

Design and Beam Test Results for the 2D Projective sPHENIX Electromagnetic Calorimeter Prototype

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Abstract—sPHENIX is a new experiment with the goal of studying the quark-gluon plasma and further understanding QCD matter and interactions. The sPHENIX detector is currently under construction at the Relativistic Heavy Ion Collider. A prototype of the sPHENIX electromagnetic calorimeter (EMCal) was tested at the Fermilab Test Beam Facility in Spring 2018 as experiment T-1044. The EMCal prototype corresponds to a solid angle of $\Delta\eta \times \Delta\phi = 0.2 \times 0.2$ centered at pseudorapidity $\eta = 1$. The prototype consists of scintillating fibers embedded in a mix of tungsten powder and epoxy. The fibers project back approximately to the center of the sPHENIX detector, giving 2D projectivity. The energy response of the EMCal prototype was studied as a function of position and input energy. The energy resolution of the EMCal prototype was obtained after applying a position dependent energy correction and a beam profile correction. Two separate position dependent corrections were considered. The EMCal energy resolution was found to be $\sigma(E)/\langle E \rangle = 3.5(0.1) \oplus 13.3(0.2)/\sqrt{E}$ based on the hodoscope position dependent correction, and $\sigma(E)/\langle E \rangle = 3.0(0.1) \oplus 15.4(0.3)/\sqrt{E}$ based on the cluster position dependent correction. These energy resolution results meet the requirements of the sPHENIX physics program.

Index Terms—Calorimeters, electromagnetic calorimetry, performance evaluation, prototypes, Relativistic Heavy Ion Collider (RHIC), silicon photomultiplier (SiPM), simulation, “Spaghetti” Calorimeter (SPACAL), sPHENIX

I. INTRODUCTION

sPHENIX is a new experiment [1] with the goal of elucidating QCD matter and interactions by studying the quark-gluon plasma (QGP) [2]–[6]. The sPHENIX detector is currently under construction at the Relativistic Heavy Ion Collider (RHIC). sPHENIX is designed to measure the QGP at a variety of length scales using various probes to provide insights into the microscopic properties of the QGP. One such probe is jets that arise from hard scattering interactions between two partons, with the energy loss of partons traversing the QGP being of particular interest. sPHENIX will allow for a detailed study of flavor dependent energy loss through a measurement of heavy flavor tagged jets, as well as open heavy flavor hadrons. Measurements of photon-tagged jets and jet substructure are also part of the sPHENIX physics program. sPHENIX will allow for measurements of jets with transverse momentum as low as 10 GeV, as well as provide measurements of both the hadronic and electromagnetic components of jets at RHIC. To

accomplish these measurements, sPHENIX is designed with a tracking system, a calorimeter system with 2π azimuthal acceptance and pseudorapidity coverage of $|\eta| < 1.1$, and the former BaBar solenoid magnet [7]. The calorimeter system consists of an electromagnetic calorimeter and a hadronic calorimeter. The use of the BaBar magnet imposed constraints on the sPHENIX detector design. In particular, the electromagnetic calorimeter was required to be compact enough to fit inside the magnet while allowing enough space for the tracking system and part of the hadronic calorimeter.

The sPHENIX electromagnetic calorimeter (EMCal) is a sampling calorimeter designed to measure the electrons, positrons, and photons in electromagnetic showers. The EMCal will also measure approximately one interaction length of hadronic showers. The EMCal has a coverage of $|\eta| < 1.1$ and full azimuth. The EMCal is segmented into *towers* of size $\Delta\eta \times \Delta\phi = 0.024 \times 0.024$, with an approximate volume of $2.5 \times 2.5 \times 14 \text{ cm}^3$, which sets the granularity of the calorimeter. The towers are defined within calorimeter *blocks* that consist of scintillating fibers embedded in a mix of tungsten powder and epoxy. Each block corresponds to a 2×2 array of towers. Each tower is equipped with a light guide coupled to silicon photomultipliers that collect the light from the fibers. The blocks are distributed in 64 sectors that describe an overall cylindrical geometry concentric with the beamline and centered at the interaction point of the particle collisions. Each side $0 < |\eta| < 1.1$ has 32 sectors distributed evenly in azimuth. Each sector has 24 rows of blocks extending along the beamline, and each row has 4 blocks along the ϕ direction. The blocks are tapered in both η and ϕ , resembling a truncated pyramid, and giving a 2D projective geometry. The blocks are further tilted such that the fibers do not project directly at the interaction point, minimizing channeling and improving energy resolution. More details about the sPHENIX detector and the EMCal can be found in reference [8].

The sPHENIX physics program requires an EMCal energy resolution equal or better than $16\%/\sqrt{E} \oplus 5\%$. This requirement is motivated by the measurement of the Upsilon states through the electronic decay channel $\Upsilon \rightarrow e^-e^+$. The electrons from these Upsilon decays are expected to produce EMCal electromagnetic showers with energies of approximately 4 to 10 GeV. In contrast, underlying event fluctuations in central Au+Au collisions would produce a comparable measurement of approximately 320 MeV [8]. The energy

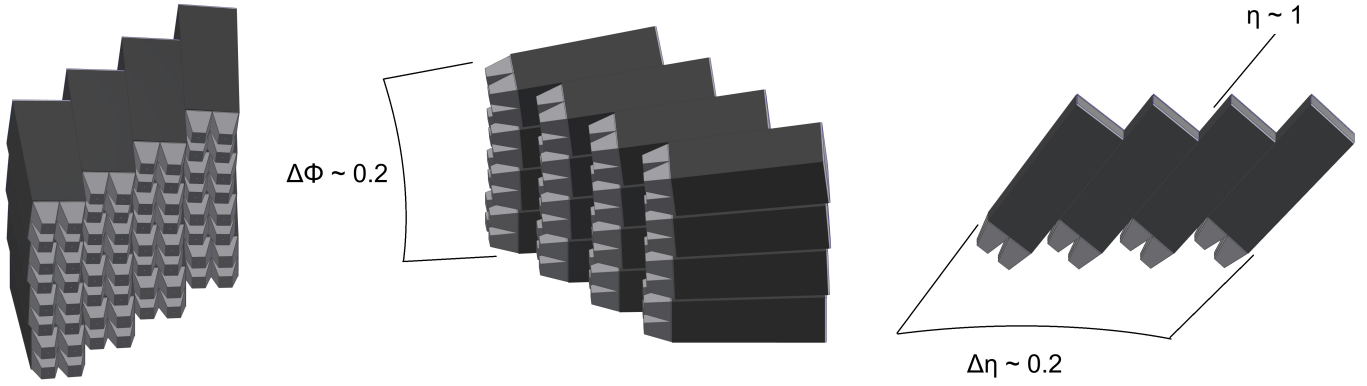


Fig. 1. EMCal prototype. The prototype consists of an array of 4×4 blocks, covering a solid angle of $\Delta\eta \times \Delta\phi = 0.2 \times 0.2$ centered at $\eta = 1$. Each block (dark gray) corresponds to a 2×2 array of towers defined by light guides (light gray).

89 resolution requirement was based on the maximum energy
90 smearing that would allow discrimination of the Upsilon states
91 against the average underlying event fluctuations.

92 A prototype of the EMCal was constructed in order to test
93 its energy resolution. The prototype corresponded to an array
94 of 8×8 calorimeter towers, or 4×4 blocks, centered at $\eta = 1$.
95 The prototype covered a solid angle of $\Delta\eta \times \Delta\phi = 0.2 \times 0.2$.
96 Figure 1 shows a schematic view of the EMCal prototype.

97 A previous prototype of the EMCal was tested in 2016
98 [9]. There are various differences between the 2016 prototype
99 and the 2018 prototype discussed in this paper. The most
100 notable difference is the projectivity of the EMCal blocks.
101 The 2016 prototype was only 1D projective (in ϕ), whereas
102 the 2018 prototype is 2D projective (in η and ϕ). The 2D
103 projectivity is a desirable feature because it improves energy
104 measurements at higher pseudorapidity. For a 2D projective
105 design, an electromagnetic shower at high pseudorapidity is
106 contained within a smaller number of towers than for a 1D
107 design, which results in a greater signal per tower and a better
108 discrimination against underlying event fluctuations. Another
109 difference between the prototypes is the pseudorapidity region
110 that they covered. While both prototypes corresponded to a
111 slice $\Delta\eta \times \Delta\phi = 0.2 \times 0.2$ of the EMCal, the 2016 prototype
112 was centered at $\eta = 0$ and the 2018 prototype was centered at
113 $\eta = 1$. The change in pseudorapidity was motivated by the fact
114 that the 2D projectivity reduces to 1D towards $\eta = 0$ because
115 the sPHENIX detector is symmetric with respect to this plane.
116 Other changes were also introduced in the 2018 prototype in
117 order to optimize the EMCal design (details in reference [8]),
118 but the 2D projectivity and the high pseudorapidity are the
119 main differences with respect to the previous prototype. The
120 final EMCal design that will be implemented in sPHENIX will
121 closely follow the design of the 2018 prototype.

122 II. PROTOTYPE ELECTROMAGNETIC CALORIMETER

123 A. EMCal Block Production

124 The EMCal blocks were produced by embedding a matrix
125 of scintillating fibers in a mix of epoxy and tungsten powder.
126 The blocks are similar to the ‘‘Spaghetti Calorimeter’’ design
127 used in other experiments [10]–[16]. The scintillating fibers are
128 as long as the block and are distributed uniformly across the

129 block’s cross section. There is a total of 2668 fibers per block.
130 The towers within a block have an area of approximately
131 $(1.1R_M)^2$, where $R_M \approx 2.3$ cm is the Molière radius. The
132 length of the towers varies with η and it has an approximate
133 value of $20X_0$, where $X_0 \approx 7$ mm is the radiation length. The
134 block density is approximately 9.5 g/cm³, with a sampling
135 fraction of approximately 2.1%.

TABLE I
EMCAL BLOCK MATERIALS

Material	Property	Value
Scintillating fiber	Saint Gobain BCF-12	
	diameter	0.47 mm
	core material	polystyrene
	cladding material	acrylic
	cladding	single
	emission peak	435 nm
	decay time	3.2 ns
	attenuation length	≥ 1.6 m
Tungsten powder	THP Technon 100 mesh	
	particle size	25-150 μ m
	bulk density (solid)	≥ 18.50 g/cm ³
	tap density (powder)	≥ 10.9 g/cm ³
	purity	$\geq 99\%$ W
	impurities ($\leq 1\%$)	Fe, Ni, O ₂ , Co, Cr, Cu, Mo
Epoxy	EPO-TEK 301	

136 The materials used to produce the blocks are listed in Table I
137 along with some of their properties. The blocks were produced
138 at the University of Illinois at Urbana-Champaign following
139 this procedure [8]:

- 140 • Scintillating fibers are dropped into mesh screens that
141 hold the fibers in place.
- 142 • The fiber-screen assembly is put into a mold.
- 143 • Tungsten powder is poured into the mold. The mold is
144 placed on a vibrating table to pack the powder.
- 145 • Epoxy is poured into the top of the filled mold, while a
146 vacuum pump is used at the bottom to extract the air as
147 well as pull the epoxy through the mold.
- 148 • The filled mold is left to dry until the mix is solid.

- The block is unmolded and machined to its final shape. A diamond tip is used to machine the readout ends of the block.

A finished EMCAL block can be seen in Figure 2. The quality assurance of the blocks included tests of density, light transmission and size. The blocks had a density ranging from 9.2 to 9.8 g/cm³. All the blocks had more than 99% fibers that successfully transmitted light. The size of the blocks deviated from the nominal dimensions by less than 0.5 mm.

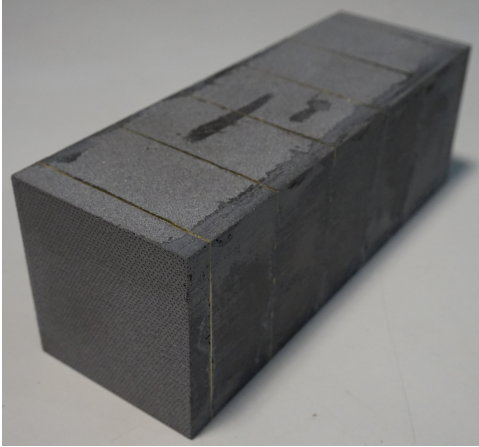


Fig. 2. EMCAL block. The block consists of scintillating fibers embedded in a mix of tungsten powder and epoxy. The blocks are tapered in two dimensions, giving a 2D projective geometry.

B. Light Collection

The light from the scintillating fibers was collected at the tower's front end (closer to the interaction point). Light guides were epoxied to the front of the blocks, while aluminum reflectors were epoxied to the back. The light guides consisted of UV transmitting acrylic with a trapezoidal shape (see Figure 3), custom made by NN, Inc. A silicone adhesive was used to couple each light guide to a 2×2 array of silicon photomultipliers (SiPM). Each SiPM (Hamamatsu S12572-015P) had an active area of 3×3 mm² containing 40K 15μm pixels, and had a photon detection efficiency of 25%. The signals from each of the four SiPMs were summed to give a single output signal from each tower. More details about the electronics are given in Section III. Figure 3 shows an EMCAL block equipped with light guides and SiPMs.

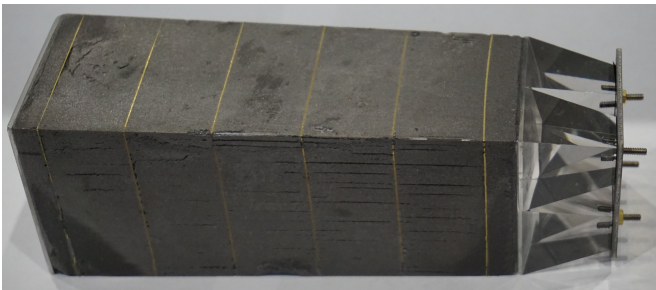


Fig. 3. EMCAL block equipped with light guides and SiPMs.

C. Assembly

Once the EMCAL blocks were equipped with light guides and SiPMs, they were stacked and epoxied together in their final positions. Since the SiPM gain is sensitive to temperature, a cooling system was used to remove the heat generated by the electronics. The cooling system consisted of multiple water coils connected to cold plates. The plates were coupled to the preamplifier boards that follow the SiPMs. Both the cooling system and electronics were controlled remotely. The EMCAL prototype can be seen in Figure 4, which shows the blocks, light guides, SiPMs, electronics and part of the cooling system.



Fig. 4. EMCAL prototype showing the EMCAL blocks, light guides, SiPMs, electronics and part of the cooling system.

III. READOUT ELECTRONICS AND DATA ACQUISITION

The summed signals from the four SiPMs from a tower were sent to a preamplifier, then shaped and driven into a digitizer. The SiPMs were operated at 4V above their breakdown voltage, which produces a gain of approximately 2.3×10^5 . A small thermistor was mounted at the center of the four SiPMs to monitor the temperature per tower. The temperature of the SiPMs was held constant within approximately 0.5°C. Since the gain temperature dependence of the SiPMs is approximately 1.5%/°C, temperature variations did not contribute significantly to the measured energy resolution. LEDs with an emission peak at 405 nm were mounted near the readout end of each tower and were used to provide a pulsed light source for calibration. Similarly, a charge injection test pulse was used to test and calibrate the readout electronics. The EMCAL prototype could operate in a nominal gain mode, or a high gain mode with 16 times the normal gain. The gain was selected through a slow control system.

The slow control system consisted of an interface board connected to a controller board. The interface board was mounted on the EMCAL prototype while the controller board was in a separate crate. The interface board contained digital-to-analog converters needed for different testing and monitoring tasks. The interface board controlled the SiPM bias and gain. Testing of the preamplifiers was controlled through the interface board

as well. The interface board also monitored leakage current and local temperature for compensation. The parameters for these testing and monitoring tasks were provided to the interface board by the controller board. An ethernet connection was used to communicate with the controller board.

Signals were digitized using a digitization system developed for PHENIX [17]. The signal waveforms were digitized using Analog-to-digital converters (ADC) at a sampling frequency of 60 MHz, followed by Field Programmable Gate Arrays. Signals were collected in Data Collection Modules and the data was finally recorded using the data acquisition system RCDAQ [18], [19]. The signals were recorded for the EMCal prototype as well as the external detectors mentioned in section IV.

IV. TEST BEAM

The EMCal prototype was tested at the Fermilab Test Beam Facility as experiment T-1044. The facility provided a particle beam, detectors such as a lead-glass calorimeter and Cherenkov counters, and a motion table in the MT6.2C area [20]. The EMCal was placed on the motion table to allow testing in different positions with respect to the beam.

The particle beam used in the experiment had energies ranging from 2 to 28 GeV and a profile size of a few centimeters, dependent on beam energy. The beam was composed mainly of electrons, muons, and pions, and their relative abundance depended on the energy [21], [22]. The beam hit the EMCal prototype with a frequency of 1 spill per min, where a spill corresponds to a maximum of approximately 10^5 particles during 4 seconds. The beam had a nominal momentum spread of $\delta p/p \approx 2\%$ for the energy range used [9], [10], [23]. A lead-glass calorimeter was used to measure the average and the spread of the beam momentum. The lead-glass calorimeter had a size of $45 \times 15 \times 15 \text{ cm}^3$ and an approximate resolution of $1.4\% \oplus 5.0\%/\sqrt{E}$ [9].

External detectors were used to discriminate electron signals from minimum ionizing particles (MIPs) and hadrons. Two gaseous Cherenkov counters were used for particle identification. The gas pressure in each Cherenkov counter was tuned to trigger only on electron signals. A hodoscope [10], [11] was placed upstream of the EMCal to determine the position of the particles in the beam precisely. The hodoscope consisted of 16 hodoscope fingers (0.5 cm wide scintillators) arranged in two arrays of 8 fingers each. One array had the hodoscope fingers arranged vertically and the other array had them arranged horizontally. The position of a hit in the hodoscope was given by a horizontal and a vertical hodoscope channel number. Each hodoscope finger was read out by an SiPM. Four veto detectors were also placed around the EMCal in order to suppress particles traveling outside the acceptance of the hodoscope. Each veto counter consisted of a scintillator coupled to a photomultiplier tube (PMT).

V. SIMULATIONS

The EMCal prototype was simulated using GEANT4 [24], [25] version 4.10.02-patch-02, with the physics list QGSP_BERT. The EMCal blocks were simulated following

their nominal design with a uniform block density. The simulations included an electron beam with a Gaussian profile. An 8 GeV beam with a standard deviation of 8 cm was used to study the prototype's energy response as a function of position. To study the prototype's energy response as a function of energy, the beam had an energy between 2 and 28 GeV and a standard deviation of 2.5 cm. For this energy dependent study, the beam was pointed between Towers A and B, which are located near the center of the prototype (see Figure 5). In the simulations, the energy deposits from the electromagnetic showers were converted into light using Birks' law [26] with constant $k_B = 0.0794 \text{ mm/MeV}$ [27]. The number of output photons was reduced by the light guide collection efficiency and then converted to number of fired SiPM pixels taking into account the SiPM saturation. The saturation was simulated by considering a Poisson distribution of photons randomly hitting the pixels and counting the total number of fired pixels. The mean of the Poisson distribution was proportional to the beam input energy, giving an energy dependent saturation effect. The number of fired pixels was converted to ADC counts and then calibrated to an input energy. The simulations were integrated into the sPHENIX analysis framework.

VI. ANALYSIS METHODS

A. Data Sets

The data sets used in this analysis correspond to a beam of electrons with energies of 2, 3, 4, 6, 8, 12, 16, 20, 24 and 28 GeV. The beam was pointed at either Tower A or Tower B (see Figure 5). In this paper, whenever Tower A or Tower B is mentioned, it is referring to the corresponding data set that had the beam centered at either of those towers.

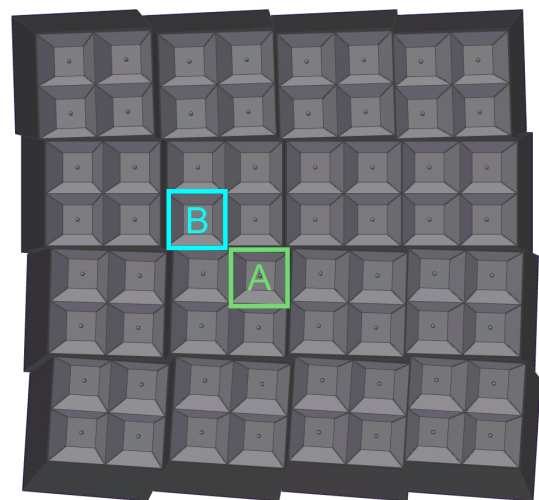


Fig. 5. Front view of the EMCal prototype showing the towers. Tower A (light green) and Tower B (light blue) are highlighted.

B. Electron Selection

Various cuts were used in order to suppress MIPs and hadrons, and select only events with single electrons. Single

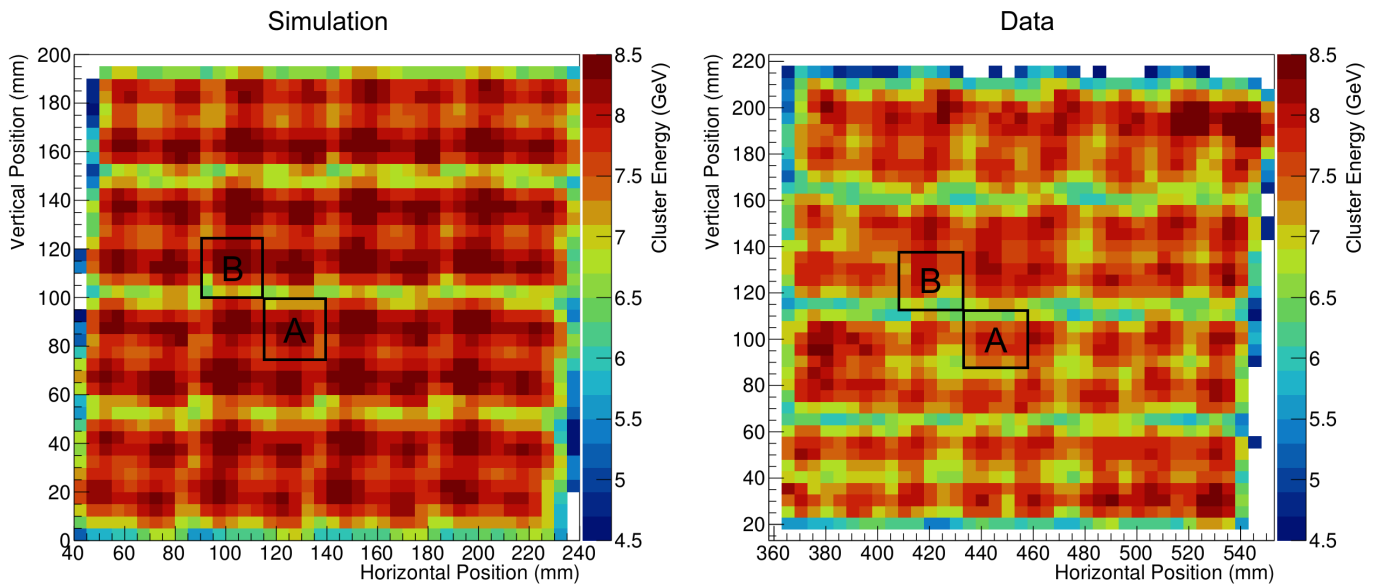


Fig. 6. Cluster Energy vs. Position for simulations (left panel) and data (right panel). The results correspond to an input energy of 8 GeV. Towers A and B are shown in black squares.

298 electrons were identified by requiring a Cherenkov cut, a
 299 vertical and horizontal hodoscope cut, and four veto cuts. It
 300 was generally assumed that the high energy peak in the energy
 301 spectra of the Cherenkov counters and hodoscope channels
 302 corresponded to the electrons. For the veto cuts, the high
 303 energy peak was assumed to correspond to particles traveling
 304 outside the beam position. The Cherenkov cut required the
 305 pulse height in the Cherenkov counters to be consistent with
 306 that of an electron. For the vertical and horizontal hodoscope
 307 cuts, the events were required to have an energy greater than
 308 50% of the peak energy in each hodoscope finger's energy
 309 spectrum. Only events with one hit in the vertical and one hit
 310 in the horizontal hodoscope fingers were considered. For the
 311 four veto cuts, the events were required to have an energy less
 312 than 20% of the peak energy in each veto detector's energy
 313 spectrum. These cuts gave a number of single electrons of
 314 approximately 5,000-50,000, depending on the energy.

315 C. Calibration

316 A preliminary calibration of the data, termed the *shower*
 317 *calibration*, was performed based on how the electromagnetic
 318 showers develop within the EMCal. A uniformity study of
 319 the EMCal prototype showed that the energy measurements
 320 depend on the transverse position within the EMCal. Figure
 321 6 shows the measured energy as a function of position for
 322 an input energy of 8 GeV, for both data and simulations. A
 323 higher energy collection efficiency is observed towards the
 324 center of the towers than at the boundaries between blocks
 325 and towers. This behavior motivated the use of secondary
 326 energy calibrations, the *position dependent correction* and the
 327 *beam profile correction*.

328
 329 The calibration procedures are as follows:

330 1) *Shower calibration*: For each event, the energy mea-
 331 sured by the EMCal was obtained as the total energy of a

5×5 cluster of towers around the maximum energy tower. 332
 The size of the cluster was selected based on the Molière 333
 radius for the EMCal blocks. A cluster of 5×5 towers 334
 contains over 95% of the electromagnetic shower energy. The energy 335
 corresponding to a cluster of 5×5 towers around the tower 336
 with the maximum energy is called the *cluster energy* and 337
 is denoted as E_{cluster} . The average cluster energy for an 8 338
 GeV electron beam *incident* at the center of each tower was 339
 reconstructed to the input energy and calibration constants 340
 were applied tower-by-tower. 341

2) *Position dependent correction*: The energy measured 342
 by the EMCal was corrected by a constant that depends on 343
 the position of the hit in the EMCal. Two different correc- 344
 tions were obtained, the difference lying in the availability 345
 of external position information. In the first, the position 346
 was determined by a horizontal and a vertical hodoscope 347
 finger, with a total of 8×8 possible positions. In the second, 348
 the position was determined by the energy averaged cluster 349
 position measured by the EMCal, discretized in 8×8 bins that 350
 matched the hodoscope. The position dependent calibration 351
 constants were obtained from 8 GeV data as described below. 352
 The procedure is the same for both the hodoscope-based and 353
 cluster-based corrections. For each of the 64 possible position 354
 bins, a histogram was filled with the cluster energy in that 355
 position. The histogram was then fit with a Gaussian of mean 356
 μ . The calibration constant for each position was obtained as 357
 8 GeV/ μ . The position dependent correction improved the 358
 energy resolution by 2-3%, depending on the energy. 359

The sPHENIX tracker can be used in place of a hodoscope 360
 to develop a position dependent correction. Since the tracker 361
 is only sensitive to charged particles, the cluster-based 362
 correction can be used for neutral particles instead. 363

364
 3) *Beam profile correction*: In the experiment, the beam 365
 had a different transverse profile at different energies. In 366

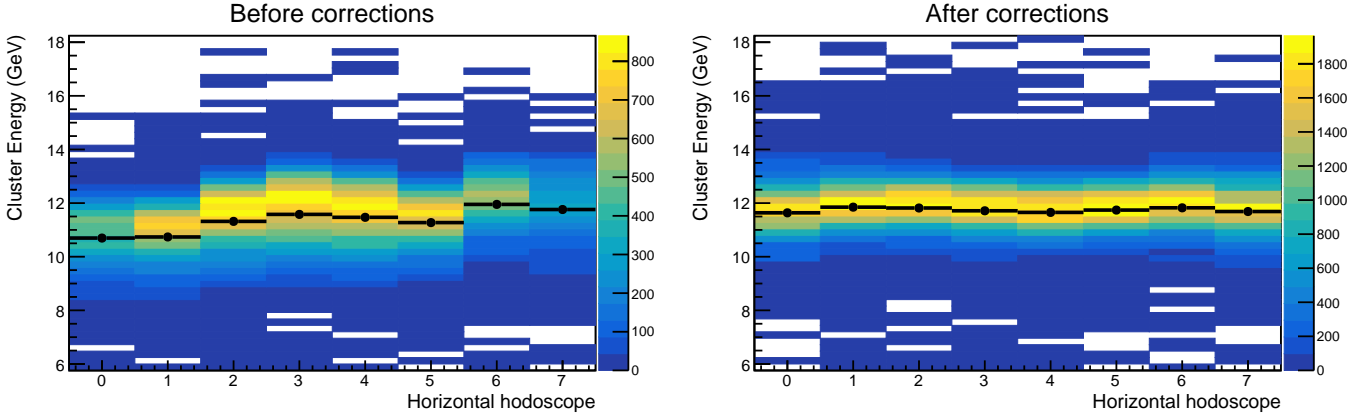


Fig. 7. Cluster Energy vs. Horizontal Hodoscope Position before (left panel) and after (right panel) applying the hodoscope-based position dependent correction and the beam profile correction. The color scale represents the number of events, while the black points correspond to the mean of the energy distributions for each hodoscope position. The data corresponds to a 12 GeV beam centered at Tower A.

367 addition to the position dependent correction, a *beam profile*
 368 *correction* was introduced in order to correct for the energy
 369 dependence of the beam profile. This correction consisted of
 370 filling the energy histograms with weights that were obtained
 371 by making the distribution of beam particles uniform as a
 372 function of position. The beam profile correction changed the
 373 energy resolution by 0.1-0.5%, depending on the energy.

374 The effects of these corrections on the energy response can
 375 be seen in Figure 7. This figure shows the cluster energy
 376 as a function of horizontal hodoscope position. The data is
 377 shown before and after applying the hodoscope-based position
 378 dependent correction and the beam profile correction. After
 379 the corrections are applied, the energy response of the EMCAL
 380 becomes more uniform.

381 The simulations also included the position dependent and
 382 beam profile corrections. The corrections were obtained using
 383 the procedure previously described, where the simulated
 384 position was discretized in 8×8 bins to mock the hodoscope.

385 VII. RESULTS AND DISCUSSION

386 Following the analysis procedure described in the previous
 387 section, the energy resolution and linearity of the EMCAL
 388 prototype was obtained for input energies ranging from 2 to
 389 28 GeV, for both simulations and data.

390 Figure 8 shows the energy resolution and linearity of the
 391 EMCAL prototype using a $2.5 \times 2.5 \text{ cm}^2$ cut centered at the
 392 tower. The $2.5 \times 2.5 \text{ cm}^2$ cut was selected based on the
 393 approximate area of a tower. The results are shown for data and
 394 simulations and include all corrections. The uncertainty bars
 395 on the data points correspond to the statistical uncertainties.
 396 The linearity was obtained as $E_{\text{cluster}} = E + cE^2$, where E
 397 is the input energy and c is a constant. The resolution was
 398 obtained as $\sigma(E_{\text{cluster}})/\langle E_{\text{cluster}} \rangle = \delta p/p \oplus a \oplus b/\sqrt{E}$, where
 399 a and b are constants, and a $\delta p/p = 2\%$ term was added to
 400 account for the beam momentum spread. Table II shows the
 401 values of the fit constants a , b , and c .

402 The resolution obtained with the cluster-based correction
 403 differs from the hodoscope-based correction by approximately
 404 0.6% in the constant term and 2.1% in the $1/\sqrt{E}$ term. Since

TABLE II
 EMCAL ENERGY LINEARITY AND RESOLUTION FOR A $2.5 \times 2.5 \text{ cm}^2$
 CUT CENTERED ON A TOWER

$$\text{Resolution fit: } \sigma(E_{\text{cluster}})/\langle E_{\text{cluster}} \rangle = 2\% \oplus a \oplus b/\sqrt{E}$$

$$\text{Linearity fit: } E_{\text{cluster}} = E + cE^2$$

	Tower	a (%)	b (% $\text{GeV}^{1/2}$)	c (GeV^{-1})
Data, hodoscope	A	3.2 ± 0.1	13.8 ± 0.2	$(-9.4 \pm 0.1) \times 10^{-4}$
Data, hodoscope	B	3.8 ± 0.1	12.8 ± 0.2	$(-10.9 \pm 0.1) \times 10^{-4}$
Data, cluster	A	2.7 ± 0.1	15.8 ± 0.3	$(-12.8 \pm 0.2) \times 10^{-4}$
Data, cluster	B	3.2 ± 0.1	14.9 ± 0.3	$(-8.6 \pm 0.3) \times 10^{-4}$
Simulation		3.04 ± 0.05	12.6 ± 0.1	$(-9.3 \pm 0.1) \times 10^{-4}$

405 the cluster-based correction depends on the position measured
 406 by the EMCAL itself and not the hodoscope, the difference
 407 in the results can potentially arise from the reduced cluster
 408 position resolution of the EMCAL at lower energy. Additionally,
 409 the energy resolution seems to be better in the simulations than
 410 in the hodoscope corrected data by approximately 0.5% in the
 411 constant term and 0.7% in the $1/\sqrt{E}$ term. These differences
 412 can arise from the lower energy collection efficiency at the
 413 boundaries between towers and blocks, as well as tower by
 414 tower variations that are not present in the simulations. The
 415 differences in the resolution results can be minimized by
 416 making a cut at the center of the towers, where the energy
 417 collection is most efficient. Figure 9 shows the linearity and
 418 resolution results using a $1.0 \times 0.5 \text{ cm}^2$ cut at the center of the
 419 towers. This figure shows better agreement between data and
 420 simulations. Table III shows the corresponding linearity and
 421 resolution fit constants.

422 Additionally, Figure 8 shows that for energies below 15
 423 GeV the resolution for Towers A and B generally agree
 424 within the statistical uncertainties, while for higher energies
 425 the resolution is consistently larger for Tower B than for
 426 Tower A. The disagreement between the resolution of the

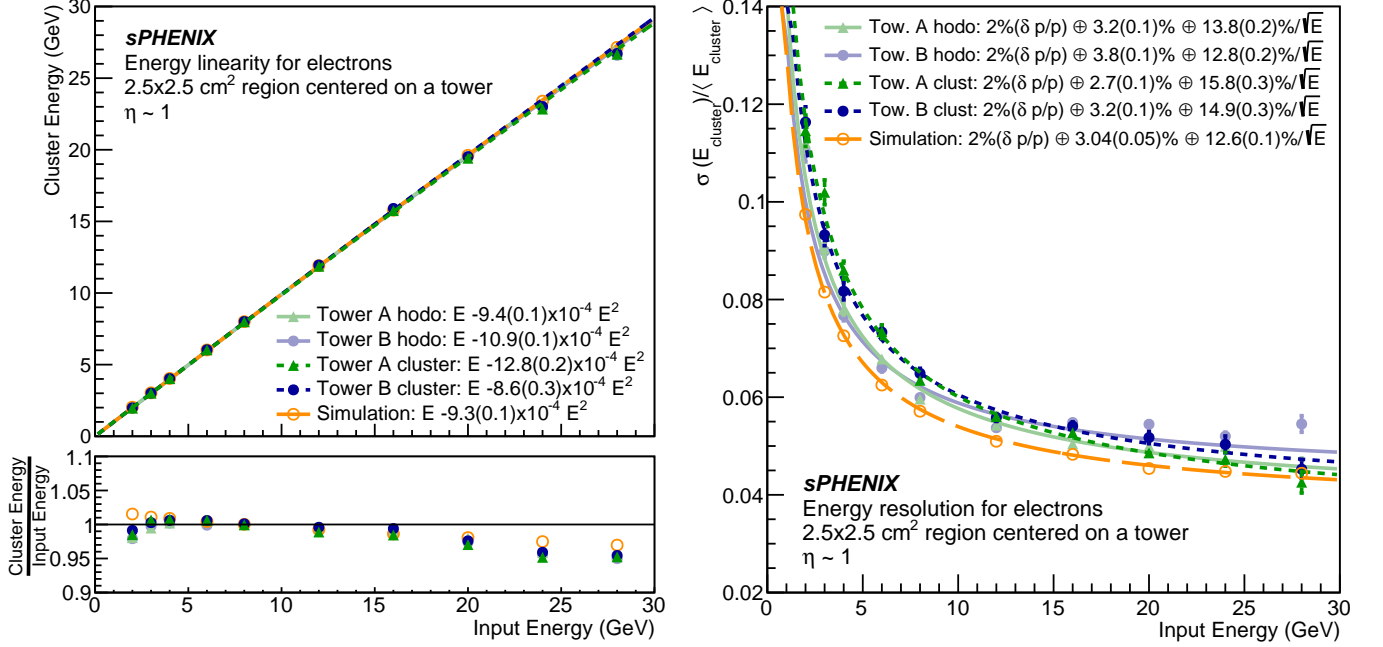


Fig. 8. Linearity and resolution of the EMCal prototype for a $2.5 \times 2.5 \text{ cm}^2$ centered on a tower. The $2.5 \times 2.5 \text{ cm}^2$ cut was selected based on the approximate area of a tower. The data corresponds to Tower A (green triangles) and Tower B (purple full circles). The data was corrected using the hodoscope-based (solid lines) and cluster-based (fine dashed lines) position dependent corrections, as well as the beam profile correction. Simulations (orange open circles, coarse dashed line) are shown for comparison and include the same corrections as the data. (top left panel) Cluster Energy vs. Input Energy. (bottom left panel) $\frac{\text{Cluster Energy}}{\text{Input Energy}}$ vs. Input Energy. The linearity was obtained as $E_{\text{cluster}} = E + cE^2$. (right panel) Energy Resolution vs. Input Energy. The resolution was obtained as $\sigma(E_{\text{cluster}})/\langle E_{\text{cluster}} \rangle = \delta p/p \oplus a \oplus b/\sqrt{E}$, where a $\delta p/p = 2\%$ term was added to account for the beam momentum spread.

427 towers above 15 GeV is observed for both the hodoscope-
 428 based and cluster-based results of Figure 8 and contributes to
 429 the fit constants of Table II. However, this disagreement is not
 430 observed when a cut at the center of the towers is used, as
 431 shown in Figure 9 and Table III.

TABLE III

EMCAL ENERGY LINEARITY AND RESOLUTION FOR A $1.0 \times 0.5 \text{ cm}^2$
 CUT AT THE CENTER OF A TOWER

Resolution fit: $\sigma(E_{\text{cluster}})/\langle E_{\text{cluster}} \rangle = 2\% \oplus a \oplus b/\sqrt{E}$

Linearity fit: $E_{\text{cluster}} = E + cE^2$

	Tower	a (%)	b ($\% \text{ GeV}^{1/2}$)	c (GeV^{-1})
Data, hodoscope	A	2.4 ± 0.2	12.3 ± 0.5	$(-12.9 \pm 0.3) \times 10^{-4}$
Data, hodoscope	B	2.3 ± 0.2	13.4 ± 0.5	$(+0.7 \pm 0.3) \times 10^{-4}$
Data, cluster	A	2.4 ± 0.2	13.2 ± 0.5	$(-10.9 \pm 0.3) \times 10^{-4}$
Data, cluster	B	2.7 ± 0.2	12.8 ± 0.4	$(-5.9 \pm 0.3) \times 10^{-4}$
Simulation		2.6 ± 0.2	11.9 ± 0.3	$(-9.1 \pm 0.3) \times 10^{-4}$

432 Comparing the 2018 results to the 2016 results of reference
 433 [9], the resolution improved for energies in the range 2 to 8
 434 GeV. In terms of the resolution fit, the $1/\sqrt{E}$ term of the reso-
 435 lution decreased by approximately 2.5% and the constant term
 436 increased by approximately 0.7%. Furthermore, the linearity

improved by approximately 1% in the 2018 prototype with
 respect to the 2016 prototype.

VIII. CONCLUSIONS

A 2D projective prototype of the sPHENIX EMCal was
 constructed and tested. The EMCal prototype's energy re-
 sponse to electrons was studied as a function of incident
 position and energy. The energy resolution and linearity of the
 EMCal prototype were obtained using two different position
 dependent energy corrections (hodoscope-based and cluster-
 based) as well as a beam profile correction. The two data
 sets used in this analysis had beam energies ranging from
 2 to 28 GeV, but one had the beam centered at Tower A
 and the other one had the beam centered at Tower B. The
 energy resolution was obtained for each tower using a cut
 of $2.5 \times 2.5 \text{ cm}^2$ centered on the tower. Based on the ho-
 doscope position dependent correction, the EMCal prototype
 was found to have a tower averaged energy resolution of
 $\sigma(E)/\langle E \rangle = 3.5(0.1) \oplus 13.3(0.2)/\sqrt{E}$. Based on the cluster
 position dependent correction, the tower averaged resolution
 was found to be $\sigma(E)/\langle E \rangle = 3.0(0.1) \oplus 15.4(0.3)/\sqrt{E}$.
 These energy resolution results meet the requirements of the
 sPHENIX physics program.

ACKNOWLEDGMENTS

This research was carried out using resources from the
 Fermi National Accelerator Laboratory (Fermilab), a HEP user

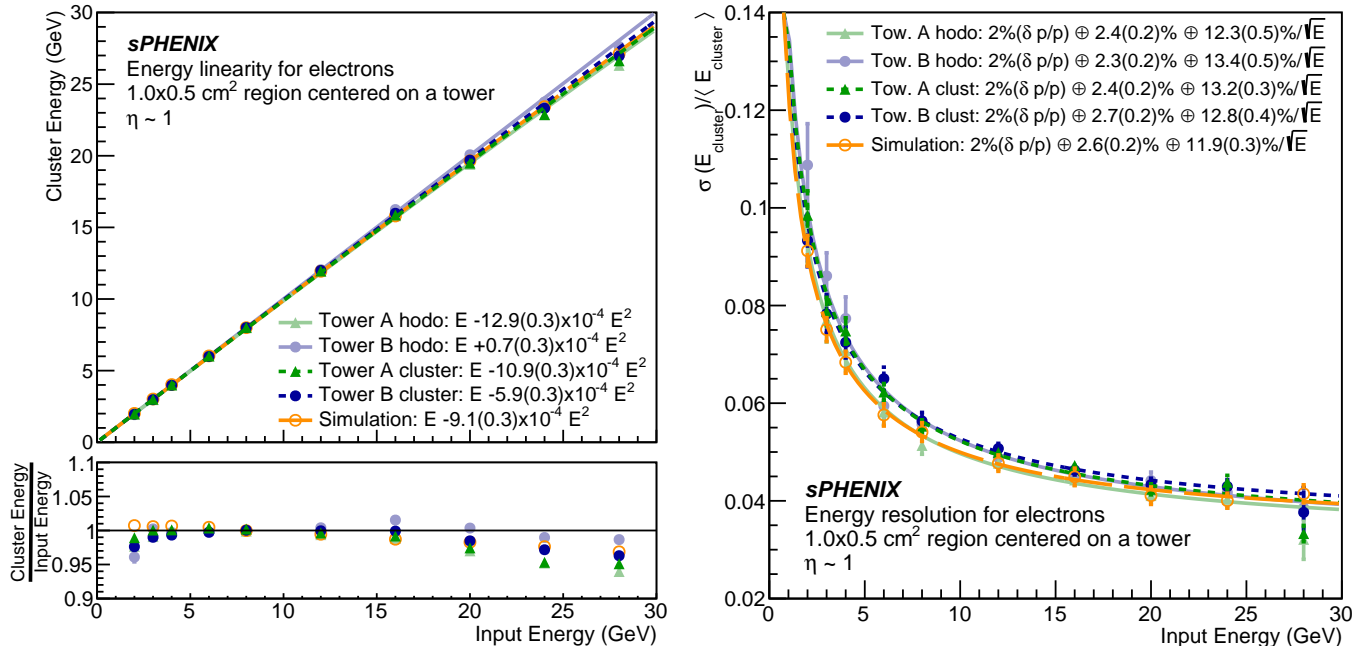


Fig. 9. Linearity and resolution of the EMCal prototype for a $1.0 \times 0.5 \text{ cm}^2$ cut at the center of a tower. The data corresponds to Tower A (green triangles) and Tower B (purple full circles). The data was corrected using the hodoscope-based (solid lines) and cluster-based (fine dashed lines) position dependent corrections, as well as the beam profile correction. Simulations (orange open circles, coarse dashed line) are shown for comparison and include the same corrections as the data. (top left panel) Cluster Energy vs. Input Energy. (bottom left panel) $\frac{\text{Cluster Energy}}{\text{Input Energy}}$ vs. Input Energy. The linearity was obtained as $E_{\text{cluster}} = E + cE^2$. (right panel) Energy Resolution vs. Input Energy. The resolution was obtained as $\sigma(E_{\text{cluster}})/\langle E_{\text{cluster}} \rangle = \delta p/p \oplus a \oplus b/\sqrt{E}$, where a $\delta p/p = 2\%$ term was added to account for the beam momentum spread.

462 facility managed by the Fermi Research Alliance, LLC, for the
 463 U.S. Department of Energy, Office of Science, acting under
 464 Contract No. DE-AC02-07CH11359. The authors would like
 465 to thank the technical staffs of the University of Illinois at
 466 Urbana-Champaign and Brookhaven National Laboratory for
 467 assisting the construction of the electromagnetic calorimeter
 468 prototype. The authors would also like to thank Dr. O. Tsai
 469 for providing the hodoscope used in the beam test.

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