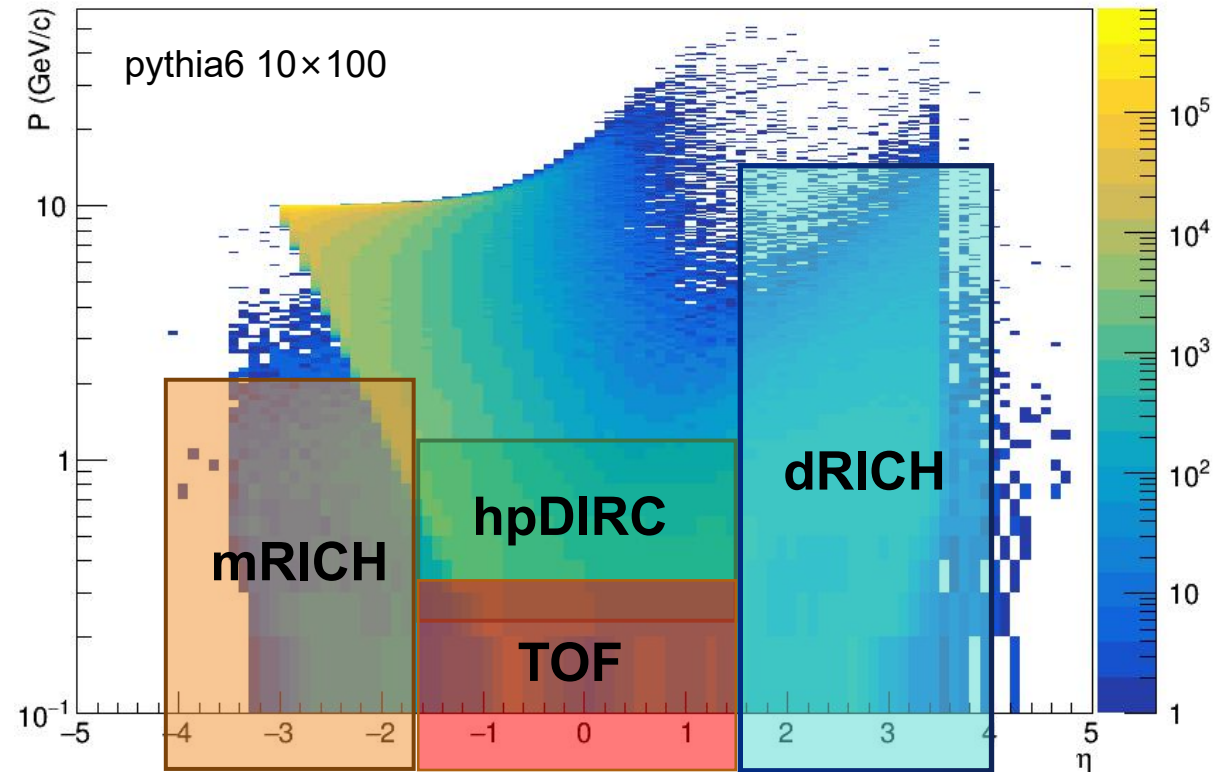
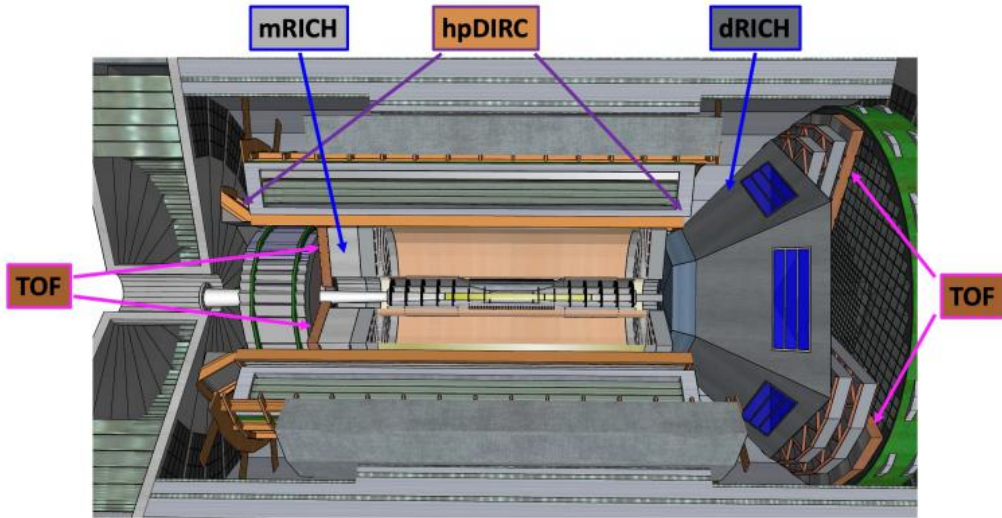


Electron identification capability at ECCE

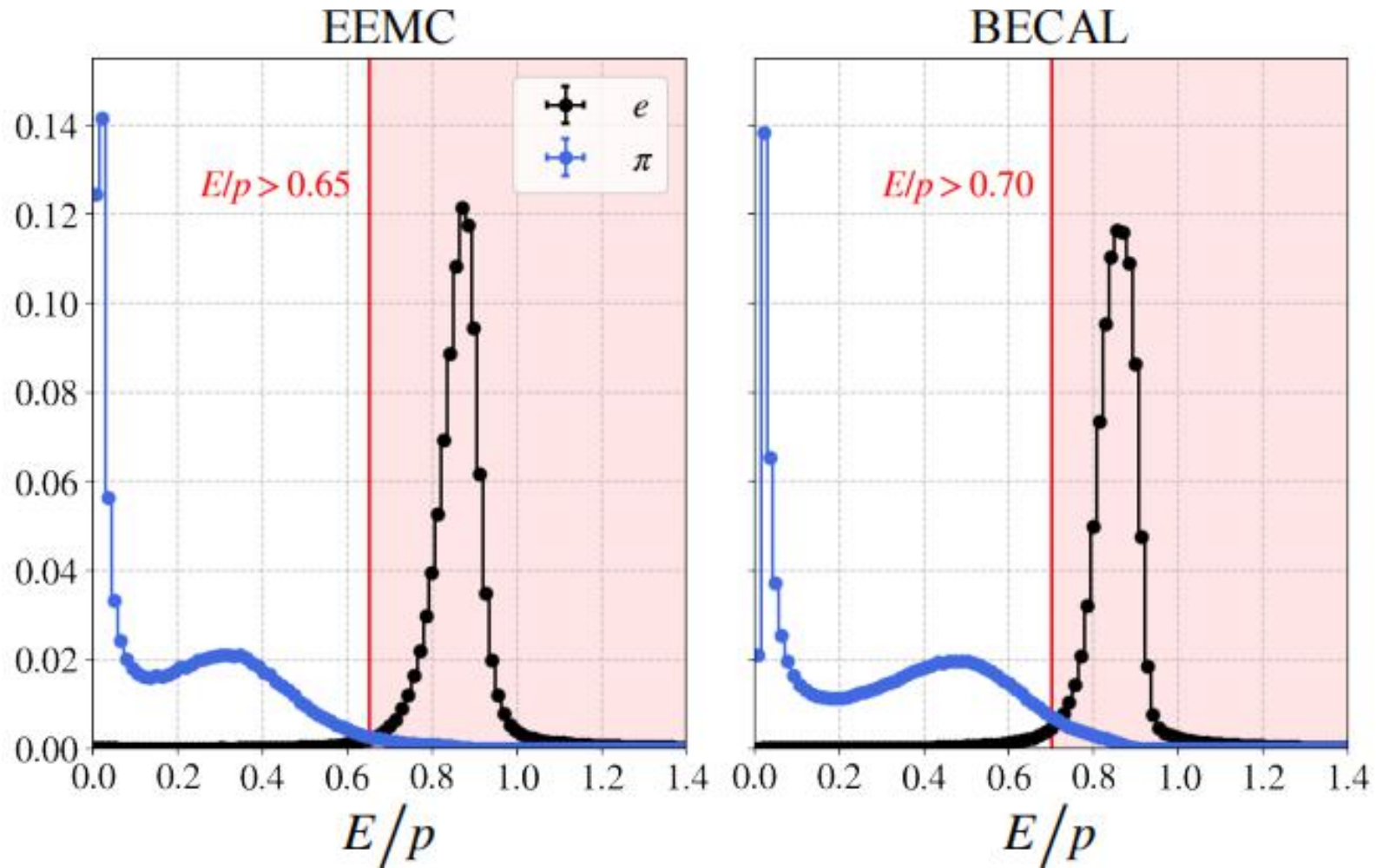
EMCal + Tracking

- ✓ The energy deposition => E/p cut
- ✓ The transverse profile of the showers
- ✓ The position resolution

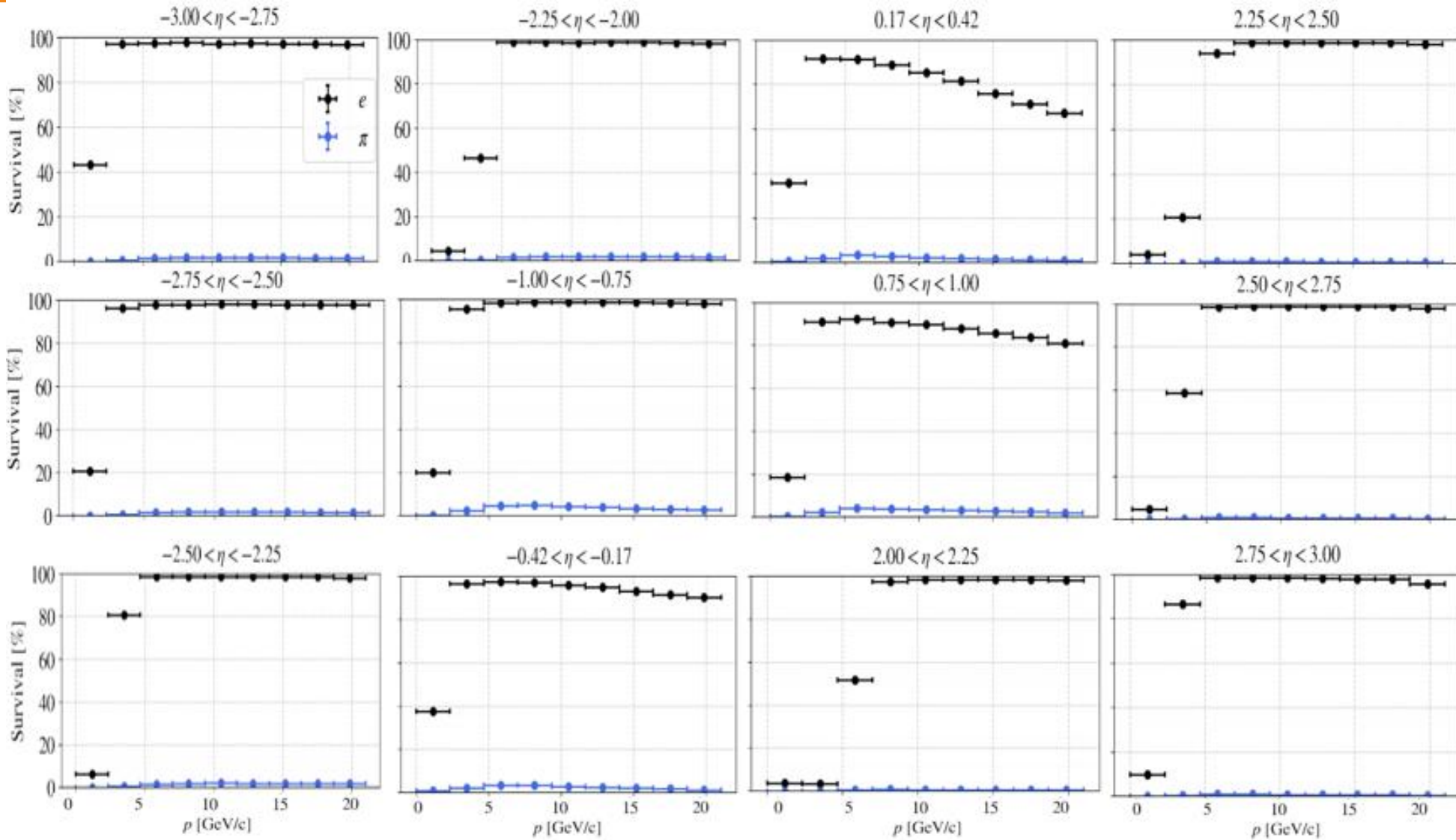
Cherenkov + TOF



Electron identification capability at ECCE



Electron identification capability at ECCE



Electron identification capability at ECCE

- **h-endcap: dRICH with two radiators (gas + aerogel)**

π/K separation up to ~ 50 GeV/c

e/π separation up to ~ 15 GeV/c

- **e-endcap: compact aerogel mRICH**

π/K separation up to ~ 10 GeV/c

e/π separation up to ~ 2 GeV/c

- **barrel: compact high-performance DIRC**

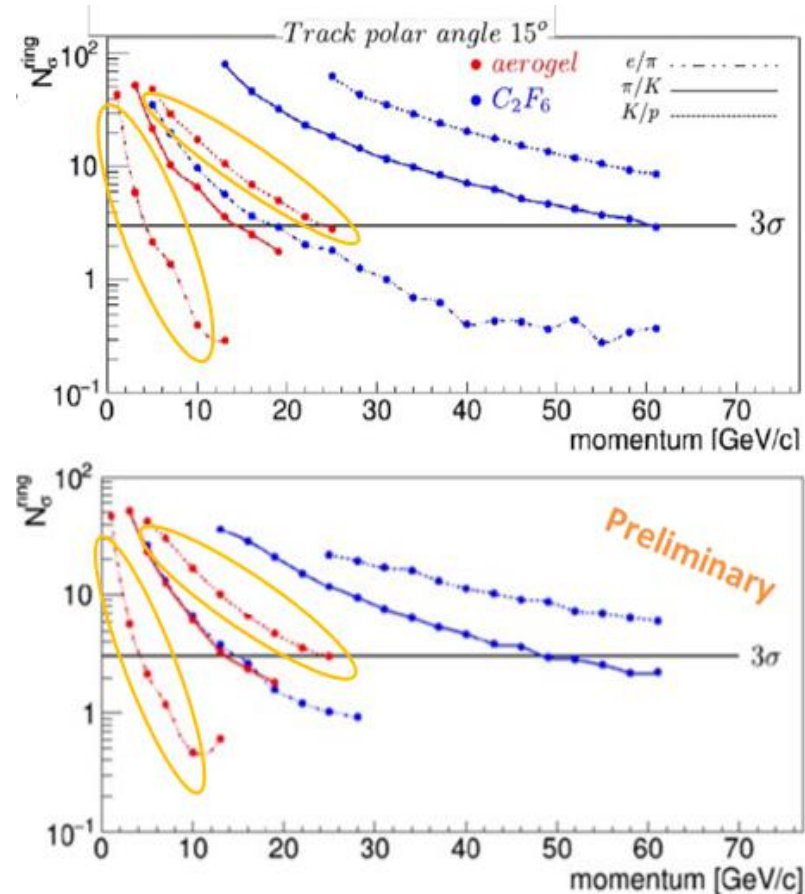
π/K separation up to $\sim 6-7$ GeV/c

e/π separation up to ~ 1.2 GeV/c

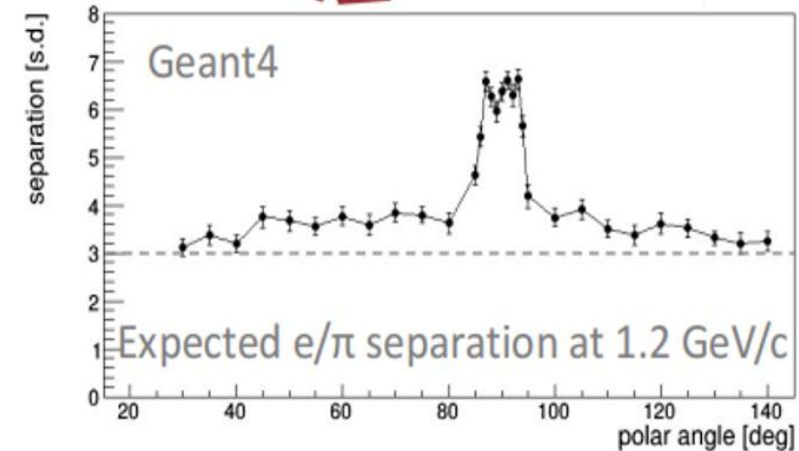
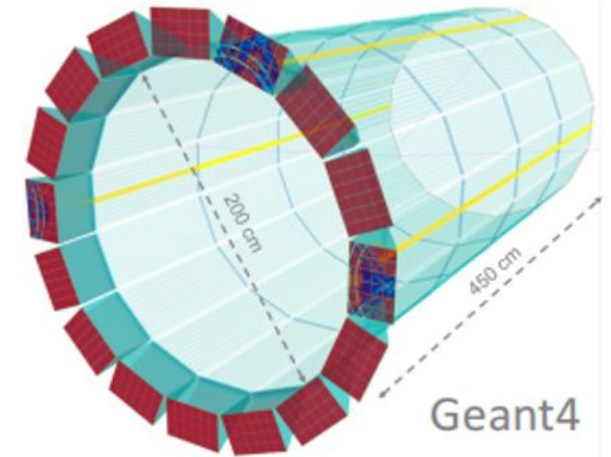
- **LGAD based TOF:**

cover lower momenta down to ~ 0.2 GeV/c

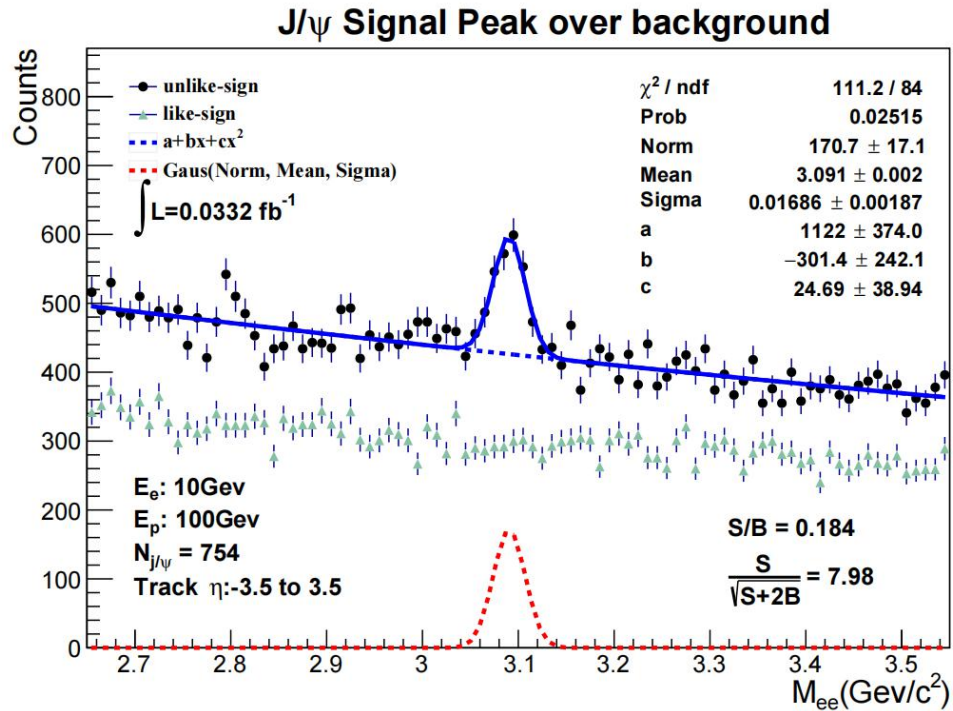
dRICH Simulated Performance



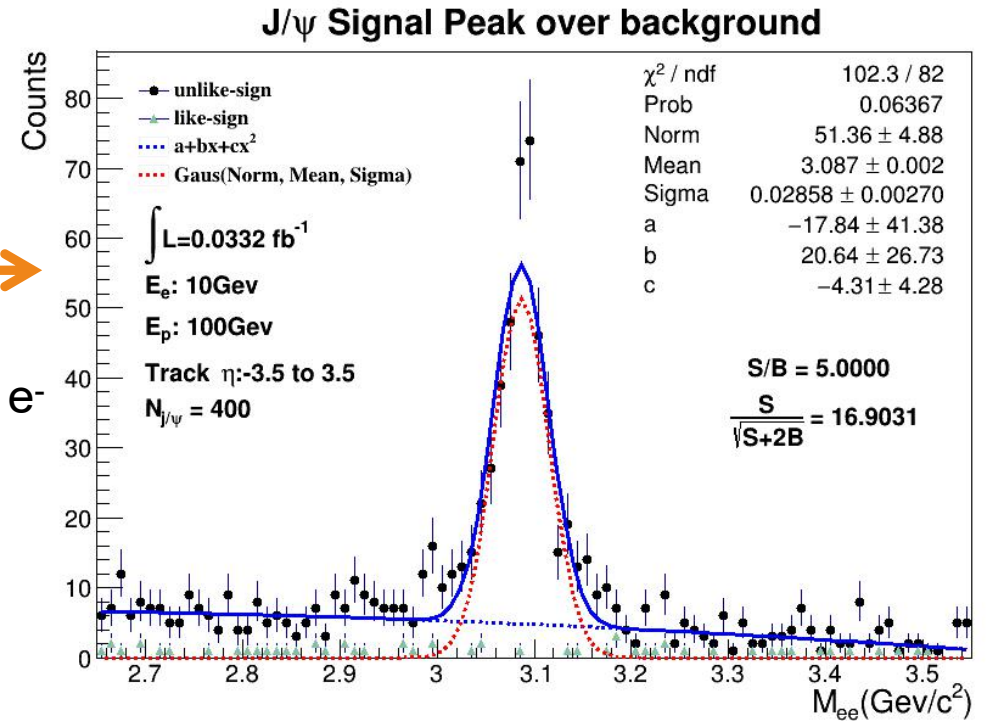
HPDIRC



J/ψ Reconstruction

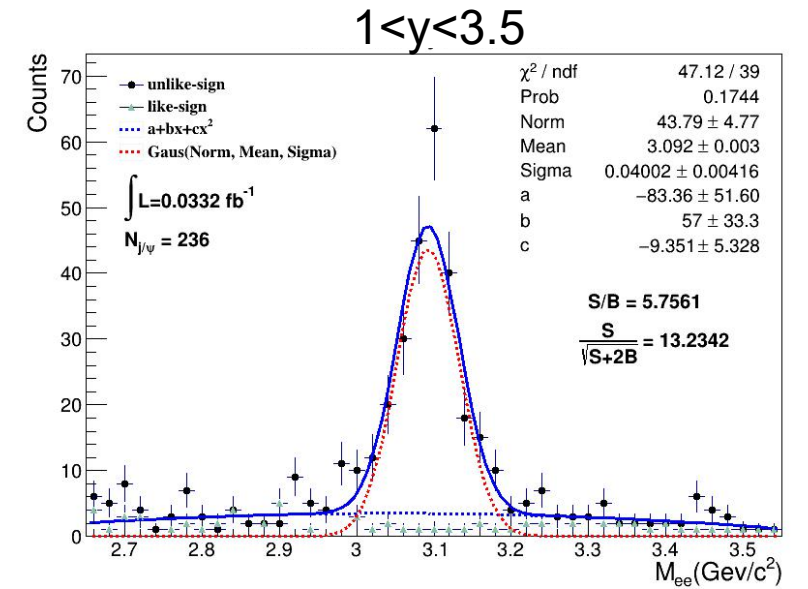
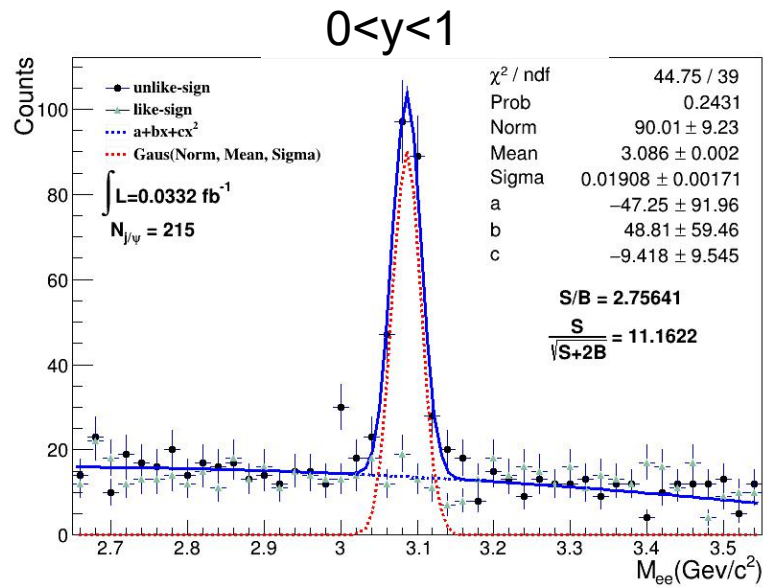
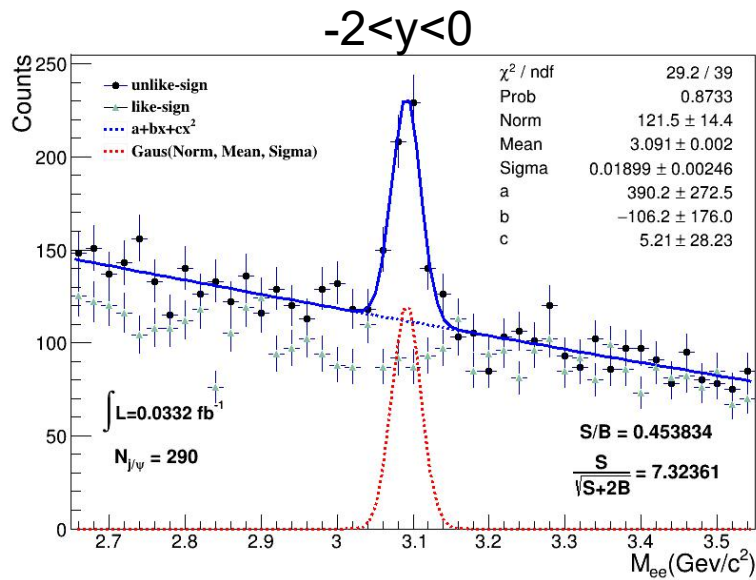


“veto”
 →
 exclude the
 beam scattering e-



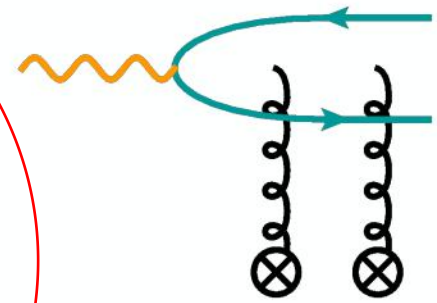
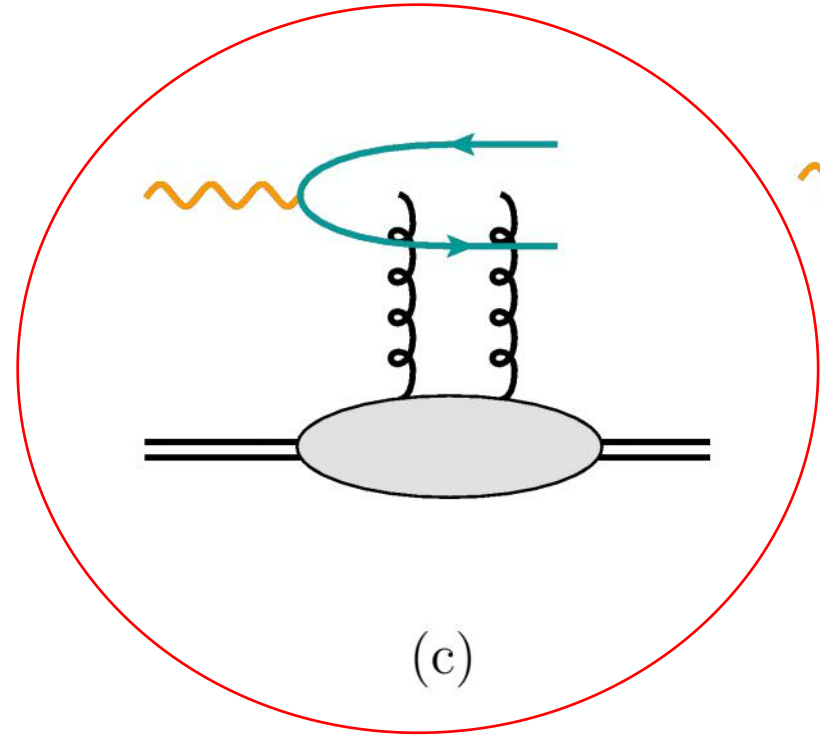
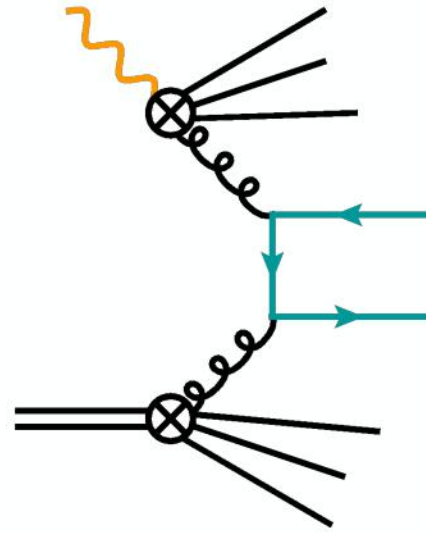
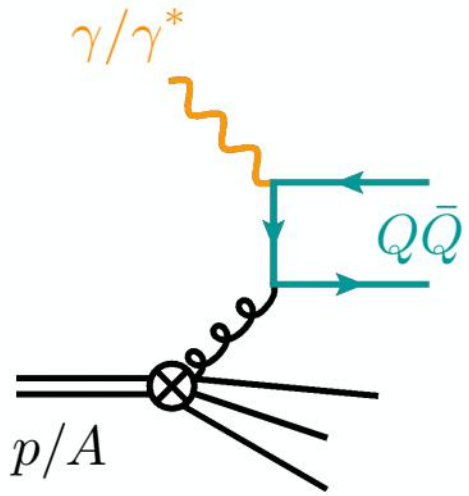
generator: pythia6 (eRHIC tuned) Full Geant simulation (fun4All)
 events: ~20million

J/ψ Reconstruction



Resolution and significance in different rapidity regions

Production mechanism of quarkonia

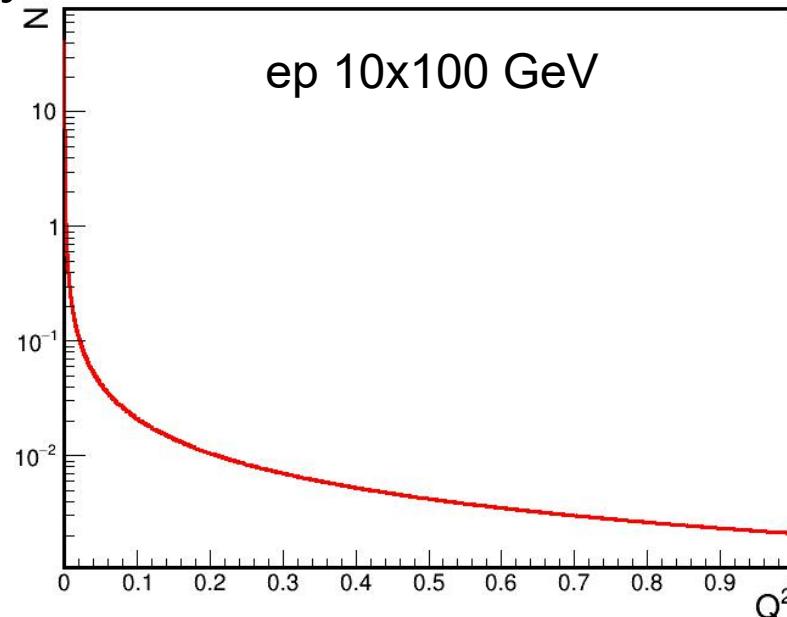


Direct, Resolve and **Exclusive** process

Yellow Report Fig. 7.95

Event feature of near threshold J/ψ photoproduction

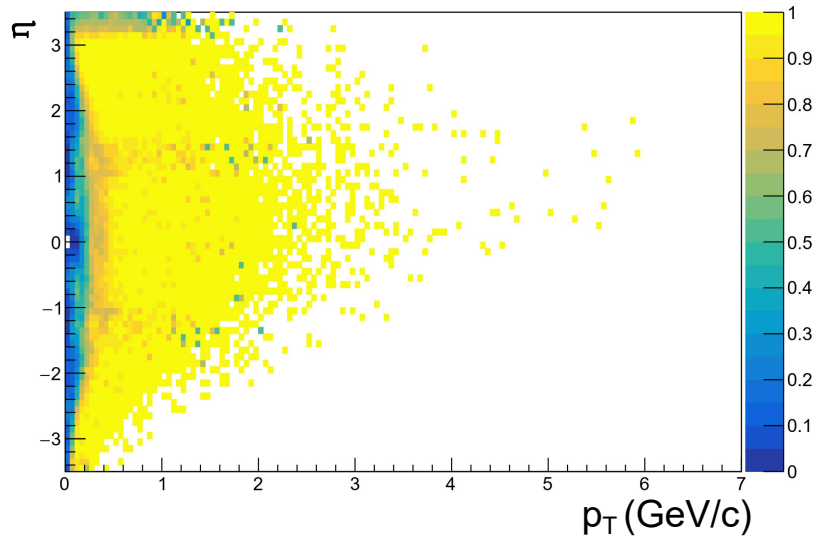
- ✓ The scattered proton or nucleus escape undetected down the beampipe at small scattering angles
- ✓ The majority of scattered electron escape undetected (Veto on $Q^2 > 1(\text{GeV}/c^2)^2$)
- ✓ No other event activity expect the electron-positron pair from J/ψ decay



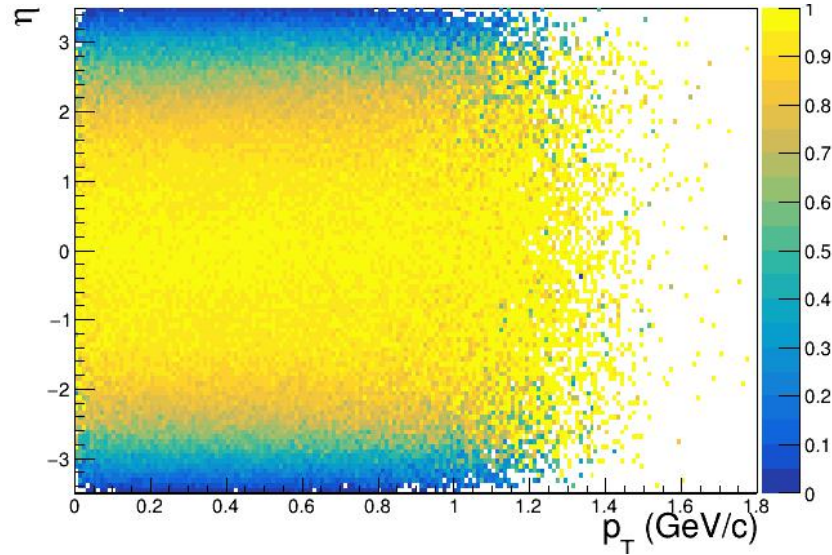
The trigger efficiency is sufficient high, assumed to be 1

Efficiency and S/B correction

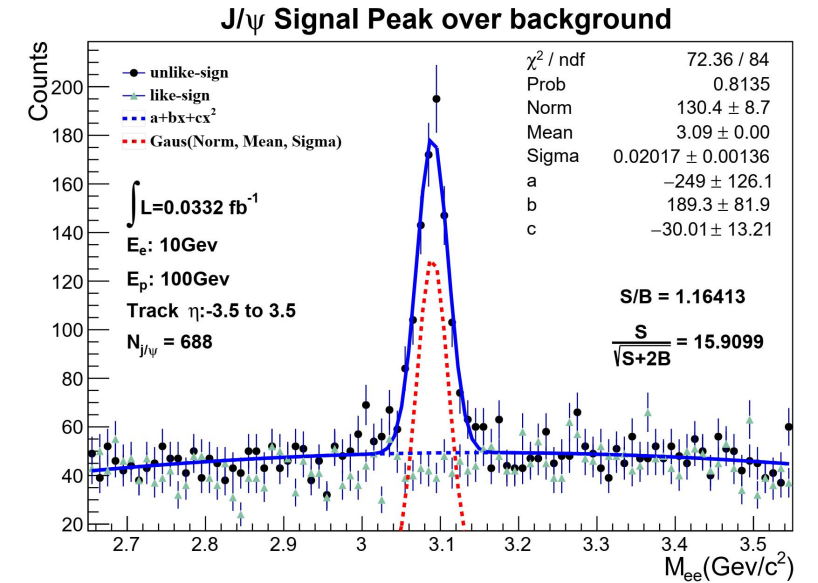
Single electron efficiency



J/ψ efficiency



S/B correction



The theoretical setup for projection (eSTARLight)

$$\sigma(eA \rightarrow eAV) = \int \frac{dW}{W} \int dk \int dQ^2 \frac{d^2 N_\gamma}{dk dQ^2} \sigma_{\gamma^* A \rightarrow VA}(W, Q^2)$$

$$\frac{d^2 N_\gamma}{dk dQ^2} = \frac{\alpha}{\pi k Q^2} \left[1 - \frac{k}{E_e} + \frac{k^2}{2E_e^2} - \left(1 - \frac{k}{E_e} \right) \left| \frac{Q_{\min}^2}{Q^2} \right| \right]$$

$$\sigma_{\gamma^* A \rightarrow VA}(W, Q^2) = f(M_V) \sigma(W, Q^2 = 0) \left(\frac{M_V^2}{M_V^2 + Q^2} \right)^n \quad n = c_1 + c_2(Q^2 + M_V^2),$$

$$\sigma(W, Q^2 = 0) = \int_{t_{\min}}^{\infty} dt \left. \frac{d\sigma(\gamma A \rightarrow VA)}{dt} \right|_{t=0} |F(t)|^2$$

Can be related to the cross section for $\sigma(\gamma+p \rightarrow V+p)$

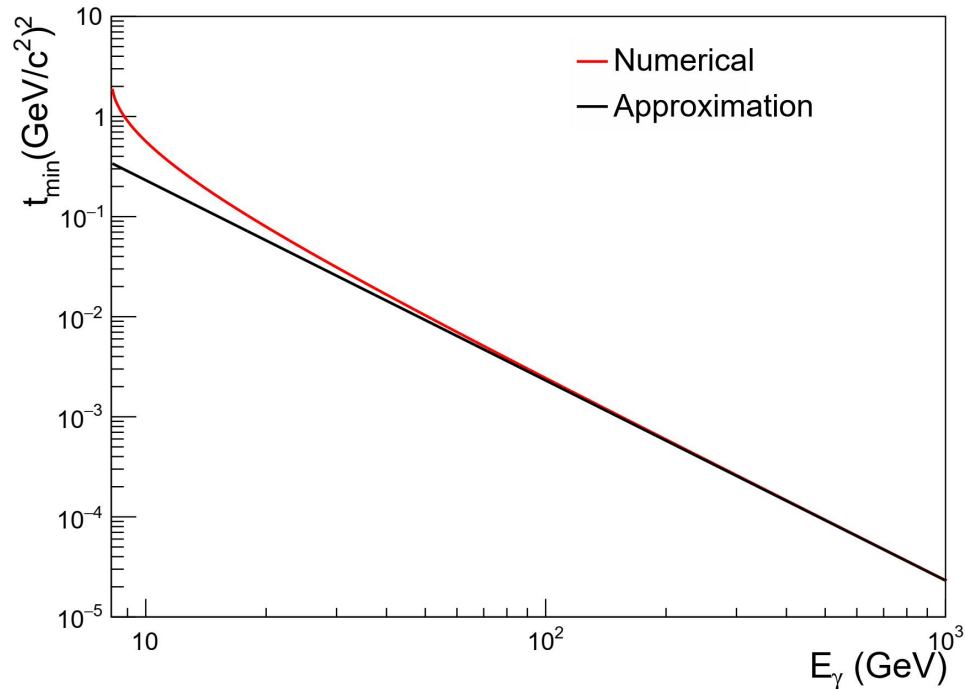
eSTARLight: Michael Lomnitz and Spencer Klein, Phys. Rev. C **99** (2019) 015203

Wangmei Zha et al, Phys. Rev. C **97** (2018) 044910

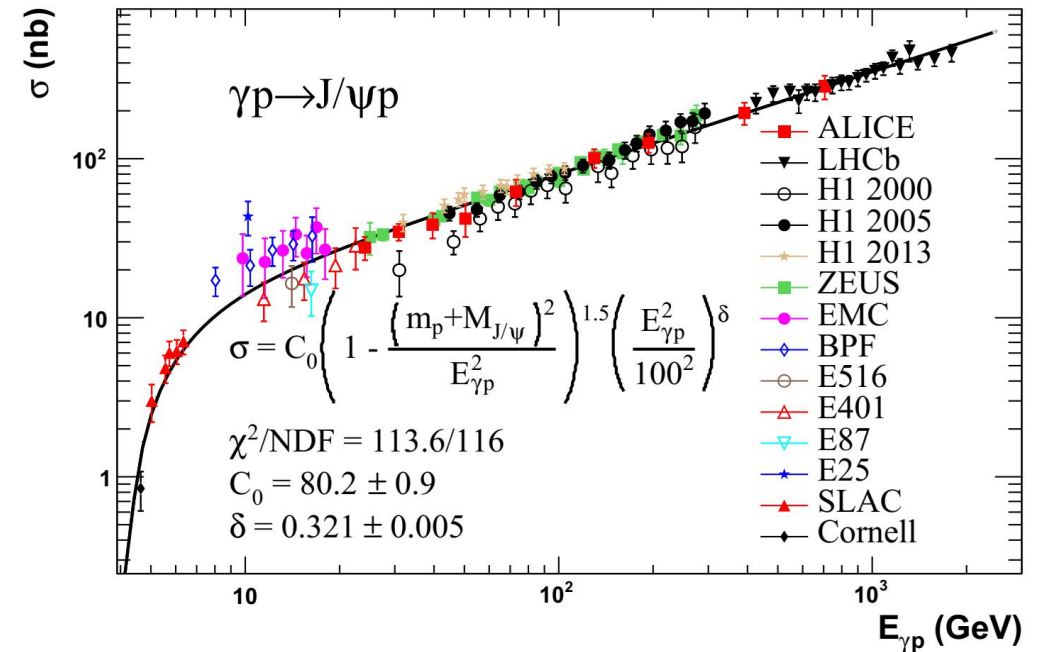
Two improvements for eSTARLight

Minimum momentum transfer

$$t_{\min} = (M_V^2/2k)^2 \text{ Approximation}$$

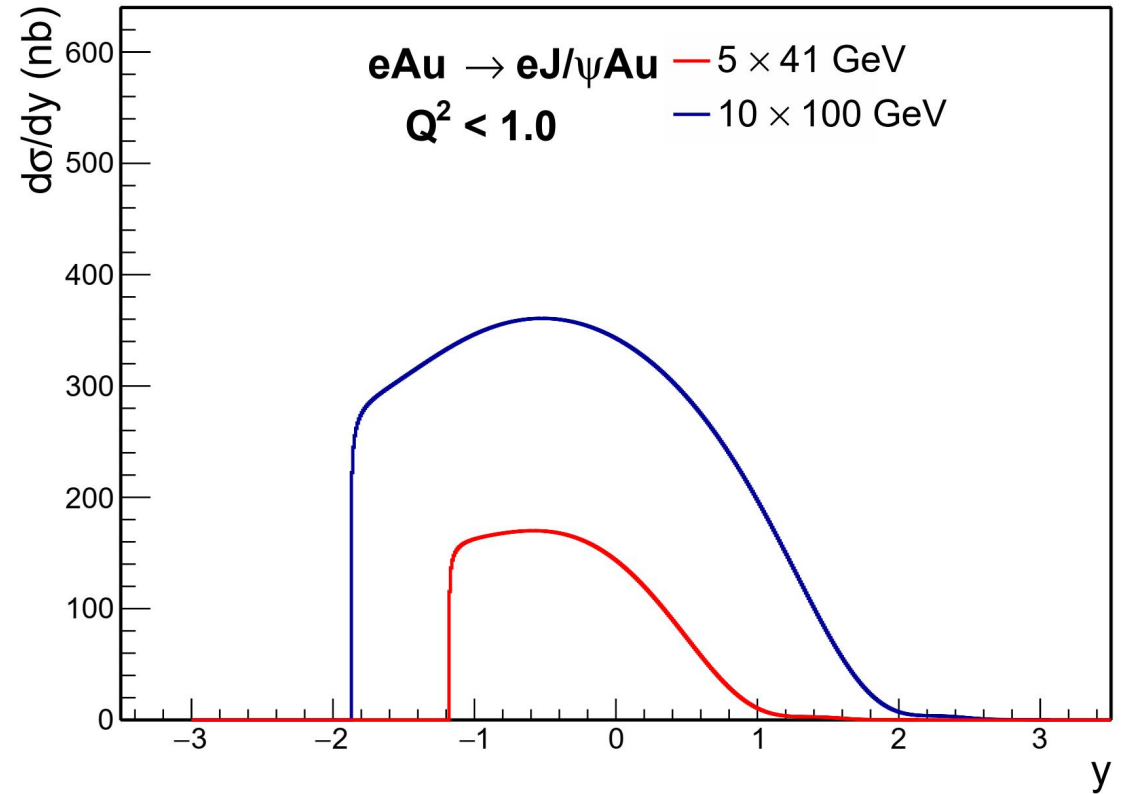
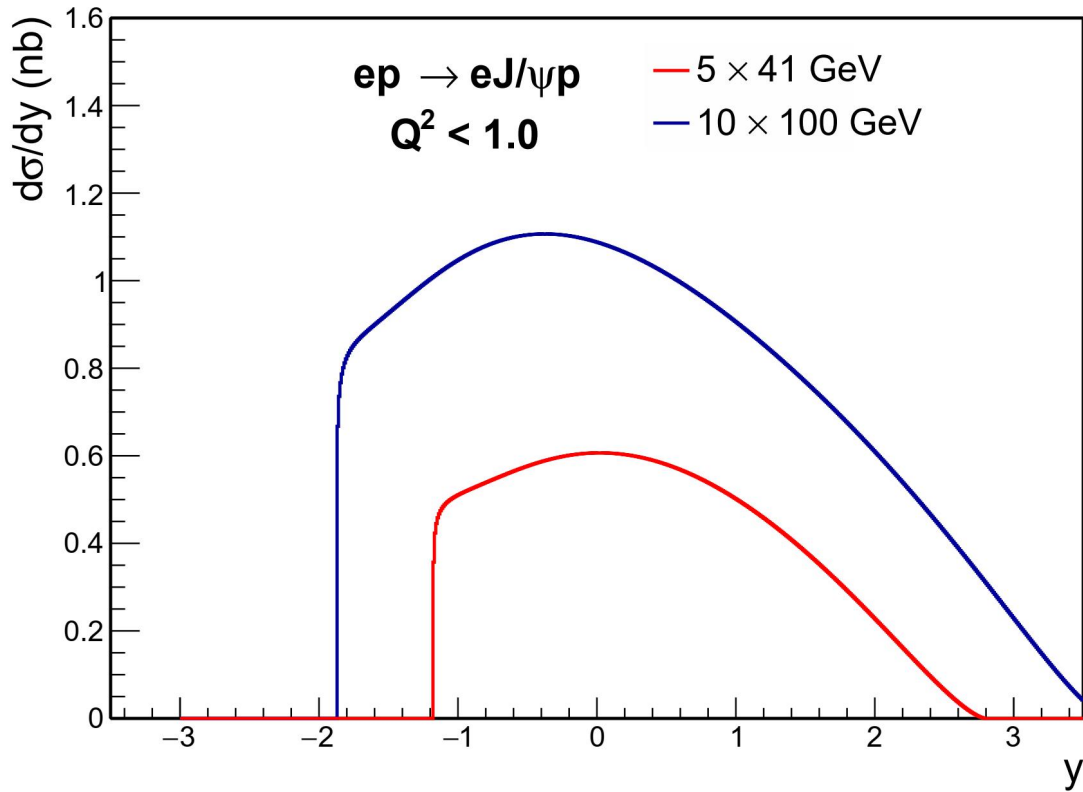


Parametrization for cross section input

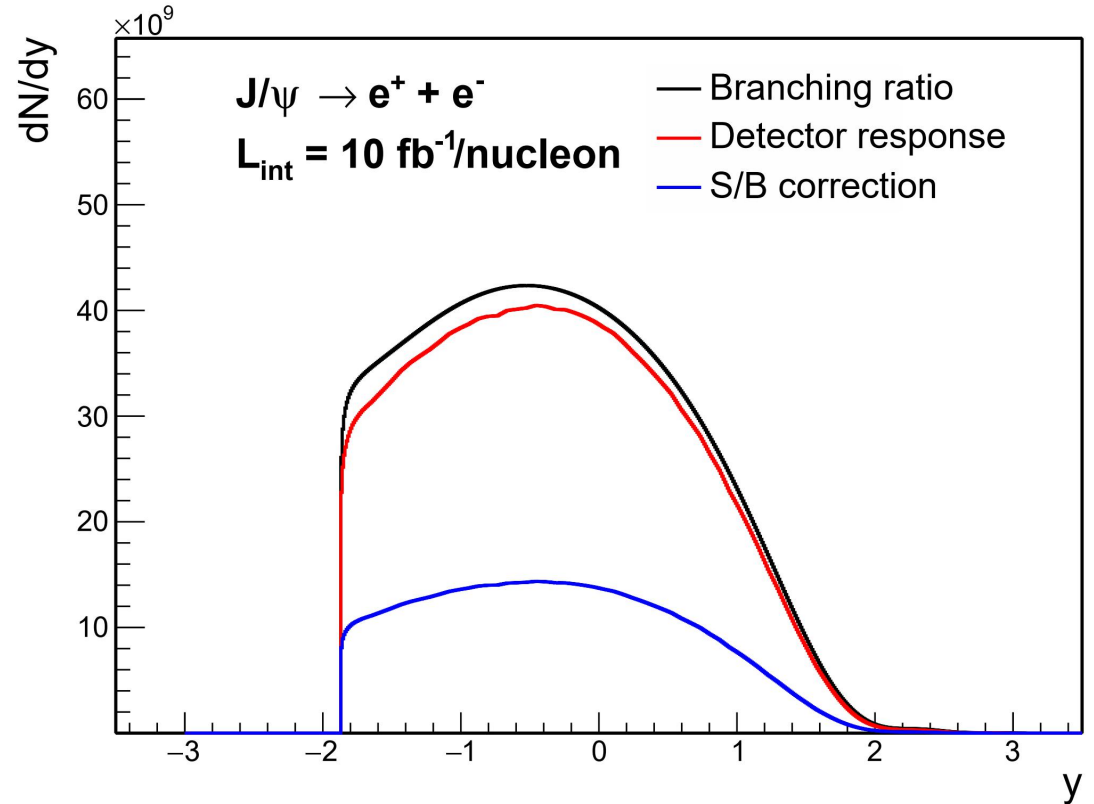
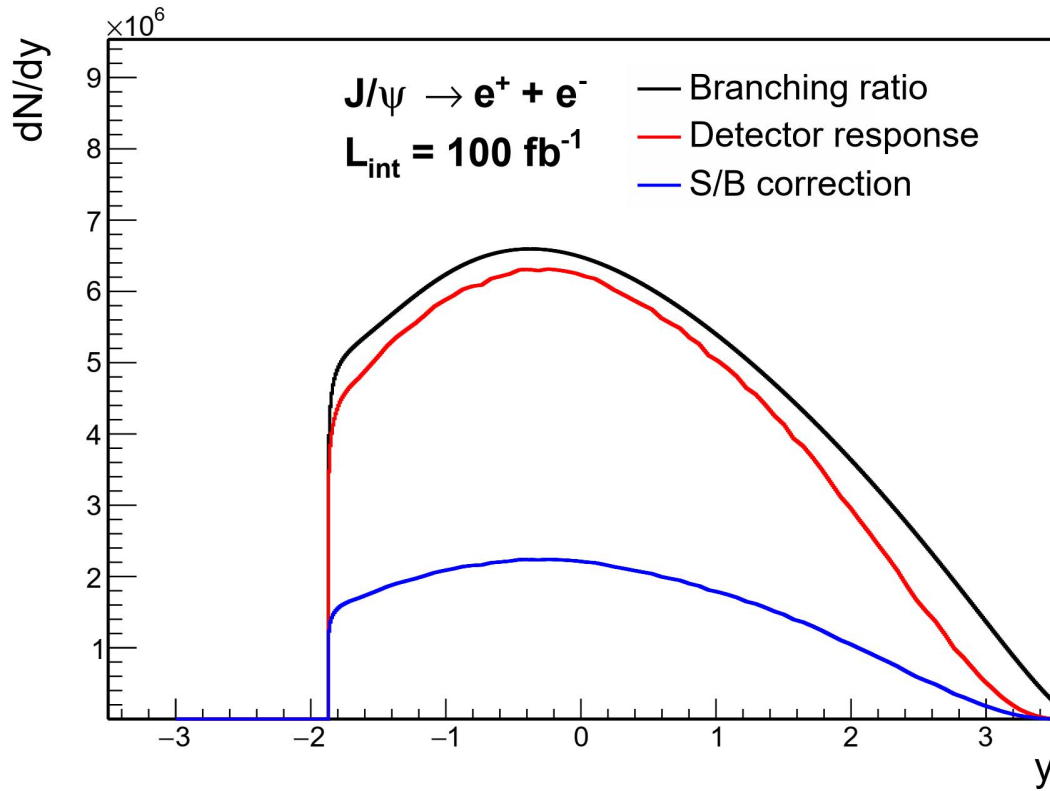


Z. Cao et al., Chin. Phys. C43 (2019) 064103

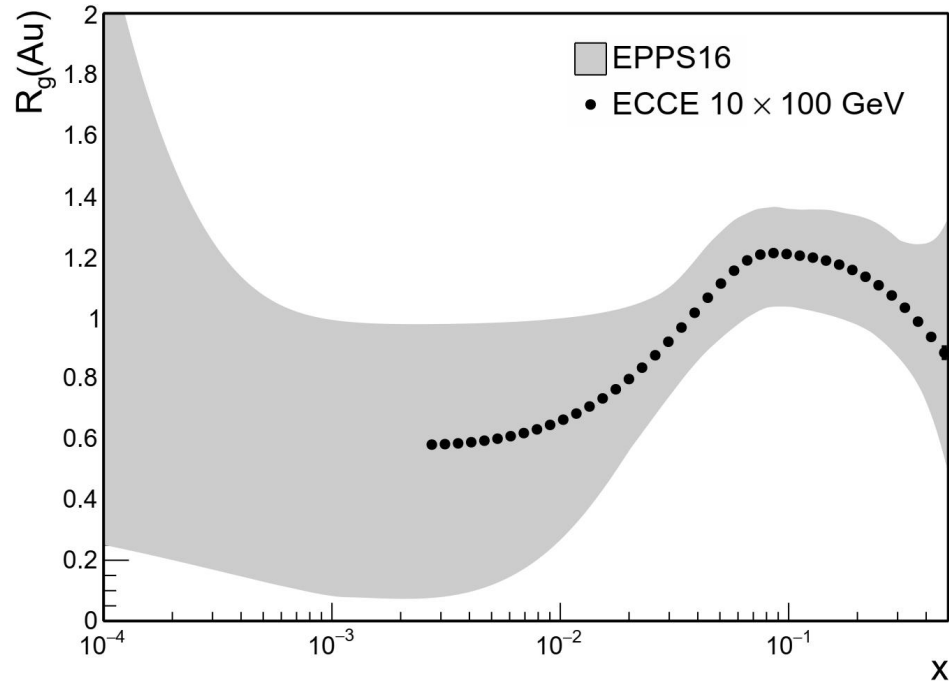
The theoretical input for ep and eAu



The projected statistics



The gluon nPDF projection



$$\left. \frac{d\sigma(\gamma A \rightarrow V A)}{dt} \right|_{t=0} = \frac{\alpha_s^2 \Gamma_{ee}}{3\alpha M_V^5} 16\pi^3 [xG_A(x, Q^2)]^2$$

$$x = \frac{M_V e^{\pm y}}{\sqrt{s}} \quad Q^2 = M_V^2/4$$

Physics Opportunities with Heavy Quarkonia at the EIC

If there is no shadowing

$$d\sigma_{\gamma A}/dt|_{t=0} = d\sigma_{\gamma p}/dt|_{t=0} \times A^2$$

Some remarks

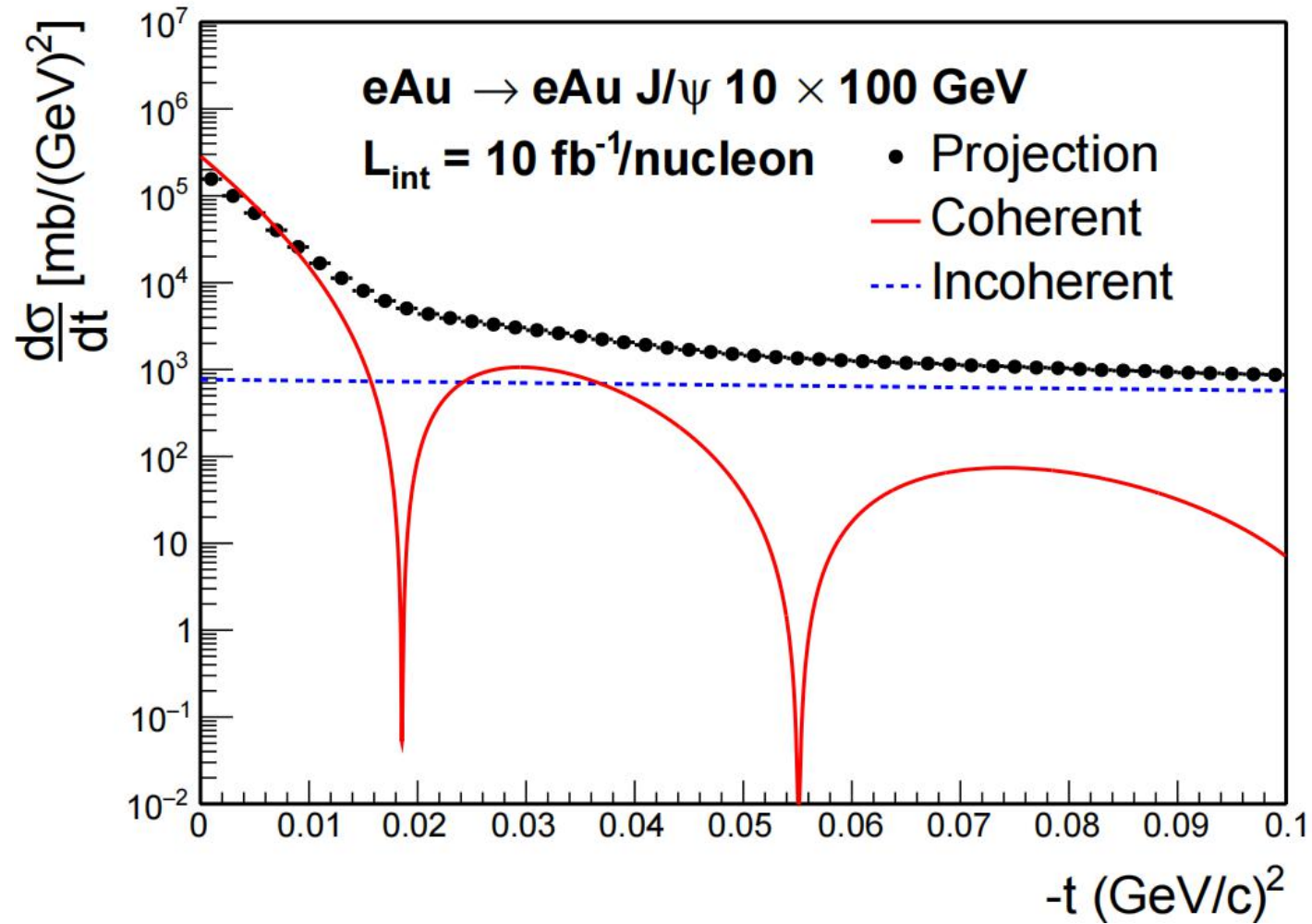
$$\rho(x) = \tan\left(\frac{\pi \Delta_{\mathbf{P}}}{2}\right)$$

$$R_{\text{skewed}} = \frac{2^{2\Delta_{\mathbf{P}}+3}}{\sqrt{\pi}} \cdot \frac{\Gamma(\Delta_{\mathbf{P}} + 5/2)}{\Gamma(\Delta_{\mathbf{P}} + 4)}$$

$$K = (1 + \rho^2(x)) \cdot R_{\text{skewed}}^2$$

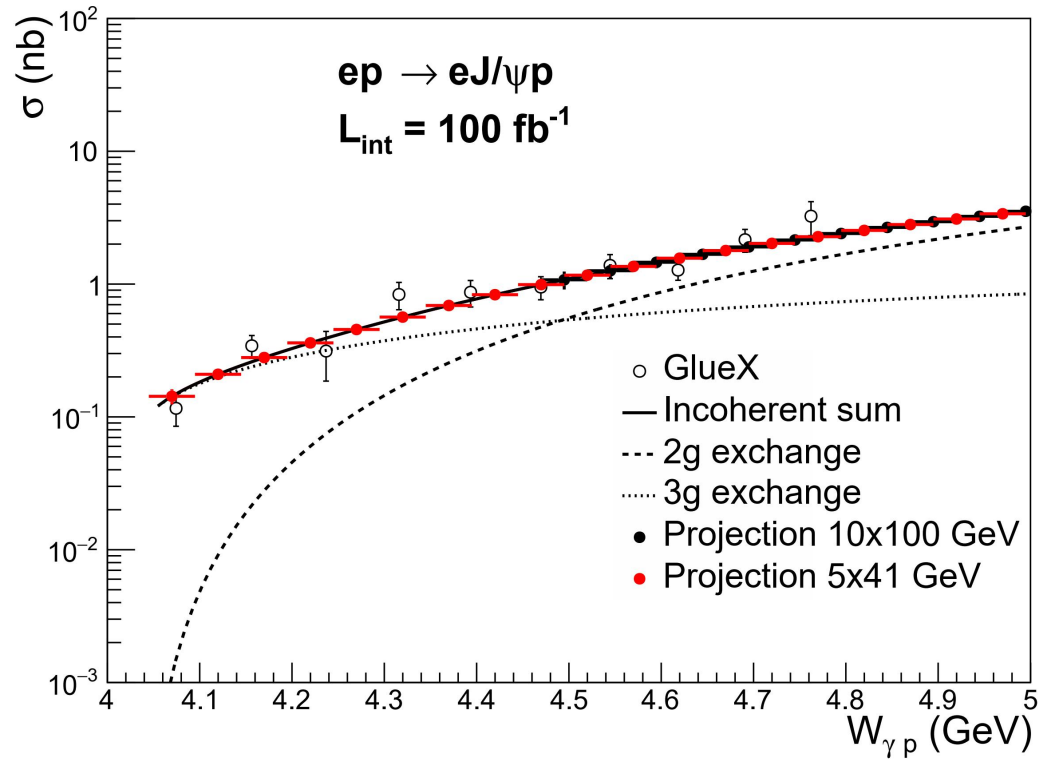
Assume the correction is the same for p and Au

The t distribution projection



The momentum resolution would wipe out the diffraction dips.

The near threshold production mechanism



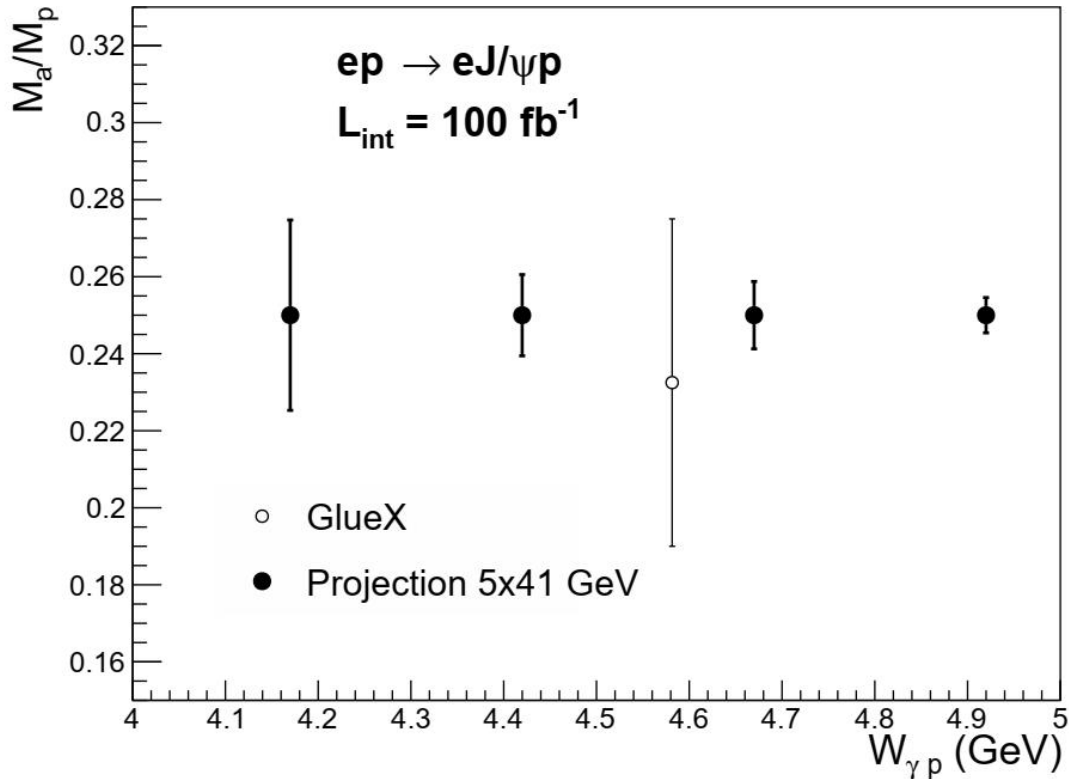
Phys. Rev. Lett. 123, 072001(2019)

$$\frac{d\sigma}{dt} = \mathcal{N}_{2g} v \frac{(1-x)^2}{R^2 \mathcal{M}^2} F_{2g}^2(t) (s - m_p^2)^2$$

$$\frac{d\sigma}{dt} = \mathcal{N}_{3g} v \frac{(1-x)^0}{R^4 \mathcal{M}^4} F_{3g}^2(t) (s - m_p^2)^2$$

Physics Letters B 498 (2001) 23–28

The trace anomaly parameter projection



$$M_q = \frac{3}{4} \left(a - \frac{b}{1 + \gamma_m} \right) M_N,$$

$$M_g = \frac{3}{4} (1 - a) M_N,$$

$$M_m = \frac{4 + \gamma_m}{4(1 + \gamma_m)} b M_N,$$

$$M_a = \frac{1}{4} (1 - b) M_N,$$

Eur. Phys. J. C (2020) 80:507

Extract the QCD trace anomaly parameter b

Summary and future plan

Summary

- Excellent electron identification probability at ECCE
- Reasonable mass resolution for J/ψ reconstruction
- High statistics to study the photoproduction of J/ψ

Future plan

- The projection results from full Geant simulation
- The projection results for Quarkonia production mechanism



Backup

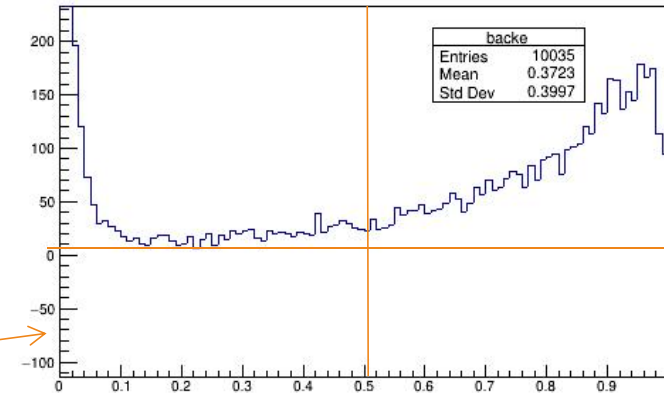
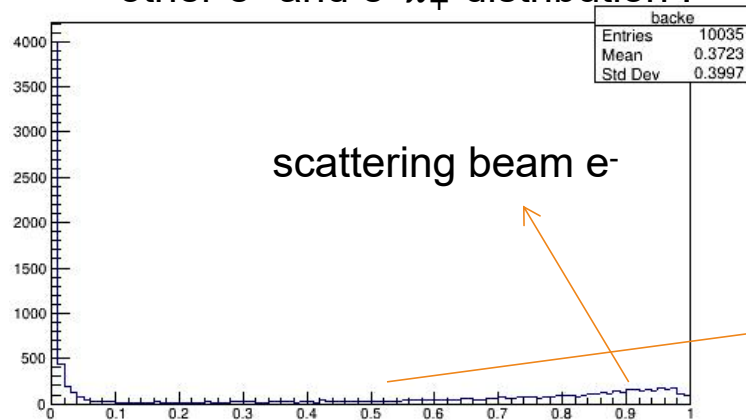
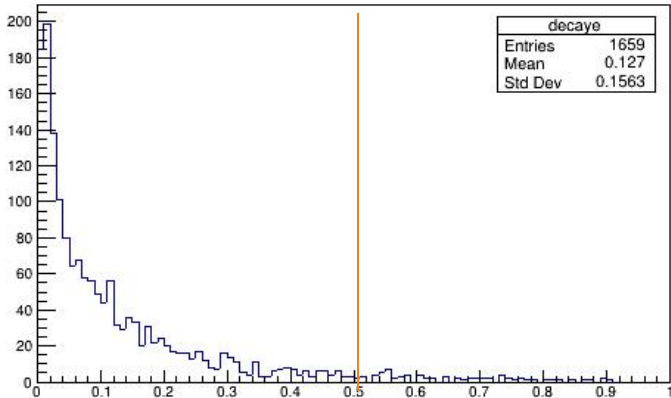
J/ψ detection (For Page15 S/B)

a forward light cone variables can be used to see scattering beam e⁻ influence

$$x_+ = \frac{b_0 + (-b_z)}{a_0 + (-a_z)} \text{ (cause beam } e^- \text{ moves along negative } z \text{ axis), } b \text{ is beam } e^-.$$

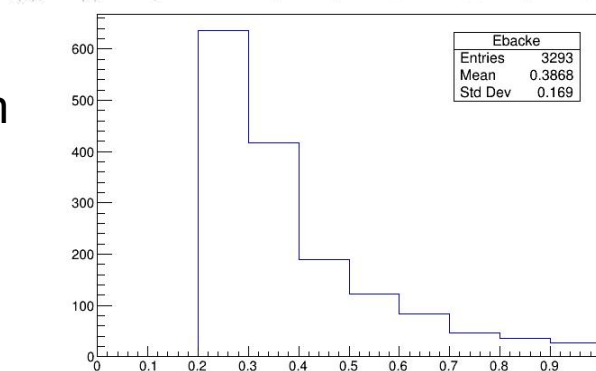
e⁺ and e⁻ of J/ψ decay x₊ distribution

other e⁺ and e⁻ x₊ distribution :

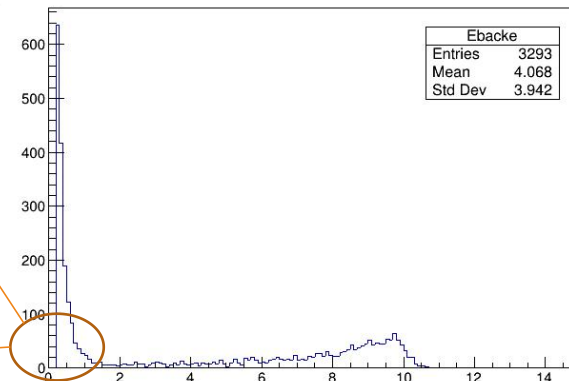


cut: $x_+ < 0.5$

influence of e⁻ from light hadron decay

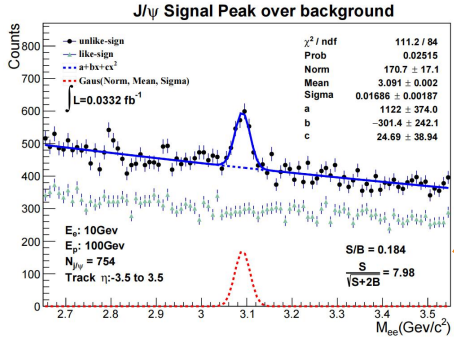


e⁺ and e⁻ of J/ψ decay x₊ distribution



other e⁺ and e⁻ x₊ distribution :

J/ψ detection (For Page15 S/B)



cut: $x_+ < 0.5$
S/B: 0.184 → 0.888

cut: E>0.6
S/B: 0.888 → 1.164

