

# How To Measure Absolute Luminosity

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References :

- [1 ] AN184, S. Belikov *et al.*
- [2 ] “Luminosity Measurements and Calculations”, K. Potter, CERN Particle Accelerator School Lectures, 1985.
- [3 ] “Luminosity Optimisation for Storage Rings with Low- $\beta$  Sections and Small Crossing Angles”, E. Keil, NIM **113**, 333 (1973).

# Overview

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- What do we mean by Absolute Luminosity? (5%)
- Why do we need to measure it? (5%)
- How do we measure it? (90%)
  - What exactly is a Vernier Scan?
  - Description of the hardware involved (DCCTs, WCMs, BPMs)
  - Sample Data

## What do we mean by Absolute Luminosity?

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- Luminosity is proportionality between event rate  $\dot{N}$  of process with cross section  $\sigma$  :

$$\dot{N} = \mathcal{L}\sigma$$

- Typically given in  $\text{cm}^{-2}\text{s}^{-1}$
- $N_y$  target particles uniformly distributed over area  $A$  of zero thickness
- Probability of single beam particle colliding head-on is  $N_y \times \sigma/A$
- If  $N_b$  beam particles, total number of collisions  $N = N_b \times N_y \times \sigma/A$   
 $\Rightarrow$  For  $k_b$  bunch pairs intersecting at frequency  $f_{rev}$ , total event rate :

$$\dot{N} = \left[ k_b \times f_{rev} \times \frac{N_b \times N_y}{A} \right] \times \sigma$$
$$\Rightarrow \mathcal{L} = k_b \times f_{rev} \times \frac{N_b \times N_y}{A}$$

For zero-length bunches with non-uniform charge distributions :

$$d\mathcal{L} = k_b \times f_{rev} \times \frac{\rho_b(x, y)dxdy \times \rho_y(x, y)dxdy}{dxdy}$$

$$\Rightarrow \mathcal{L} = k_b \times f_{rev} \times \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \rho_b(x, y)\rho_y(x, y)dxdy$$

## What do we mean by Absolute Luminosity?

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For Gaussian beams with  $N_b$  and  $N_y$  particles colliding head-on :

$$\begin{aligned}\mathcal{L} &= k_b \times f_{rev} \times \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \rho_b(x, y) \rho_y(x, y) dx dy \\ &\Rightarrow k_b \times f_{rev} \times \frac{N_b \times N_y}{2\pi \sqrt{(\sigma_{bx}^2 + \sigma_{yx}^2)(\sigma_{by}^2 + \sigma_{yy}^2)}}\end{aligned}$$

- Luminosity probably 2nd most important machine parameter, after beam energy

## Why do we need to measure Absolute Luminosity?

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Many PHENIX papers discuss the measurement of an invariant cross-section :

$$E \frac{d^3\sigma}{dp^3} \approx \frac{N}{\mathcal{L}} \frac{1}{2\pi p_T} \frac{1}{\Delta p_T \Delta y} \frac{1}{\eta_{trig}} \frac{1}{\eta_{recon}}$$

- We usually need  $\mathcal{L}$  to extract cross-sections
- Knowing  $\sigma$  and  $\mathcal{L}$  useful for calibrating detector efficiencies, geometric acceptances, ...
- Required to calibrate detectors (BBC, ZDC) which can then be used as luminosity monitors  $\dot{N}_{monitor} = \mathcal{L} \sigma_{monitor} \Rightarrow \mathcal{L} = \dot{N}_{monitor} / \sigma_{monitor}$
- Important feedback for CAD for tuning beam, evaluating performance

## How do we measure Absolute Luminosity?

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$$\mathcal{L} \approx k_b \times f_{rev} \times \sum_{i=1}^{k_b} \frac{N_b^i \times N_y^i}{2\pi \sqrt{(\sigma_{bx}^{i2} + \sigma_{yx}^{i2})(\sigma_{by}^{i2} + \sigma_{yy}^{i2})}}$$

- To measure the number of ions  $N_b$  and  $N_y$  in each colliding bunch pair we use :
  - DCCT = D.C. current transformer
  - WCM = Wall Current Monitor
  - $k_b \approx 60$ ,  $f_{rev} \approx 78.2$  kHz
- Beam overlap integral (effective width and height) determined with vernier scan
  - Requires knowledge of event rates (BBC, ZDC) versus beam displacement
  - Beam displacement measured with Beam Position Monitors (BPMs)

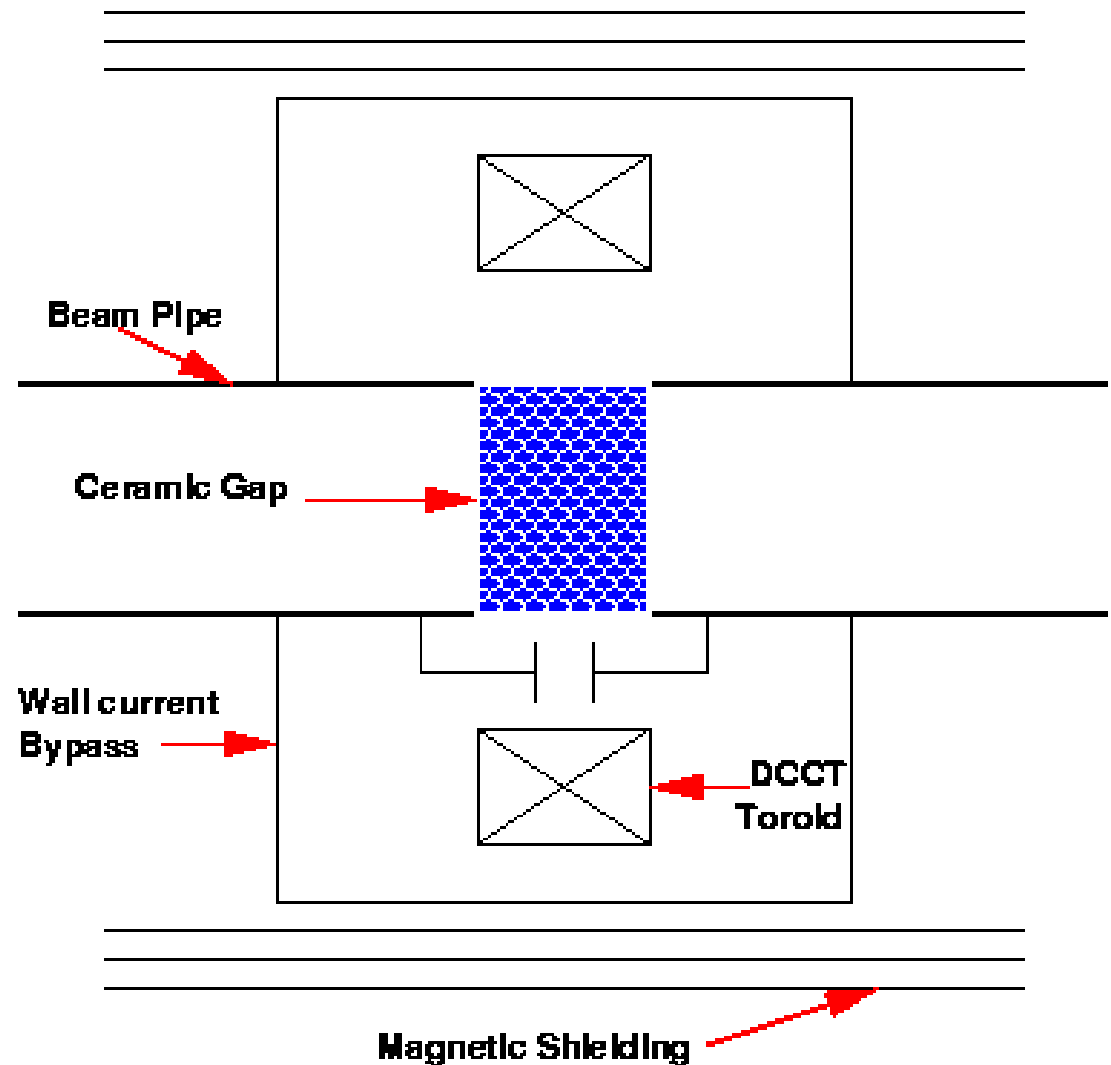
## Beam Ion measurement with the DCCTs

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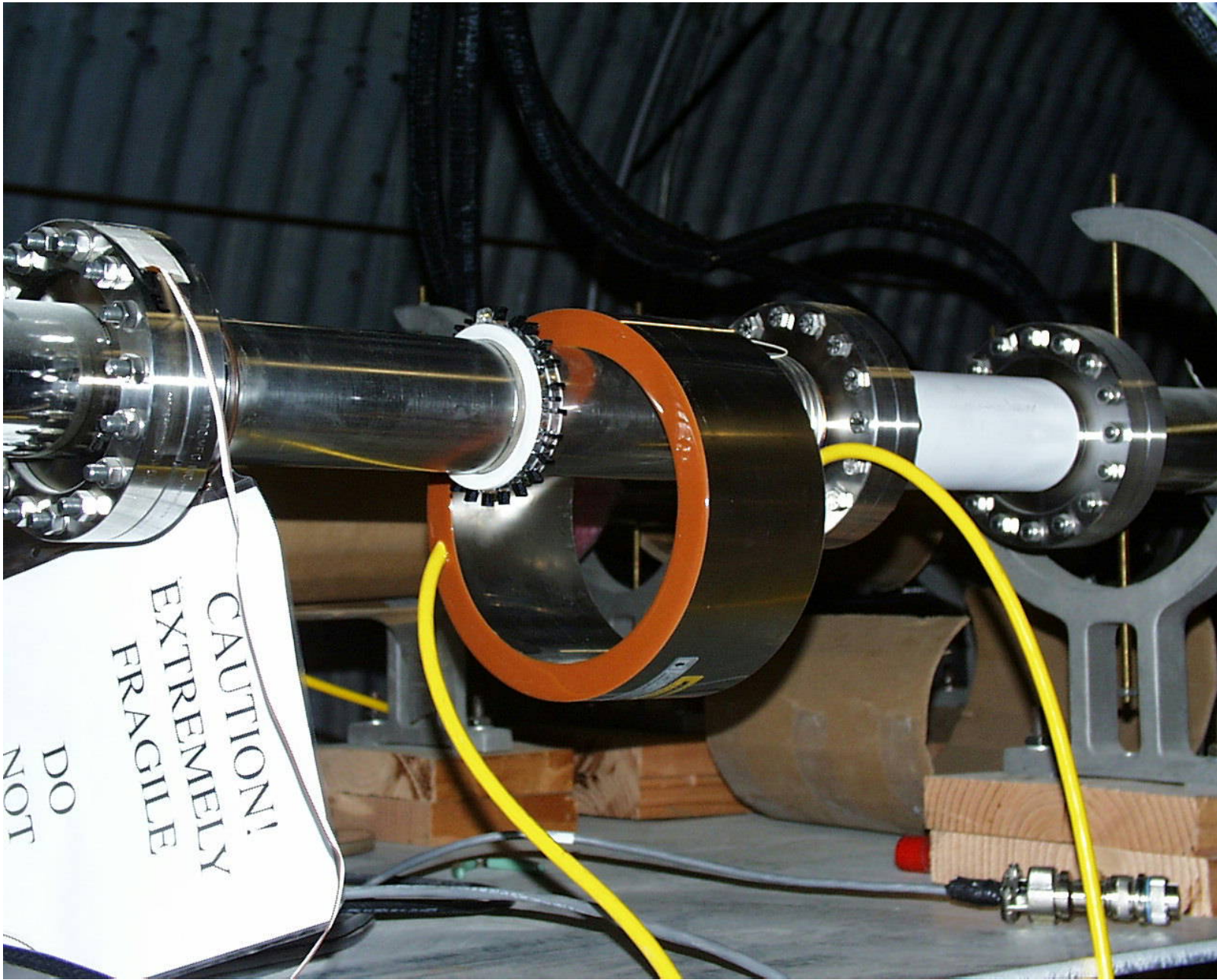
- D.C. current transformer designed in 1969 by Klaus Unser for the ISR
- One DCCT for blue and one for yellow at 2 o'clock
- Consists of high permeability toroid around beam with three windings
- One set of coils drives toroid into positive then negative saturation
- Sense windings detect flux changes, produce symmetric positive and negative voltage
- If beam current  $\neq 0$ , acts like shift in  $H$ , coil driven to different parts of  $B - H$  curve - positive and negative swings no longer balanced
- Second harmonic of sense signal is feedback to third winding to cancel flux from beam
- Voltage across precision resistor in series with third winding is proportional to beam current
- Measures average current over integration period of  $\approx 1$  second (hence DC) with  $\approx 0.2\%$  absolute accuracy,  $50 \mu\text{A}$  rms resolution,  $\pm 0.01\%$  linearity error
- Long integration period means insensitive to bunched/unbunched beam, can't measure charge in individual bunches
- Calibrated with precision current source at startup
- Demagnetized after each beam dump to minimize "memory" effects

# Beam Ion measurement with the DCCTs

- Insulating ceramic break in beam-pipe shunts wall currents around toroid
- Magnetic shielding excludes external magnetic fields
- Only beam current threads core



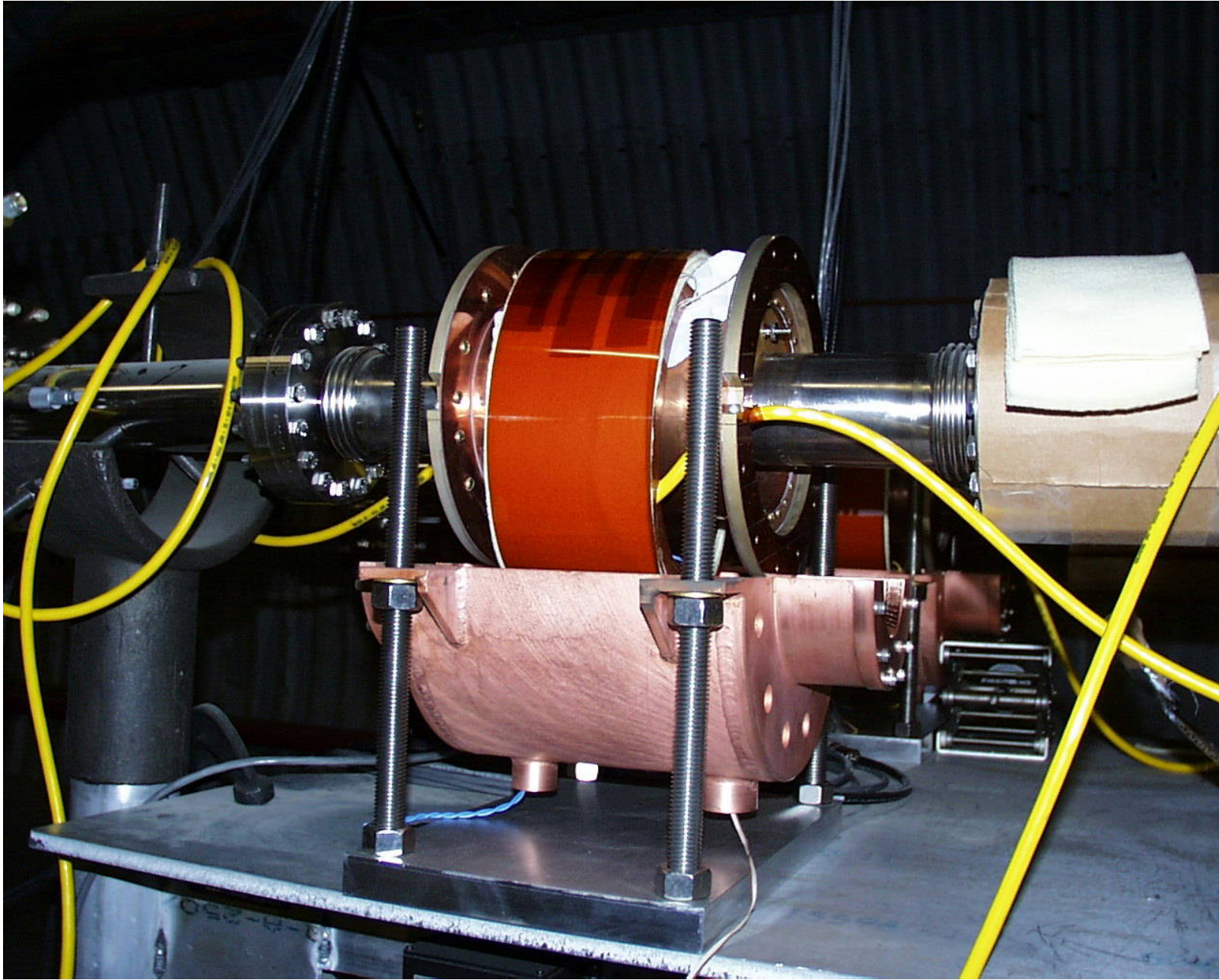
## Beam Ion measurement with the DCCTs





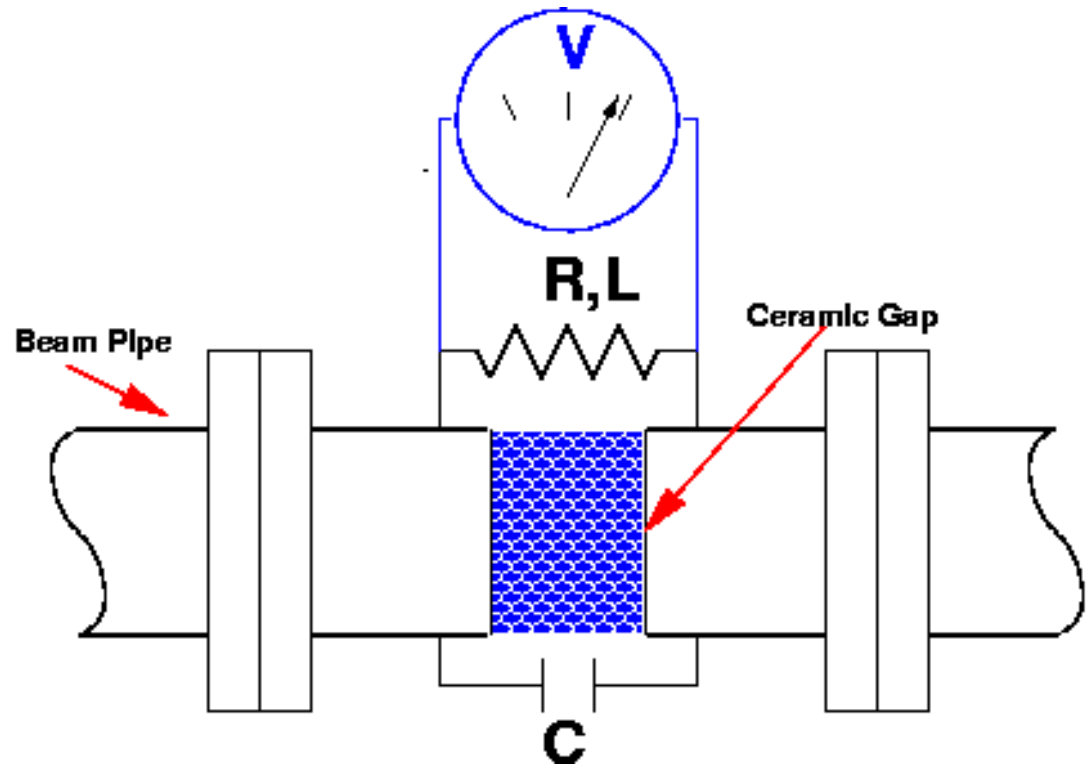
## Beam Ion measurement with the DCCTs

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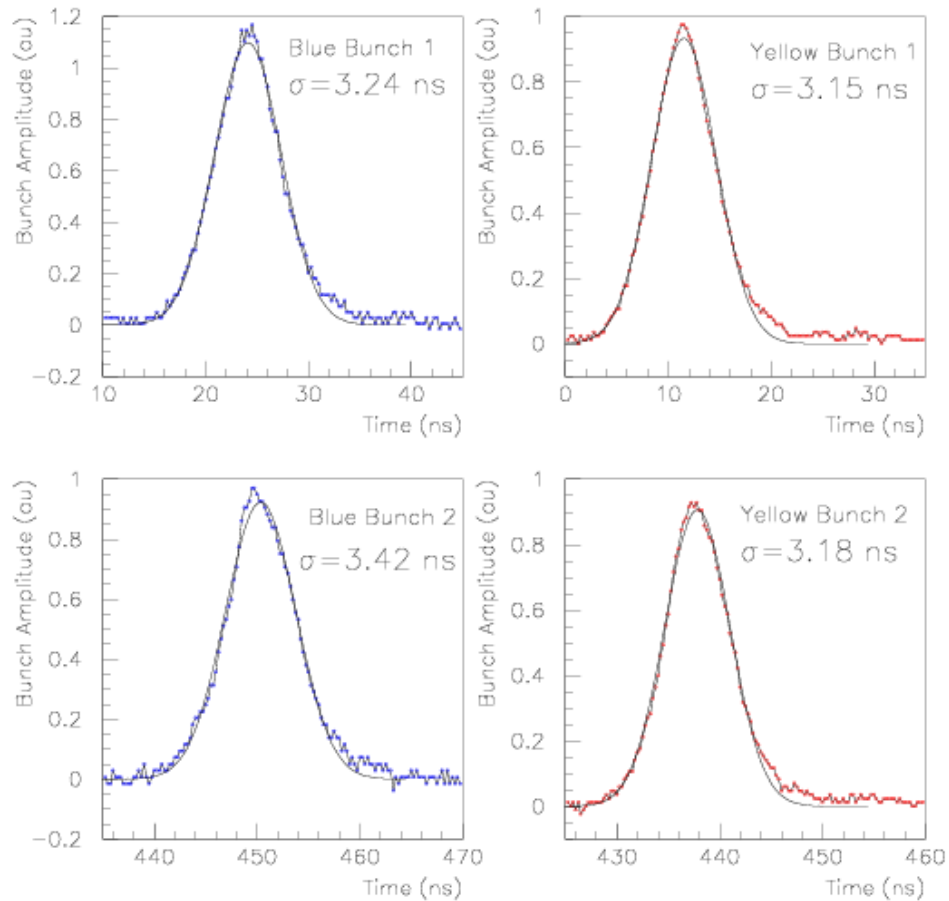


# Beam Ion measurement with the WCMs

- Insulating ceramic break in beam-pipe forces image wall currents through resistors
- $\omega_{\text{low}} \approx R/L$
- $\omega_{\text{high}} \approx 1/RC$ , few GHz
- Sensitive only to bunched beam
- Less accurate than DCCT
- Calibrate WCM with DCCT each fill after ramp when no unbunched beam
- Can measure bunch longitudinal profiles

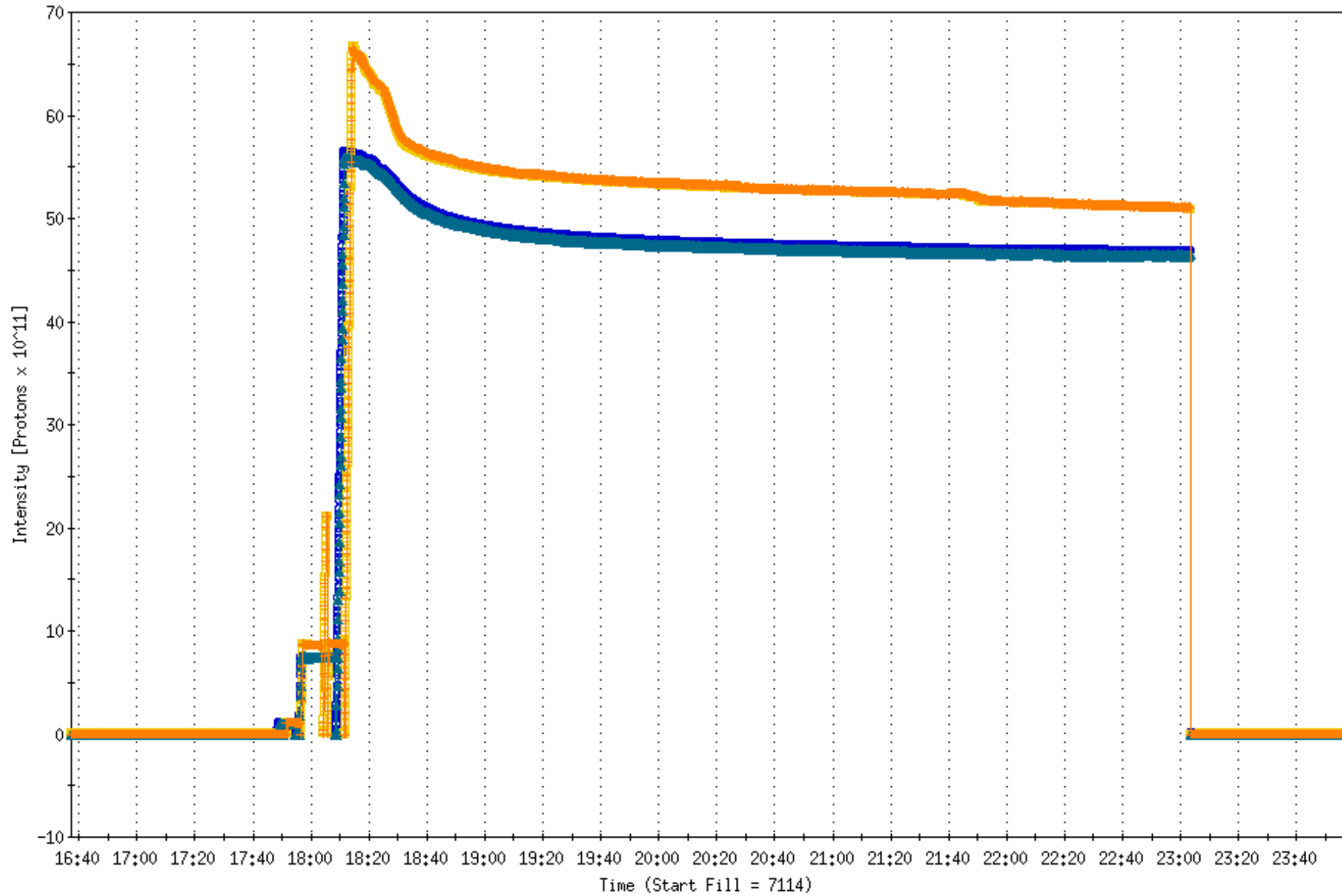


# Beam Ion measurement with the WCMs



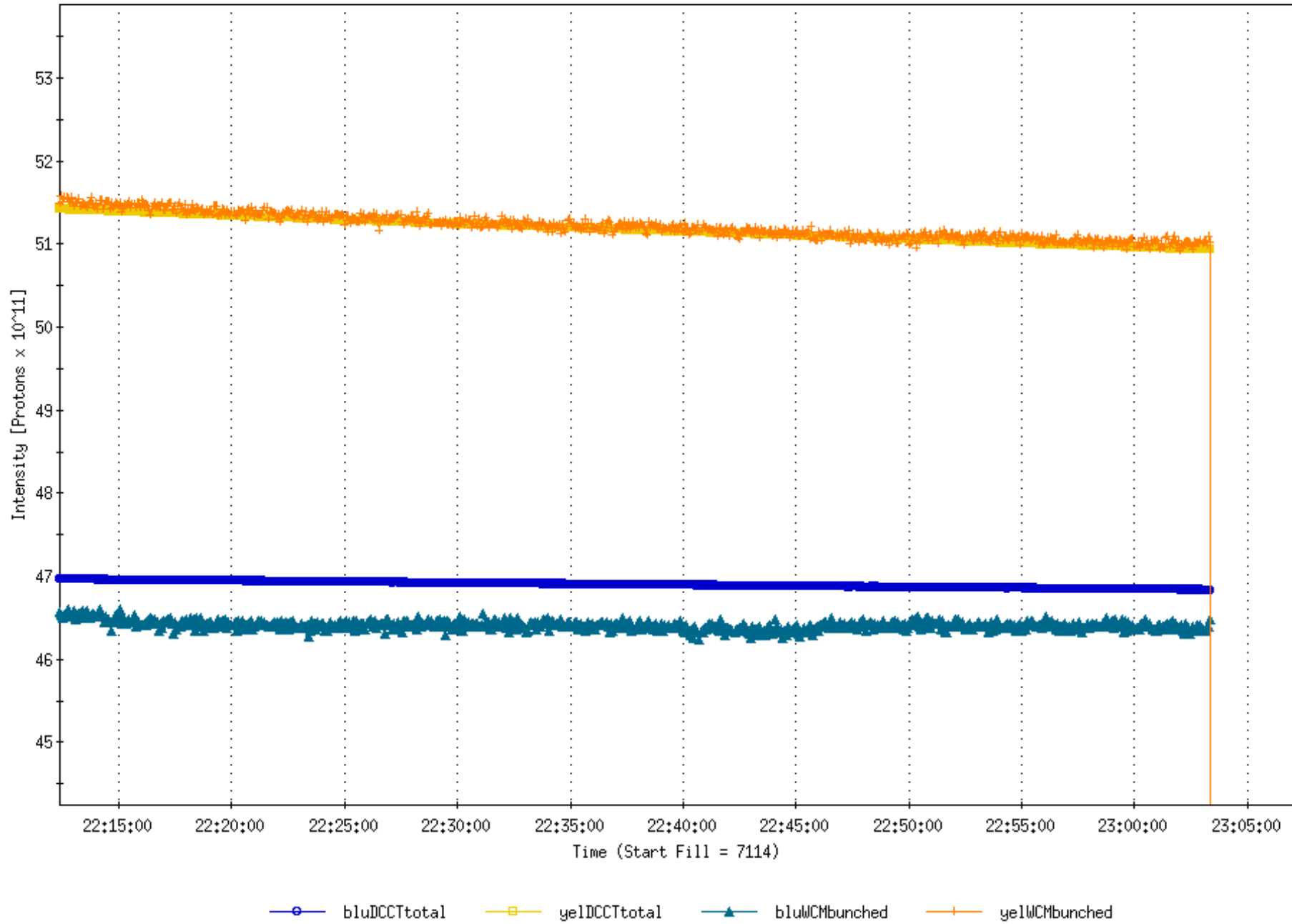
- Wall current monitor samples profile in 0.25 nsec bins
- Wide bunches seem to correspond to wide vertex
- 1st pair : 47 cm; 2nd pair : 51 cm
- Still to do : predict vertex dist. from WCM
- Alternating between bunches  $\rightarrow$  0.125 nsec bins  $\rightarrow$  improved vertex reconstruction

# Calibration of the WCMs with the DCCTs after Ramp



bluDCCTtotal    yelDCCTtotal    bluWCMbunched    yelWCMbunched

# WCMs exhibit more noise than DCCTs



## What is a Vernier Scan?

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Recall  $\mathcal{L}$  of identical zero-length bunch pairs with non-uniform charge distributions :

$$d\mathcal{L} = k_b \times f_{rev} \times \frac{\rho_b(x, y) dx dy \times \rho_y(x, y) dx dy}{dx dy}$$

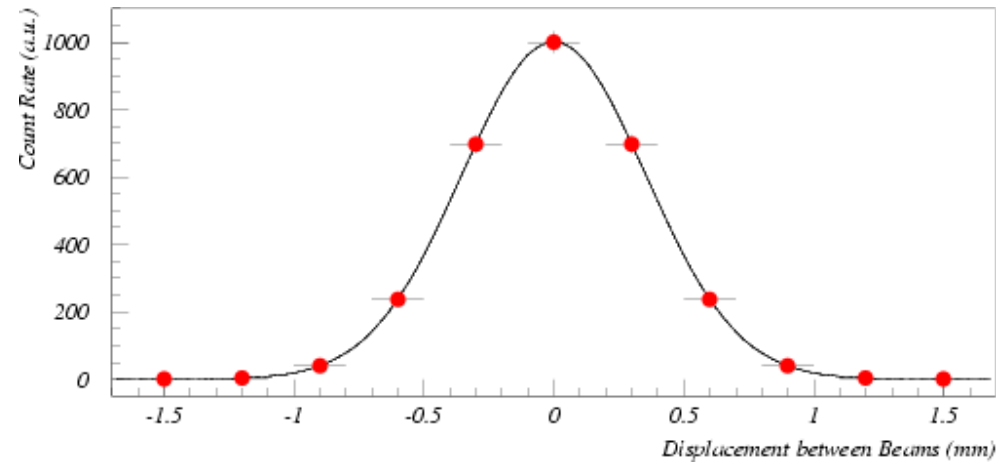
Let  $N_b = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \rho_b(x, y) dx dy$  be number of ions in blue bunch :

$$\begin{aligned} \mathcal{L} &= k_b \times f_{rev} \times N_b \times N_y \times \frac{\int \int \rho_b(x, y) \times \rho_y(x, y) dx dy}{\int \int \rho_b(x, y) dx dy \int \int \rho_y(x, y) dx dy} \\ &\Rightarrow k_b \times f_{rev} \times N_b \times N_y \times \frac{\int \rho_b(x) \times \rho_y(x) dx}{\int \rho_b(x) dx \int \rho_y(x) dx} \quad (\text{ribbon - like beams}) \end{aligned}$$

- Simon van der Meer showed how to measure the overlap integral in 1968 at ISR
- Detailed knowledge of each beam's transverse profile is not required

## What is a Vernier Scan?

- Consider event rate when beams displaced by  $h$  from each other :



Count Rate( $h$ ) =  $\eta \int \rho_b(x)\rho_y(x - h)dx$  so area under counting rate curve :

$$\begin{aligned}
 \text{Area} &= \eta \times \int \int \rho_b(x)\rho_y(x - h)dzdh \\
 &= \eta \times \int \rho_b(x) \left[ \int \rho_y(x - h)dh \right] dx \quad \text{but} \quad \left[ \int \rho_y(x - h)dh \right] = \int \rho_y(x)dx \\
 &= \eta \times \int \rho_b(x)dx \int \rho_y(x)dx
 \end{aligned}$$

- Maximum count rate at  $h=0 \Rightarrow \dot{N}_{\max} = \eta \times \int \rho_b(x)\rho_y(x)dx$

$$\Rightarrow \frac{\int \rho_b(x)\rho_y(x)dx}{\int \rho_b(x)dx \int \rho_y(x)dx} = \frac{\dot{N}_{\max}}{\text{Area}} \quad !!!$$



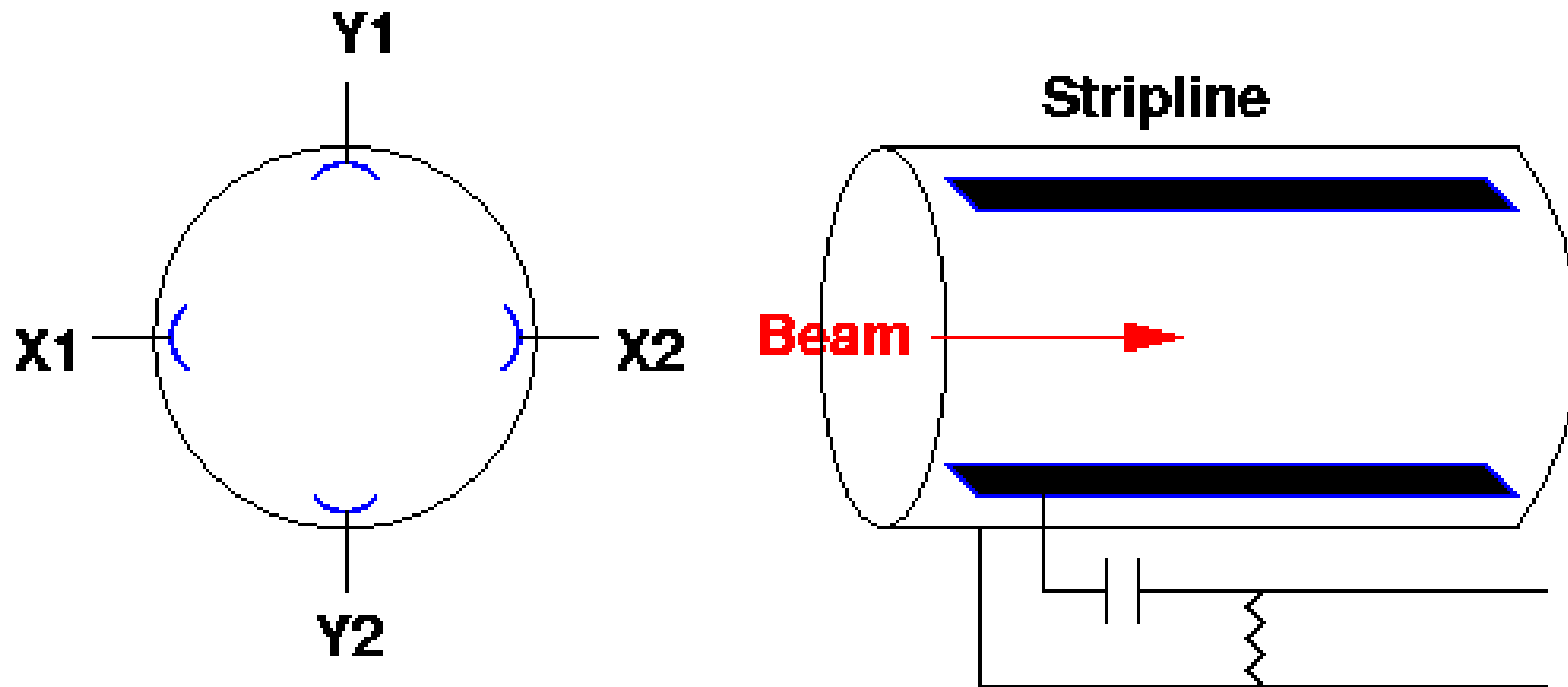
## How do we do a Vernier Scan?

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- Measure count rate, step beam by putting known bump in orbit, measure rate again, ...
- Rates measured with ZDC and BBC - multiple collisions not a significant problem yet
- Beam displacement known from models of beam transport through accelerator
- Displacements measured with Beam Profile Monitors (BPMs) at  $\pm 8$  meters from IP

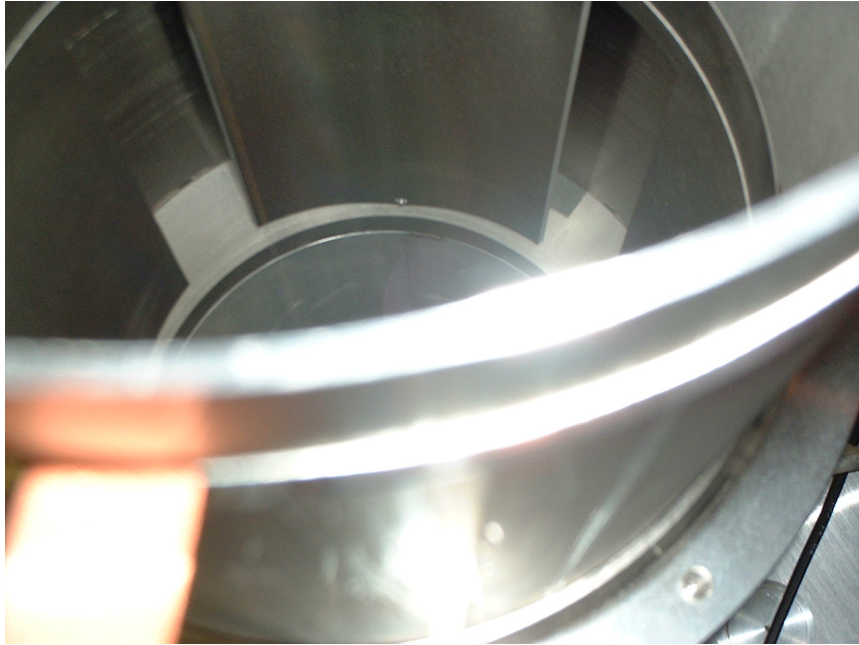


## How do the BPMs work?



- Beam Current induces voltage(t)  $\propto$  derivative of  $I_b(t)$
- BPM electronics find peak of first of two double humps
- Magnitude of voltage is sensitive to total charge and position w.r.t. stripline
- X Position  $\approx X2 - X1 / X2 + X1$
- Many possible offsets : electrical center  $\neq$  mechanical center  $\neq$  orbit center ...

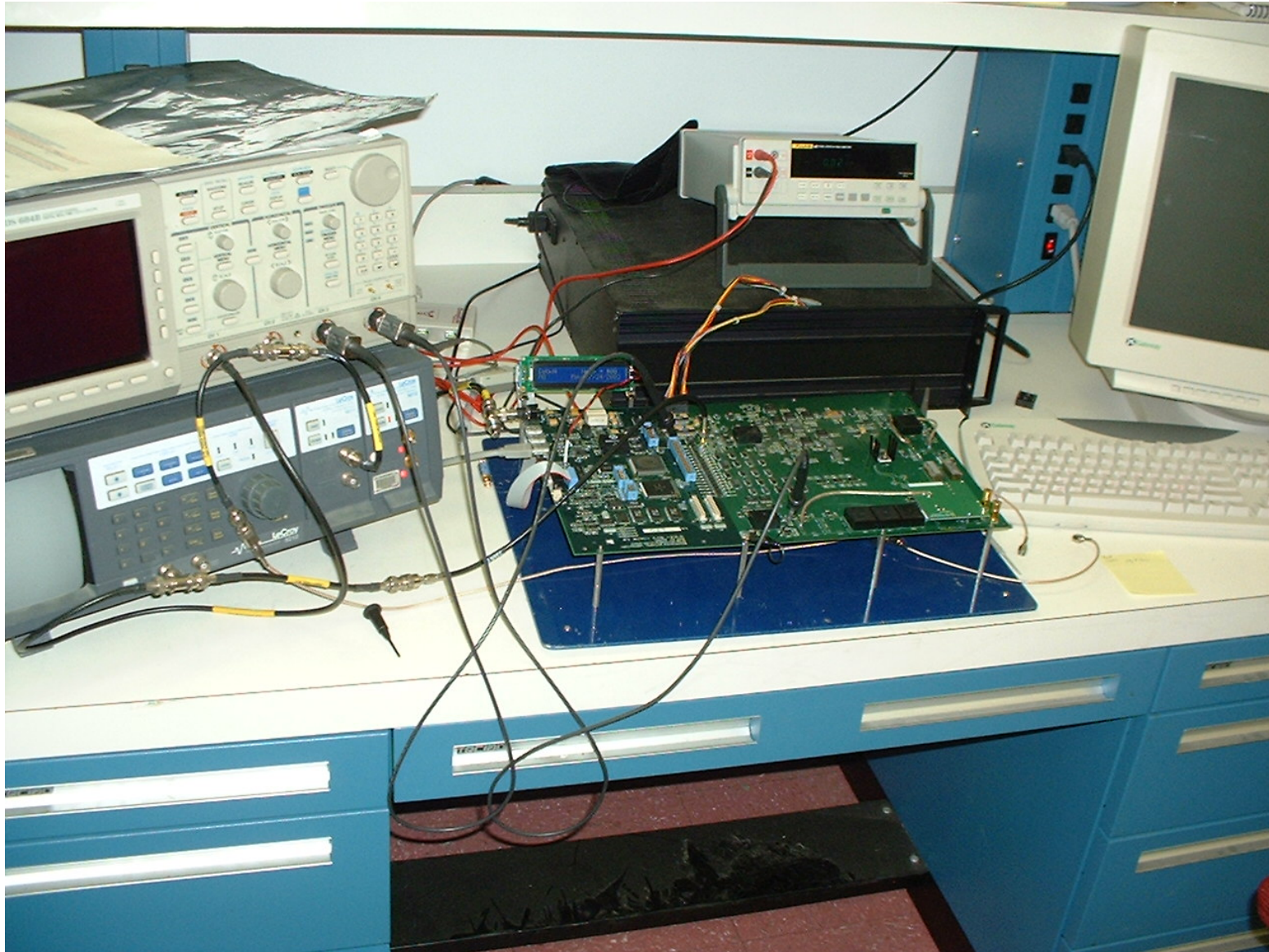
## Some Examples of Poor Photography



- More than 700 BPMs in RHIC and AGS
- Must make electrical connection from 4K to room temperature



# Calibration of BPM electronics



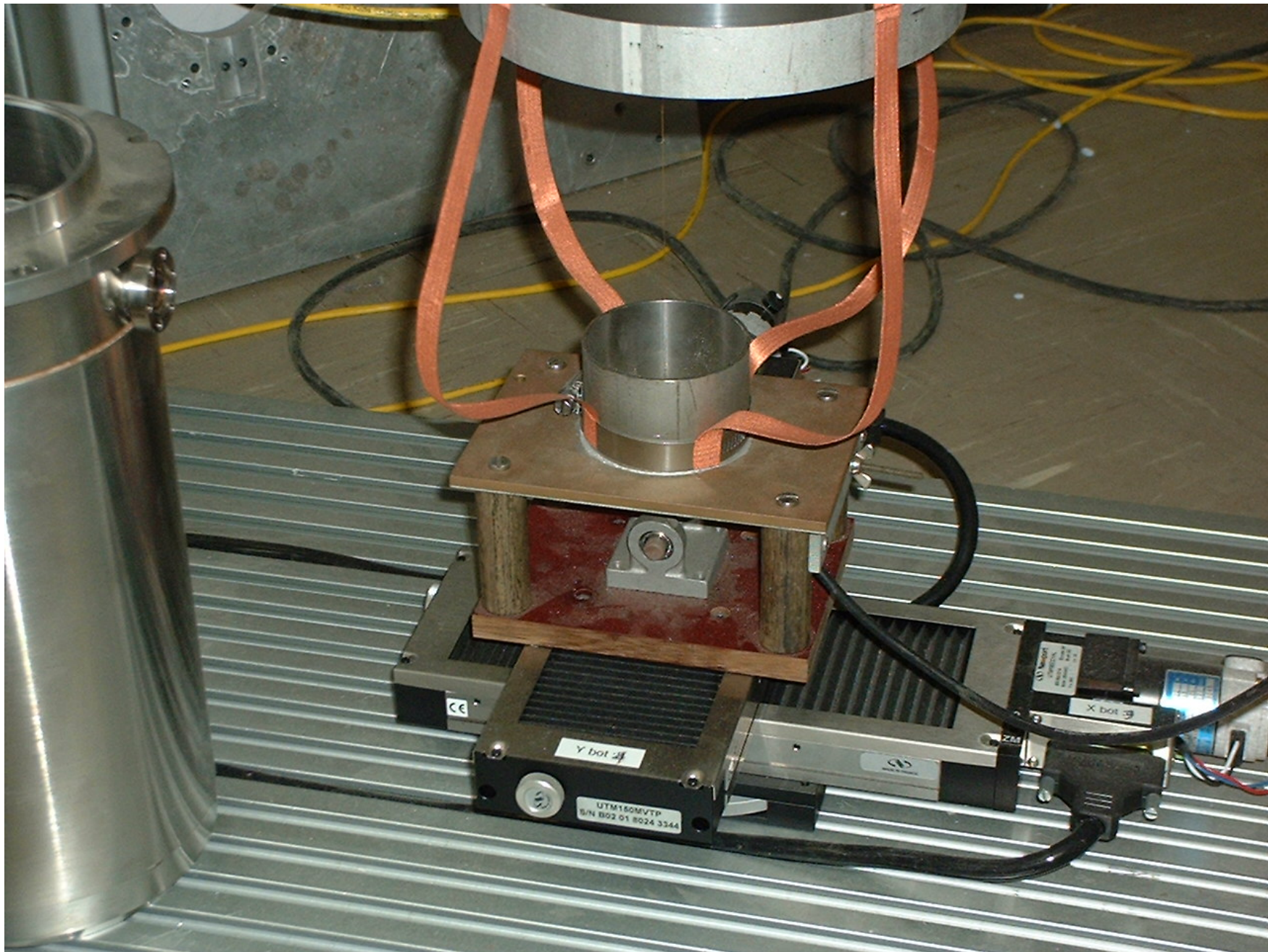
# BPM wire scanning calibration apparatus

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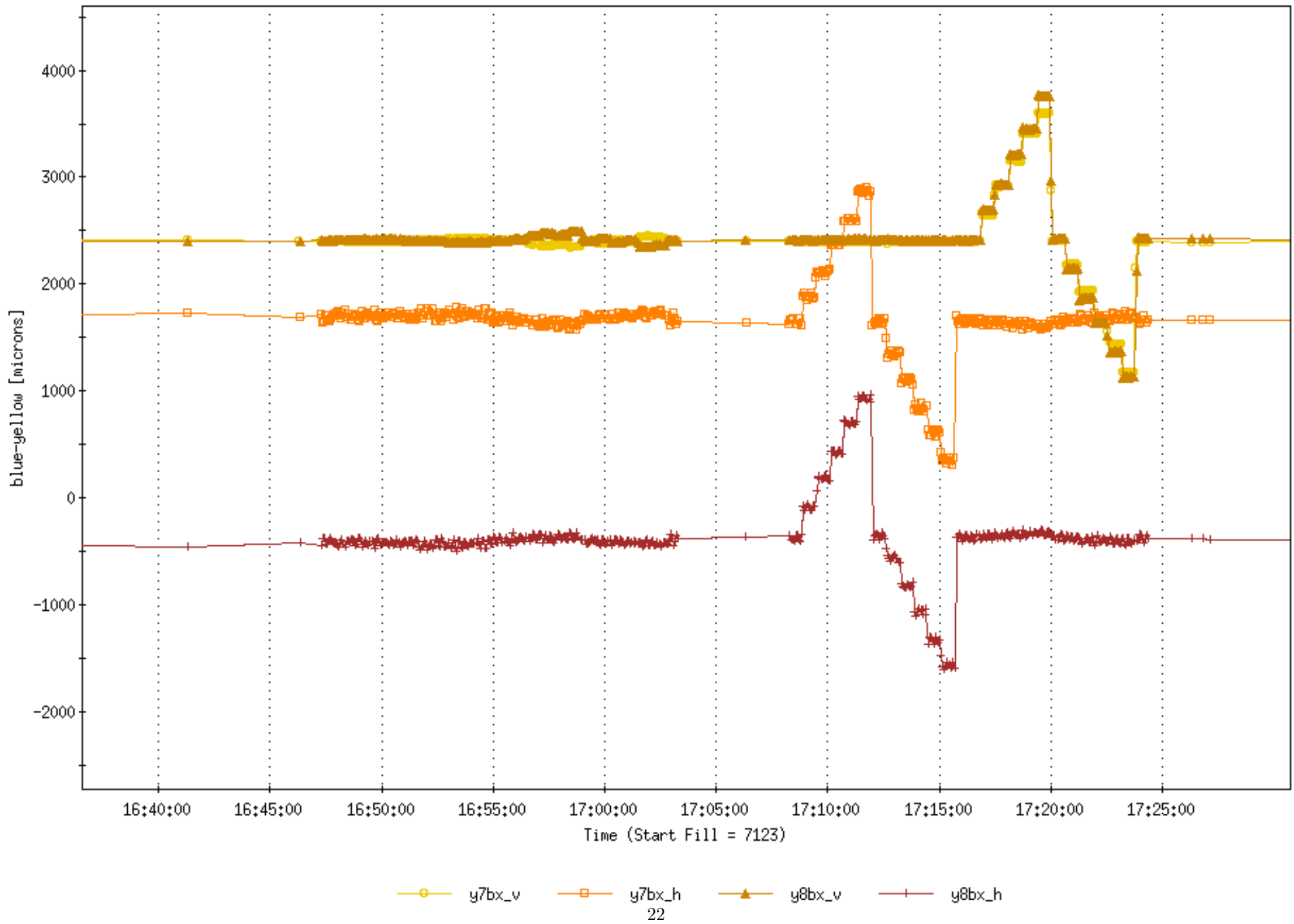




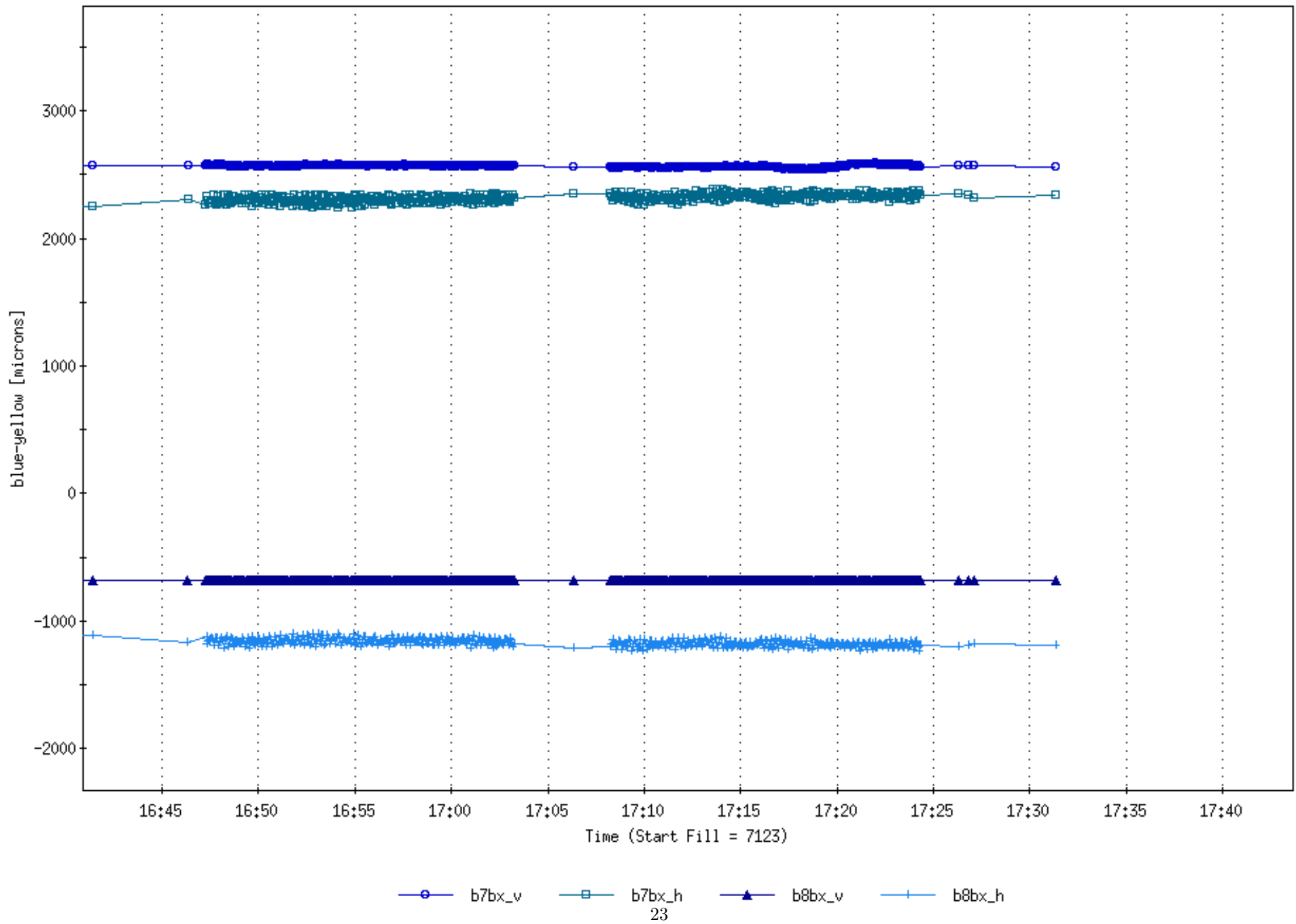
# BPM wire scanning calibration apparatus closeup



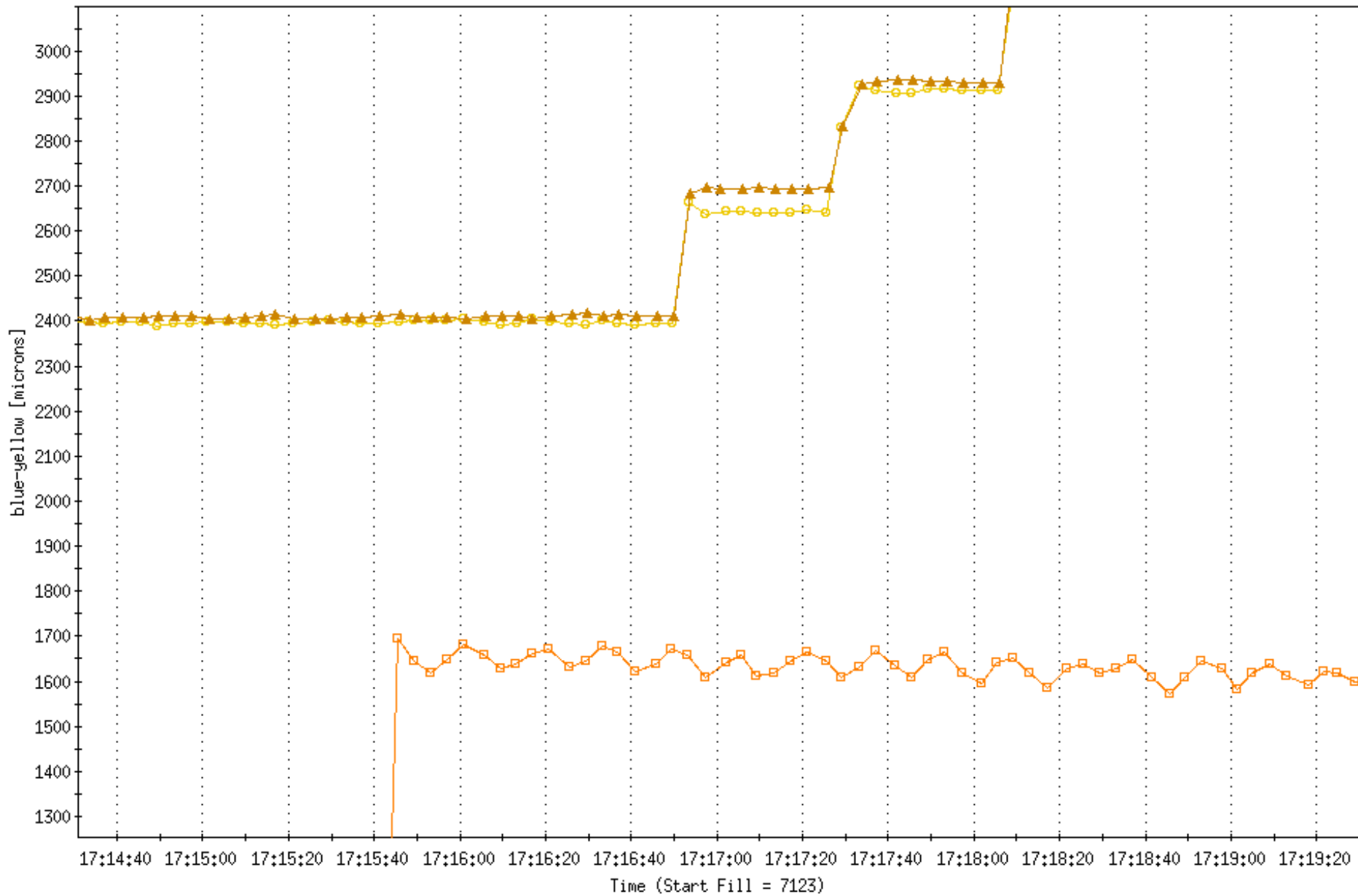
# Vernier Scan 174762 showing BPMs of Yellow beam



# Vernier Scan 174762 showing BPMs of Blue beam : Small coupling



# BPMs suggest beam is oscillating in horizontal



y7bx\_v y7bx\_h y8bx\_v y8bx\_h

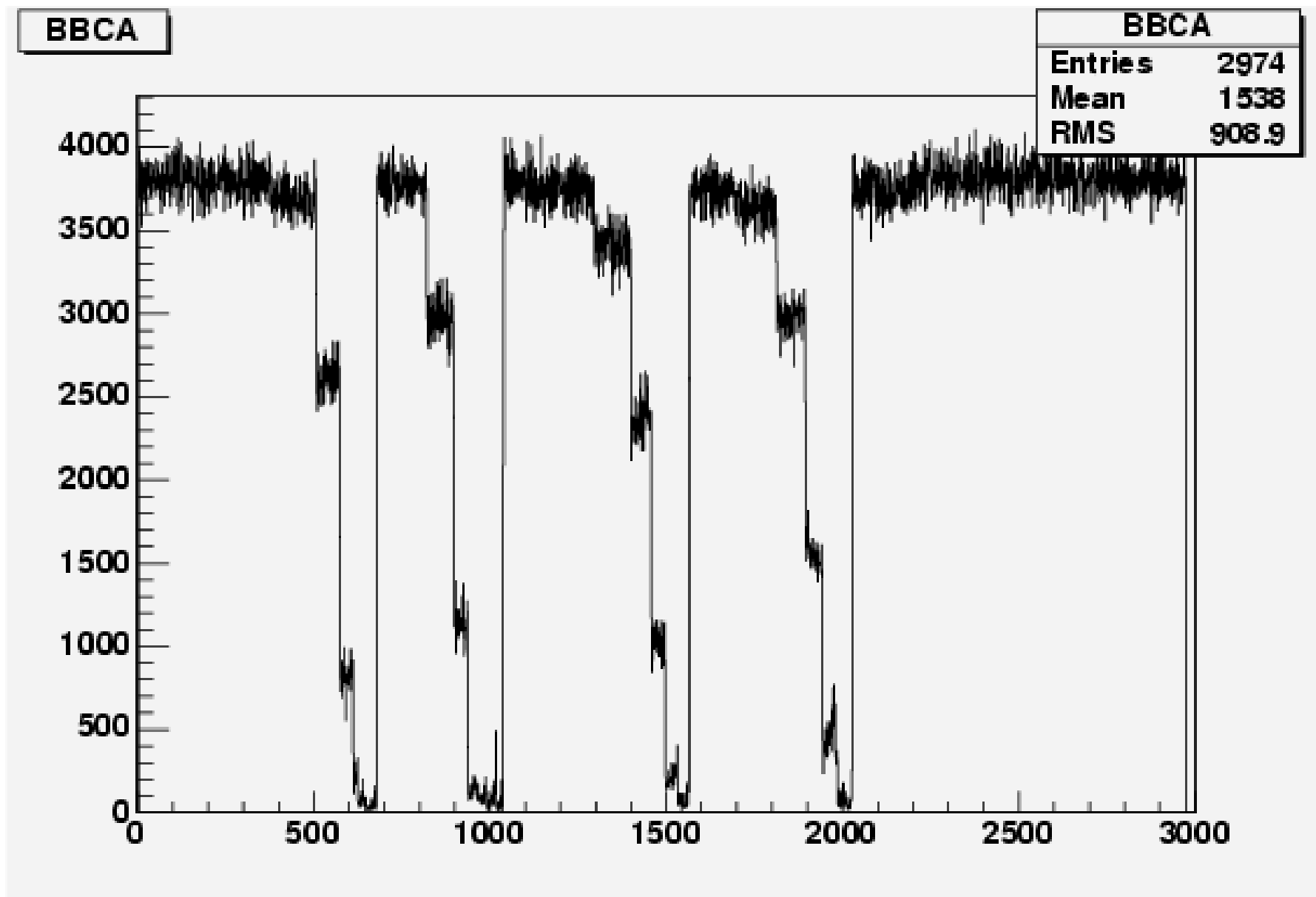


## Event Rates during Vernier Scan recorded in GL1P Scalers

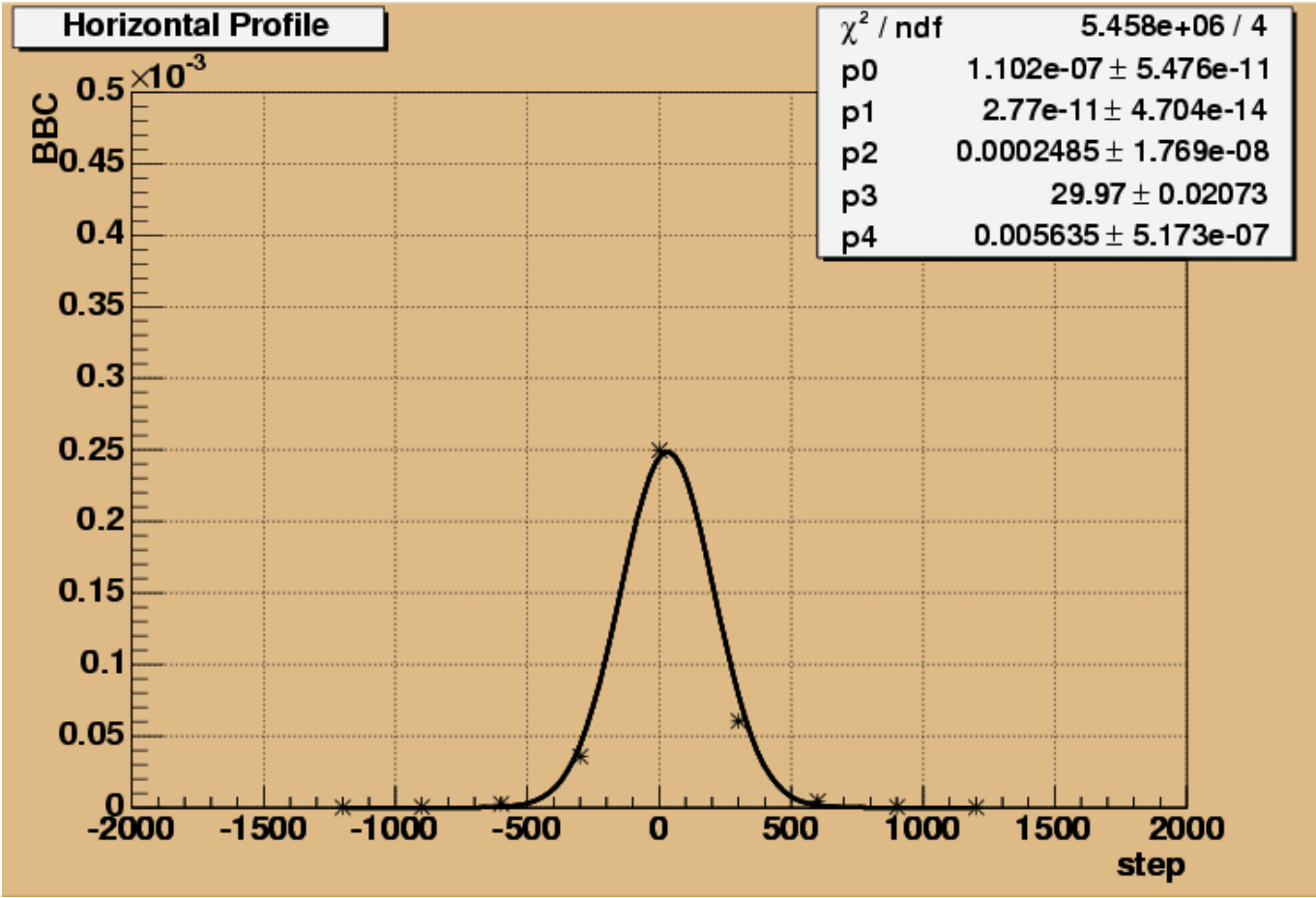
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- Put BBCLL1, Clock, ZDCWide, ZDCNS in GL1P scalers (live counts)
- Bin and plot BBCLL1/Clock, ZDCWide/Clock, and ZDCNS/Clock
- Find steps in rates corresponding to beam motions
- Plot Event Rate versus Beam displacement
- Find plot area, peak, combine with calibrated WCM data, insert into expression for  $\mathcal{L}$

# Sample Data (Thanks to Robert Bennett, Oleg Eysler)



# Sample Data (Thanks to Robert Bennett, Oleg Eysler)



## Many Sources of Uncertainty (see AN184)

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- Beams never overlap maximally - never measure  $\dot{N}_{\max}$
- Determination of Area requires beam profile assumptions or many measurement points
- Emittance blow-up during vernier scan, coupling of beam motions
- Non-zero crossing angle, Beam position and charge uncertainties
- Multiple collision effects, afterpulsing, backgrounds
- Rotated beam axis (should measure a few stripes)
- Factors affecting calibration of BBC as  $\mathcal{L}$  monitor : smearing effects, vertex position dependent efficiencies, “hourglass effect”, ...

## Accounting for Crossing Angle and Hourglass Effect (see Keil)

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$$\mathcal{L} \propto \int \rho_b(x, y, s) \rho_y(x, y, s) dx dy ds$$
$$\rho_b \propto \frac{N_b}{2\pi\sigma_0(1 + s^2/\beta_0^2)} \exp\left(-\frac{x^2 + y^2}{2\sigma_0^2(1 + s^2/\beta_0^2)}\right),$$
$$\rho_y \propto \frac{N_y}{2\pi\sigma_0(1 + s'^2/\beta_0^2)} \exp\left(-\frac{x'^2 + y'^2}{2\sigma_0^2(1 + s'^2/\beta_0^2)}\right)$$

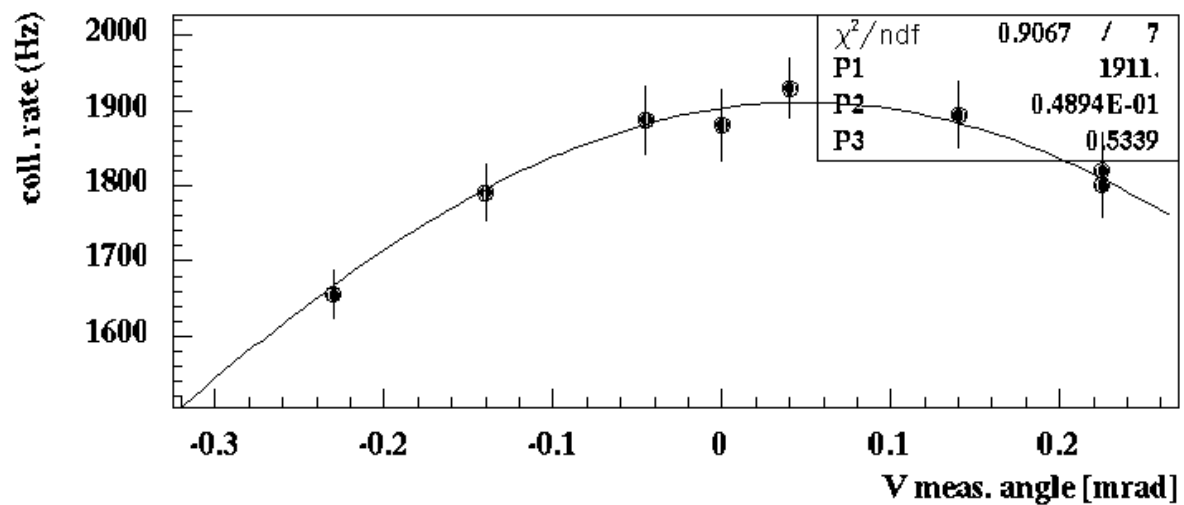
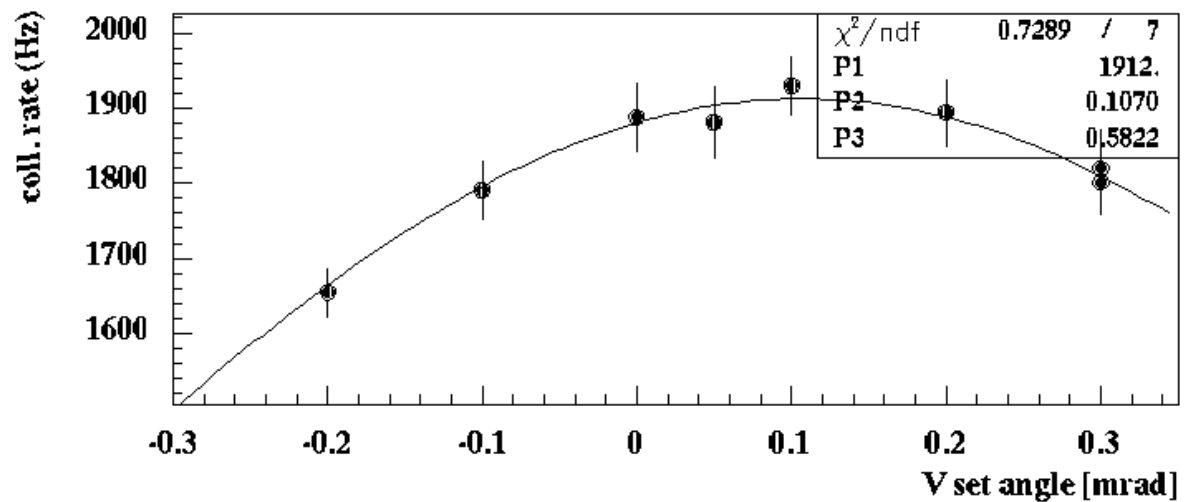
- $x', y'$  related to  $x, y$  by crossing angle
- Focusing account for by  $\sqrt{1 + s^2/\beta_0^2}$

$$\sigma_0 \approx \sqrt{\frac{\epsilon\beta_0}{6\pi\gamma}} \approx \sqrt{\frac{20\pi \times 10^{-6} \times 1}{6\pi 100}} \approx 200 \mu m$$

- Crossing angle effect has been measured (angle scan)
- Focussing effect leads to ZDCWide beam width appearing wider than ZDCNarrow (see John Lajoie's plots)

# Preliminary Results from Angle Scan (From Angelika)

## 7112 vertical angle scan



Thanks for your attention