

Heavy Flavor and Quarkonia Physics at sPHENIX

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The sPHENIX detector being constructed at BNL's Relativistic Heavy Ion Collider (RHIC) will begin measuring a plethora of Heavy Flavor and Quarkonia observables with unprecedented statistics and kinematic reach at RHIC energies starting in 2023. This includes the largest recorded sample of b-hadron decays from Heavy Ion collisions at RHIC, allowing for precise probes of the QGP using charm and beauty quarks. These measurements are enabled by the excellent vertexing of the MVTX detector, timing of the INTT, precision tracking by the TPC, and the EM and hadronic calorimetry systems, the latter of which is deployed for the first time at RHIC. The sPHENIX collaboration has created the reconstruction software stack as well as realistic data simulations, which allow for testing and optimization of the software and physics selections.

Keywords: *Open Heavy Flavor, Quarkonia, Heavy Ion Collisions, sPHENIX*

1 Introduction

The super Pioneering High Energy Nuclear Interaction eXperiment (sPHENIX) [1] is a relativistic heavy-ion experiment stationed at the Relativistic Heavy Ion Collider (RHIC) at Brookhaven National Laboratory (BNL). The sPHENIX barrel consists of several layers of detectors encompassing a full 2π azimuthal coverage with a pseudorapidity range of $\eta \leq 1.1$, and allows for a 15 kHz trigger rate with additional data taken in streaming mode. It also includes a 1.4 T superconducting solenoid magnet repurposed from the BaBar experiment. Based on current projections and progress, sPHENIX is poised to begin data taking in early 2023 with an initial Au+Au sample at $\sqrt{s_{NN}} = 200$ GeV, followed by p+Au and p+p samples at the same $\sqrt{s_{NN}}$ in 2024, and a larger Au+Au sample in 2025.

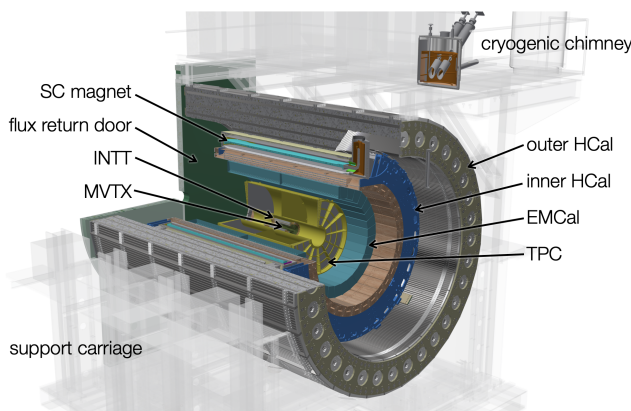


FIGURE 1. Rendering of the sPHENIX Detector with section cut-away. EPD and MBD not shown but would exist near the end caps of the barrel. [2]

The tracking system for the sPHENIX detector begins 2.3 cm from the center of the beam pipe with the MAPS-based micro-VerTeX detector (MVTX), followed radially

outward by the INTERmediate silicon strip Tracker (INTT), Time Projection Chamber (TPC), and TPC Outer Tracker (TPOT). Beyond the tracking detectors, the electromagnetic calorimeter (EMCal) and inner hadronic calorimeter (iHCal) systems continue outward while still within the BaBar magnet until the barrel is finally surrounded by the outer hadronic calorimeter (oHCal) outside the magnet. Additionally, sPHENIX employs Event Plane Detector (EPD) and Min Bias Detector (MBD) systems at the end caps of the barrel. A rendering of the complete detector can be seen in Figure 1.

sPHENIX aims to study in detail the quark gluon plasma (QGP) produced at RHIC as well as the nature of parton structure, parton energy loss, and mass dependent energy loss using its high precision next generation detectors. Exploring kinematic regions unavailable to and overlapping with the Large Hadron Collider (LHC), particularly at low p_T , additionally combine to establish a high level of importance for the experiment relative to the DOE/NSF NSAC 2015 Nuclear Physics Long Range Plan [3]. To accomplish its goals, the sPHENIX physics program is divided into four main components: jets and photons, open heavy flavor, upsilons, and cold QCD physics. The open heavy flavor and upilon physics components will be emphasized here.

2 Tracking Detectors and Performance

The sPHENIX tracking system is composed of three main detector systems: the MVTX, the INTT, and the TPC. Together these detectors are the keys to unlocking sPHENIX's heavy-flavor tagged jet and heavy-flavor physics programs.

Starting at the innermost of these three, the MVTX is based closely on the ALICE ITS inner barrel and performs precise vertexing using three layers of monolithic active pixel sensors (MAPS) ($\sim 5 \mu\text{m}$ space point precision per pixel for tracks with $p_T > 1$ GeV/c). The MAPS used in sPHENIX have a 5-10 μs integration time, which is a significant im-

64 improvement over the 180 μs integration time of the previous¹¹⁸
 65 generation of MAPS as used in the earlier runs of the STAR¹¹⁹
 66 experiment [4] [5]. The MVTX begins 2.3 cm radially out-¹²⁰
 67 ward from the center of the beam pipe and extends to 3.9 cm.¹²¹

68 Moving outward radially at 6 to 12 cm from the center of,¹²²
 69 the beam pipe resides the INTT, a silicon strip detector sur-¹²³
 70 rounding the MVTX made of two rings of two silicon strip,¹²⁴
 71 layers (one oriented in the ϕ direction, the other in the z direc-¹²⁵
 72 tion). Using space point resolution between that of the finer,¹²⁶
 73 MVTX and coarser TPC, the INTT allows for better pattern
 74 recognition and interpolation between the two. In addition,
 75 it is the only detector with single-beam-crossing timing res-
 76 olution, allowing us to associate a hit with a single bunch
 77 crossing. This will be crucial to sPHENIX's ability to asso-
 78 ciate fully reconstructed tracks with the event that produced
 79 them.

80 Lastly, the TPC sits between 20 to 78 cm radially out-
 81 ward from the center of the beam pipe. Filled with a 50:50
 82 mixture of Ne-CF₄ gas with an electron drift velocity of 8
 83 cm/ μs , it operates as the primary tracker for the experiment
 84 and provides precise momentum resolution. This resolution
 85 is critical for the experiment's goal of being able to distin-
 86 guish the $\Upsilon(2\text{S})$ and $\Upsilon(3\text{S})$ states with the $\sim 1.2\%$ momentum
 87 resolution needed in the 4-8 GeV/ c p_T range to accomplish
 88 this. Additionally, it is designed to be able to measure dis-
 89 placed tracks originating from a D or B meson decay, further
 90 enhancing the open heavy-flavor physics program. The TPC
 91 is currently set to operate in such a way that ion backflow is
 92 minimized which, combined with the shorter gas length com-
 93 pared to other larger TPCs such as at STAR, also results in re-
 94 duced particle identification (PID) capabilities using dE/dx .

95 With this tracking system, the software is designed to
 96 be able to calibrate and reconstruct Au+Au events within a
 97 few weeks of data taking. To be able to accomplish this,
 98 sPHENIX uses the fast and experiment-agnostic A Common
 99 Tracking Software (ACTS) from the ATLAS experiment [6].
 100 Track seeding for the silicon detectors is done separately
 101 from TPC track seeding with the former using ACTS, the lat-
 102 ter using cellular automation and a Kalman filter developed
 103 for sPHENIX based on the algorithm developed for the AL-
 104 ICE TPC, and the final combined track fitting using an ACTS
 105 Kalman filter that matches tracks using η , ϕ , and position at
 106 the beam line.¹²⁷

107 Recent simulation results project tracking efficiency in¹²⁹
 108 the 85-90% range above a p_T of 2 GeV after requiring 3¹³⁰
 109 MVTX hits for heavy flavor analysis using these methods.¹³¹
 110 Momentum resolution compatible with the necessary preci-¹³²
 111 sion needed for separating the first three Upsilon states is also¹³³
 112 observed, as well as p_T resolution below 3.5% at 20 GeV¹³⁴
 113 which is sufficient for making important jet physics measure-¹³⁵
 114 ments at high p_T . With continued improvements in track¹³⁶
 115 seeding, silicon track matching, and TPC clustering, these¹³⁷
 116 results are expected to continue to improve as the experiment¹³⁸
 117 moves towards data taking in early 2023.¹³⁹

3 Quarkonia Physics

Clear separation of the $\Upsilon(1\text{S}, 2\text{S}, 3\text{S})$ states is a key deliver-
 able of the sPHENIX experiment. Determining the centrality
 and p_T dependence of the nuclear modification factors for
 each of these states are critical measurements for the compar-
 isons sPHENIX intends to make between RHIC and the
 LHC due to the important differences shown in temperature
 profiles of the states from hydrodynamic calculations at dif-
 ferent collision energies.

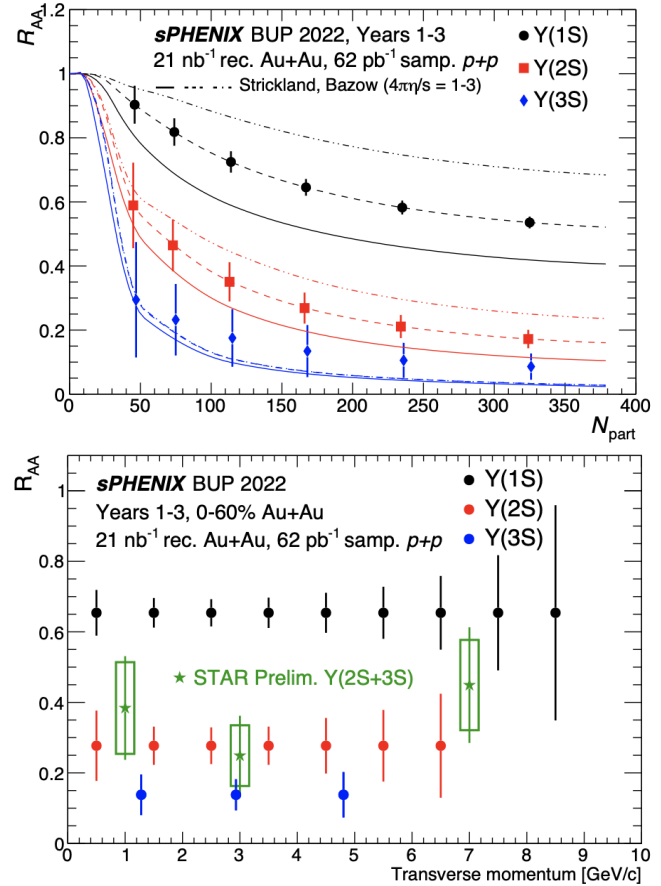


FIGURE 2. Projected statistical uncertainties for the centrality (top) and p_T (bottom) dependent R_{AA} measurements of the $\Upsilon(1\text{S}, 2\text{S}, 3\text{S})$ states at sPHENIX [2] [7] [8].

Figure 2 shows the current projected statistical uncertainties for these measurements using sPHENIX simulation data. The 1S and 2S R_{AA} projected measurements in the top figure are taken from predictions from Strickland and Bazow, while the 3S measurement is taken to be approximately half that of the 2S measurement as seen in the recent CMS measurement at the LHC [7]. In the bottom figure, the current best knowledge of Upsilon suppression from STAR is also included with the projected uncertainties for R_{AA} measurements as a function of transverse momentum in 0-60% Au+Au collisions.

It is also important to note sPHENIX's unique projected ability at RHIC to separate the first three Upsilon states. This has never been done before at RHIC, and has only recently

140 been performed by CMS at the LHC. It was previously ex-157
141 pected that the Drell-Yan background would be comparabl-158
142 with the 3S signal in Au+Au collisions at RHIC energies, but-159
143 the new CMS measurement suggests a heightened possibility-160
144 for sPHENIX to explore the systematics of the Upsilon 3S-161
145 state.

through D meson and b -jet R_{AA} measurements. By combin-
ing the reconstructed meson states and full jet, sPHENIX
is capable of covering a broad kinematic range for probing
the b -quark coupling with the QGP and providing stringent
constraints on model parameters.

146 4 Open Heavy Flavor Physics

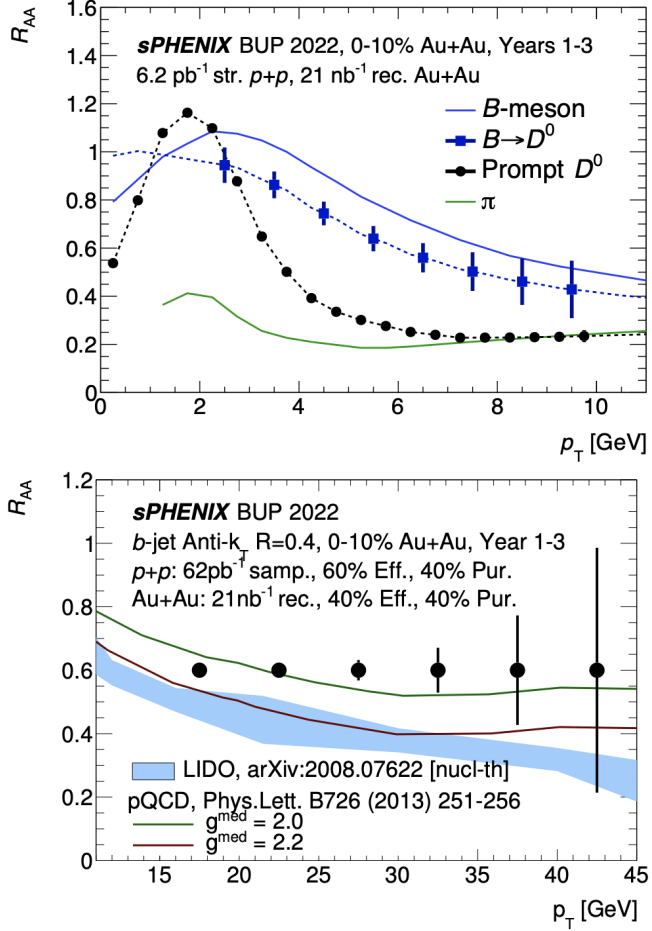


FIGURE 3. Projected statistical uncertainties for prompt and non-162
prompt D^0 meson (top) and b -jet (bottom) nuclear modification-163
factor measurements as a function of p_T in 0-10% Au+Au collis-164
ions [2]. 165 166

147 Because of their large masses, heavy-flavor quarks (c and b)-167
148 are predominantly created in initial hard scattering at high Q^2 -168
149 which allows for important insight into the initial stages of-169
150 heavy ion collisions and comparison with calculable produc-170
151 tion rates in perturbative QCD. The high precision and high-171
152 data rate of sPHENIX allows for improvements in measure-172
153 ments of these probes that are often limited by the rarity of-173
154 heavy-flavor signals in heavy ion collisions at RHIC energies.-174

155 Figure 3 shows the projected statistical uncertainties-175
156 for sPHENIX's observation of mass-dependent energy loss-176

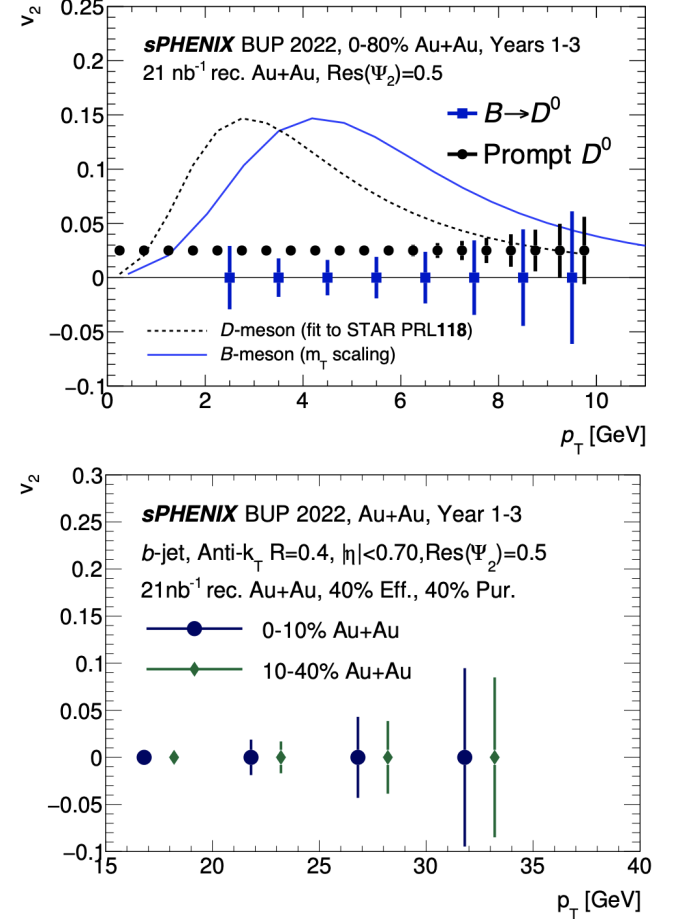


FIGURE 4. Projected statistical uncertainties for prompt and non-
prompt D^0 meson (top) and b -jet (bottom) elliptic flow measure-
ments as a function of p_T in Au+Au collisions at $\sqrt{s} = 200$ GeV
[2].

In addition to these R_{AA} measurements, sPHENIX plans
to make precise elliptic flow measurements for charm and
bottom mesons as well as b -jets in order to obtain insight into
how the heavy-flavor quark couples to the QGP medium. The
current projected statistical uncertainties for these measure-
ments can be seen in Figure 4. Due to the better constraints
on bottom quarks, as opposed to lighter charm quarks, in the
theoretical modeling of this measurement [9] [10], sPHENIX
plans to significantly constrain the heavy quark diffusion
transport parameter of the QGP as well as its temperature
dependence through making very precise bottom measure-
ments over a wide p_T range. At the higher end of the p_T
spectrum above 10 GeV, the bottom plot in Figure 4 demon-
strates sPHENIX's ability to examine the path-length depend-
ent energy loss of the b -quark as well.

177 sPHENIX also aims to investigate charm hadronization
 178 in the QGP through production ratio measurements. Recent
 179 results at both RHIC and the LHC suggest an enhancement
 180 of the Λ_c baryon compared to the D^0 meson, but RHIC is
 181 missing a baseline p+p measurement of that ratio as well as
 182 a broad in p_T and precise Au+Au measurement. The high
 183 statistics and precision of sPHENIX will allow for these mea-
 184 surements to be taken at RHIC so that questions remain-
 185 ing regarding differences in model predictions for charm
 186 hadronization in the QGP at RHIC energies can finally be
 187 resolved. Projected uncertainties for these measurements can
 188 be seen in Figure 5. It is important here to note as well that
 189 sPHENIX does not have π , K , p particle identification (PID)
 190 due to a lack of a time of flight detector and poor dE/dx reso-
 191 lution in the TPC, but the Λ_c baryon can still be reconstructed
 192 well enough to perform this analysis using geometric and
 193 momentum-based measurements.

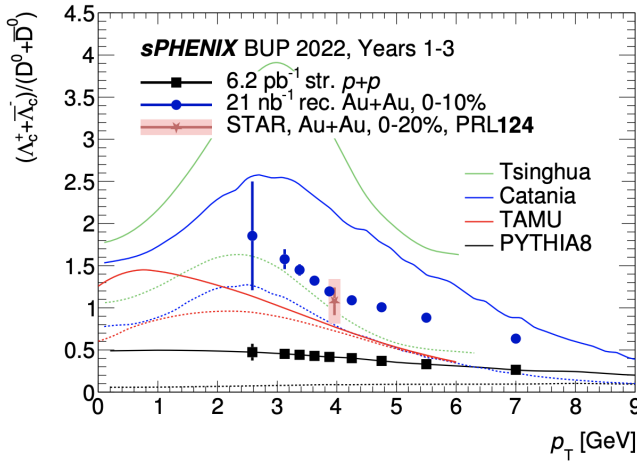


FIGURE 5. Projected Λ_c/D meson production ratio uncertainties in central Au+Au and p+p collisions at sPHENIX overlaid with STAR's recent comparable measurement [2].

194 4.1 HF Reconstruction with KFPARTICLE

195 As mentioned in the previous section, the sPHENIX compact
 196 TPC is designed for minimizing the ion backflow in the
 197 streaming mode, which limits its capability for providing
 198 dE/dx -based hadron PID. Without these PID capabilities,
 199 sPHENIX needs a way to perform HF reconstructions; the
 200 KFPARTICLE package is the answer. This tool is a current
 201 industry standard originally developed for the CBM exper-
 202 iment that is also used in other comparable experiments
 203 such as ALICE and STAR [11]. Using the pre-established
 204 external package, sPHENIX has been able to add a user
 205 interface, the ability to unpack tracks and vertices in the
 206 KFPARTICLE format, and the selection requirements and
 207 combinatorics needed to run the package effectively.

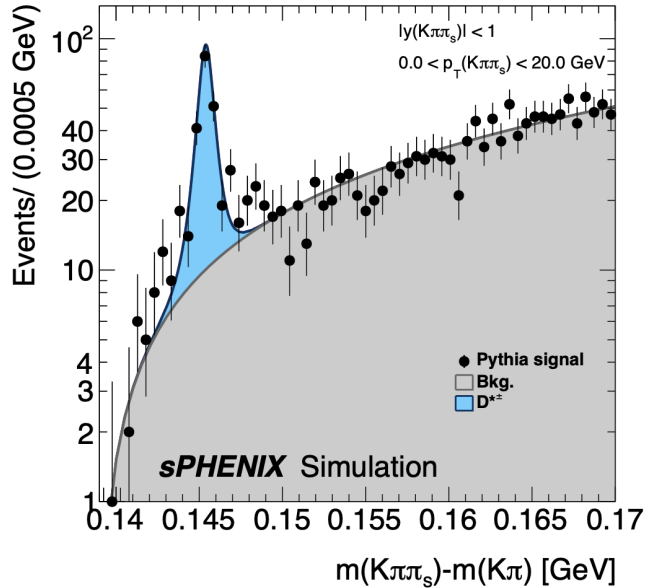
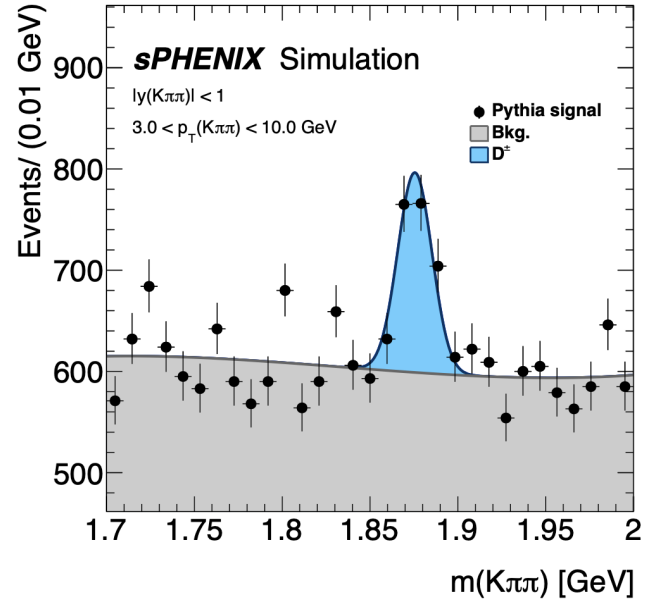
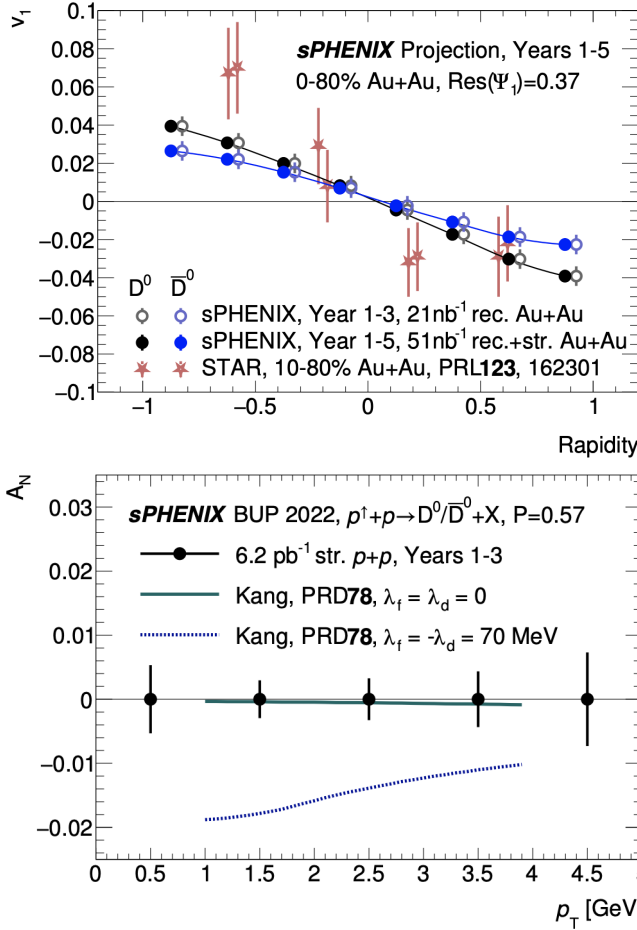


FIGURE 6. Mock Data Challenge 1 reconstruction of the D^\pm invariant mass (top) and $D^{*\pm}$ invariant mass (bottom) using the KFPARTICLE package implemented in the sPHENIX framework. Reconstruction performed on 50 million event $c\bar{c}$ sample from $\sqrt{s} = 200$ GeV p+p collisions without pileup [12].

In order to test the efficacy of this package in sPHENIX's software framework, a mock data challenge (MDC) was undertaken by the collaboration to evaluate the performance of its current simulation and reconstruction software. Significantly for this discussion, a reconstruction of the D^\pm and $D^{*\pm}$ meson was performed using sPHENIX reconstruction software and the sPHENIX implementation of the KFPARTICLE package. As can be seen in Figure 6, clear invariant mass peaks for each of these mesons were able to be reconstructed from a 50 million event $c\bar{c}$ sample from p+p collisions at $\sqrt{s} = 200$ GeV without pileup. These figures along with

219 others that were evaluated during the MDC give the group 238
 220 confidence in the efficacy of the KFParticle package within 239
 221 the sPHENIX framework, as well as putting the group on the 240
 222 right track to be ready for data-taking on Day 1.

223 4.2 D^0/\bar{D}^0 Separation



265 FIGURE 7. Projections of D^0/\bar{D}^0 uncertainties for v_1 (top) and 266
 267 transverse single spin asymmetry (bottom) measurements with 268
 269 sPHENIX [2].

224 The KFParticle package will additionally be critical in allow- 269
 225 ing for the separation of the D^0 and \bar{D}^0 mesons without ef- 270
 226 fective PID. The better sPHENIX is able to separate these 271
 227 heavy flavor mesons, the better the measurements shown in 272
 228 Figure 7 can be taken and refined to determine properties 273
 229 of the QGP and cold QCD. In the top figure, separate mea- 274
 230 surements of D^0 and \bar{D}^0 v_1 are projected to show a split- 275
 231 ting across the sPHENIX rapidity range due to a predicted 276
 232 transient magnetic field effect that is odd under charge con- 277
 233 jugation. Compared to the precision of the STAR measure- 278
 234 ment shown in the same figure, sPHENIX has the potential to 279
 235 significantly enhance knowledge of the initial magnetic field 280
 236 formed in heavy ion collisions through constraining the pa- 281
 237 rameters of this predicted effect. In the bottom figure, the

241 current sPHENIX projection for transverse single spin asym-
 242 metry, A_N , is made averaged over both D^0 and \bar{D}^0 measure-
 243 ments. If sufficient separation of the mesons is possible, this
 244 measurement can instead be separated allowing for further
 245 constraints on the model parameters in the trigluon correla-
 246 tion function [13].

245 Current work on separating the D^0 and \bar{D}^0 mesons is
 246 ongoing through another simulated data sample using p+p
 247 collisions at RHIC energies with pileup. Without PID, the
 248 reconstructed invariant mass of the D^0 becomes a powerful
 249 tool for separation. Using KFParticle to obtain daughter candi-
 250 date tracks, the K^\pm and the π^\pm mass assumption can be
 251 made for both tracks and, using the reconstructed momen-
 252 tum, used to create two different reconstructed mass values.
 253 Truth matched results have shown that the incorrect track
 254 mass assumptions will lead to a significant widening of mass
 255 distribution around the true mass of the D^0 , allowing cuts to
 256 be made on each of the two reconstructed masses to label the
 257 mother as a D^0 or \bar{D}^0 candidate. Basic cuts as well as multi-
 258 variate analyses are being tested using this method currently
 259 and preliminary work shows promising results which may allow
 sPHENIX to perform these measurements effectively.

260 5 Timeline and Progress

261 Fabrication, assembly, and installation of the sPHENIX de-
 262 tector systems are well underway in preparation for data-
 263 taking beginning in early 2023. In the assembly hall at RHIC,
 264 the outer support structure, outer hadronic calorimeter, and
 265 BaBar magnet are all firmly in place. This can be seen, mi-
 266 nus the now-completed top of the support structure, in Fig-
 267 ure 8. Outside the assembly hall, the remaining detectors
 268 are also rapidly approaching installation readiness. The inner
 269 hadronic calorimeter is completely assembled and installed
 270 in its support ring with pre-assembly-hall testing completed.
 271 All electromagnetic calorimeter sectors have been assembled
 272 and are nearly completely tested while installation practice
 273 is ongoing. The tracking detectors (MVTX, INTT, TPC, and
 274 TPOT) are all also well along the way towards being ready
 275 for installation in the coming months, and are still undergo-
 276 ing testing and assembly.

277 Although there is still significant work remaining for
 278 sPHENIX to be ready for day-1 data taking, the current
 279 progress puts the collaboration in position to be able to align
 280 with its projected commissioning and data taking schedule as
 281 shown in Figure 9.



FIGURE 8. Outer Hadronic Calorimeter completely assembled and installed in the sPHENIX assembly hall at RHIC

Year	Species	$\sqrt{s_{NN}}$ [GeV]	Cryo Weeks	Physics Weeks	Rec. Lum. $ z < 10$ cm	Samp. Lum. $ z < 10$ cm
2023	Au+Au	200	24 (28)	9 (13)	3.7 (5.7) nb ⁻¹	4.5 (6.9) nb ⁻¹
2024	$p^\uparrow p^\uparrow$	200	24 (28)	12 (16)	0.3 (0.4) pb ⁻¹ [5 kHz] 4.5 (6.2) pb ⁻¹ [10%-str]	45 (62) pb ⁻¹
2024	p^\uparrow +Au	200	-	5	0.003 pb ⁻¹ [5 kHz] 0.01 pb ⁻¹ [10%-str]	0.11 pb ⁻¹
2025	Au+Au	200	24 (28)	20.5 (24.5)	13 (15) nb ⁻¹	21 (25) nb ⁻¹

FIGURE 9. Proposed run schedule in the sPHENIX beam use proposal. Further details given in [2].

6 Conclusion

sPHENIX is rapidly approaching data taking and first analyses coming at the start of 2023. The collaboration has made great efforts to keep assembly and installation projects on schedule, while also preparing analysis and reconstruction modules to be ready for the first data taken. With these efforts, sPHENIX is prepared to make extremely precise and accurate measurements critical to the stated goals of the 2015 Long Range Plan for Nuclear Science [3] that also encompass numerous first-ever measurements made at RHIC such as the nuclear modification factor of the Upsilon 3S state, a baseline p+p measurement of the Λ_c baryon to D^0 meson production ratio, and possible separation of D^0 and \bar{D}^0 v_1 .

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