

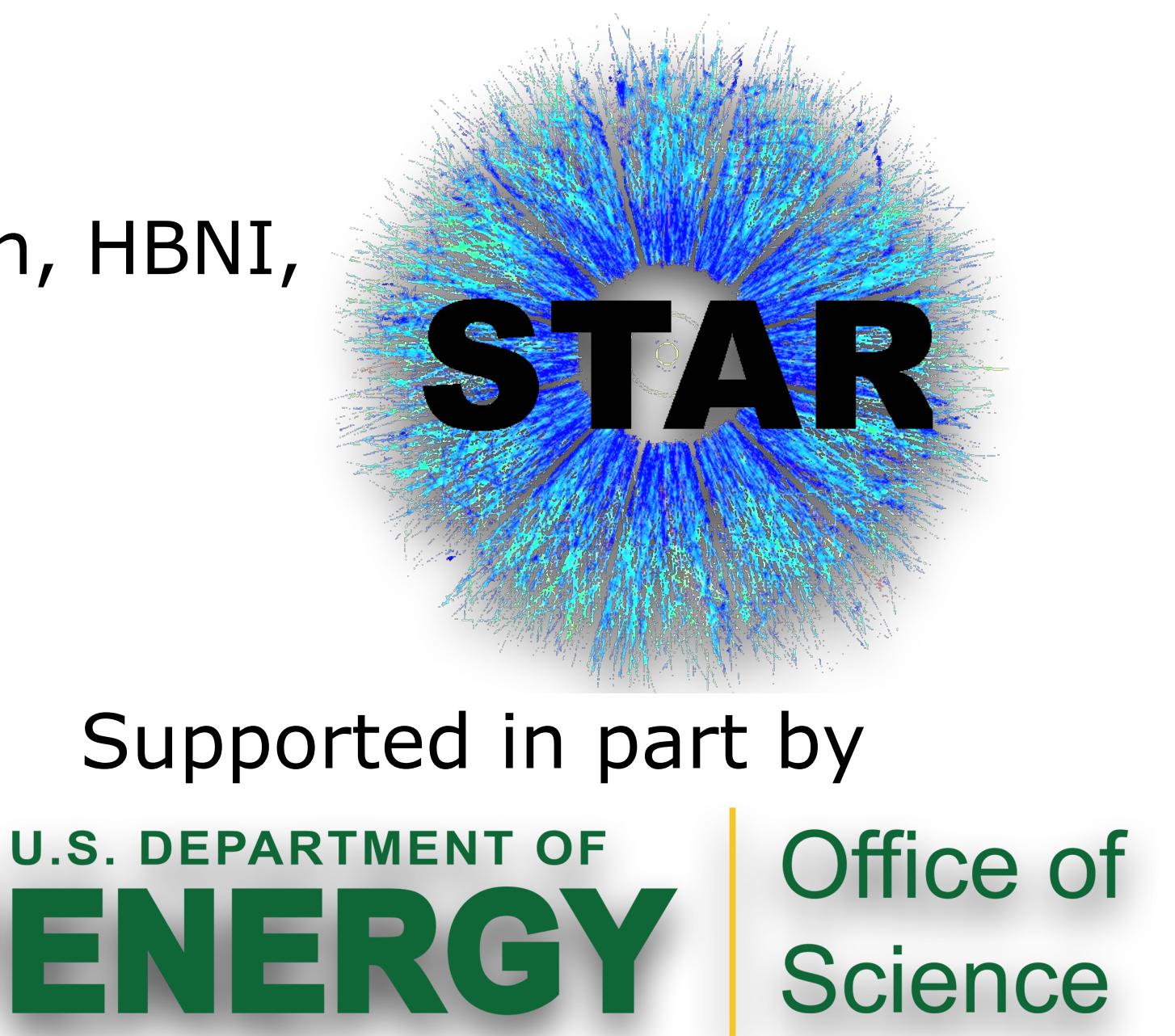
# Deuteron Number Fluctuations and Proton-deuteron Correlations in High Energy Heavy-ion Collisions in STAR Experiment at RHIC



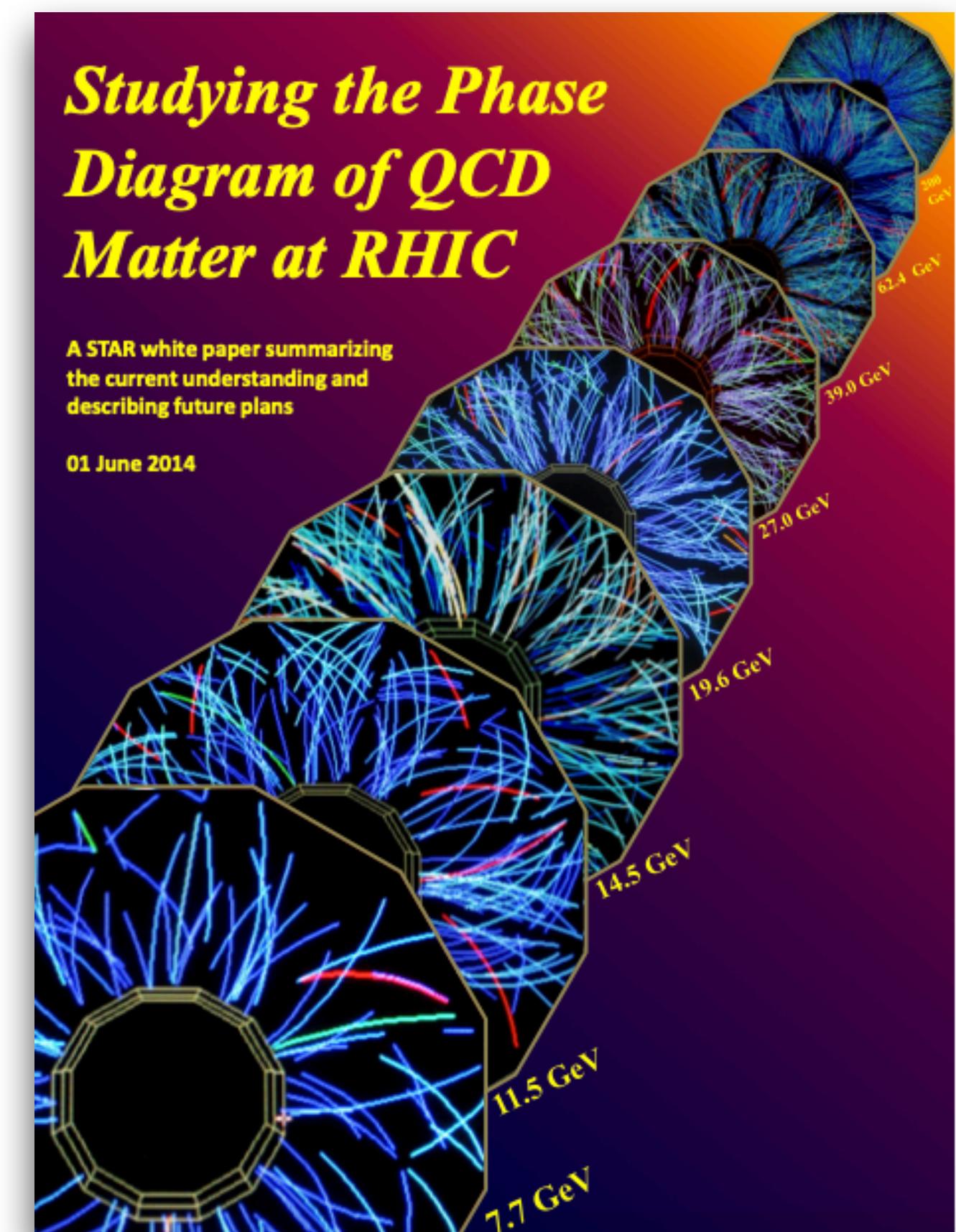
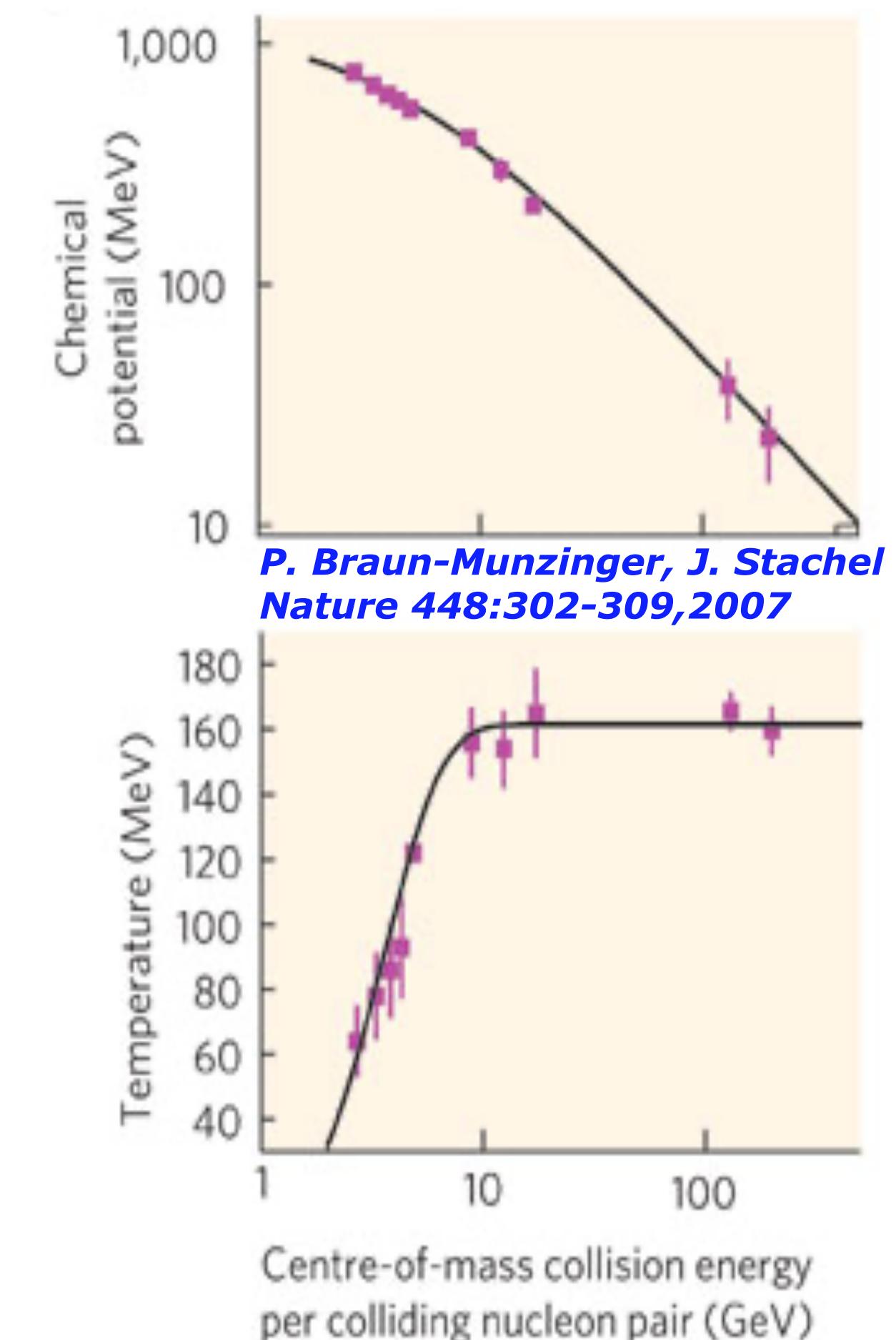
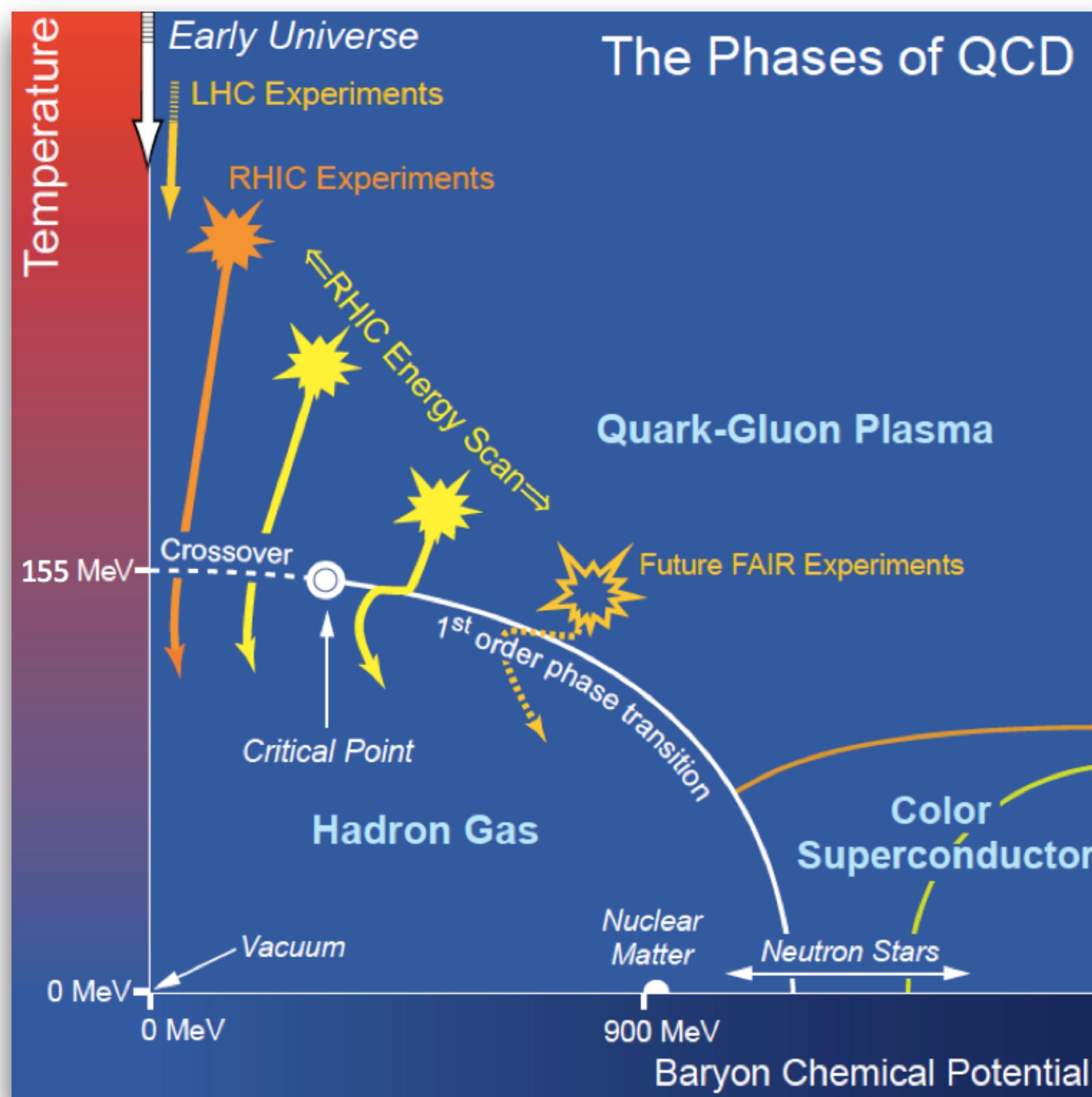
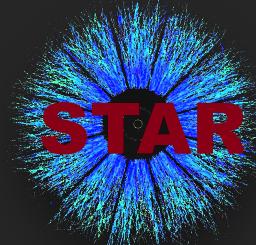
Debasish Mallick,  
On Behalf of the STAR Collaboration,  
National Institute of Science Education and Research, HBNI,  
Jatni, India

## Outline:

- Introduction
- Motivation and Observables
- STAR and Analysis Method
- Results
- Summary



# Introduction



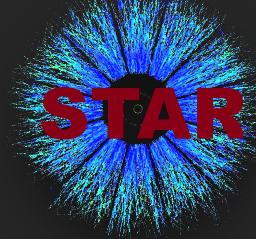
<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0493>  
[https://drupal.star.bnl.gov/STAR/files/BES\\_WPII\\_ver6.9\\_Cover.pdf](https://drupal.star.bnl.gov/STAR/files/BES_WPII_ver6.9_Cover.pdf)

Goal: Study the phase diagram of QCD.

Beam Energy Scan (BES): Varying beam energy varies temperature ( $T$ ) and baryon chemical potential ( $\mu_B$ ).

Fluctuations in conserved quantities are sensitive to phase structure and critical point.

# Light Nuclei Synthesis

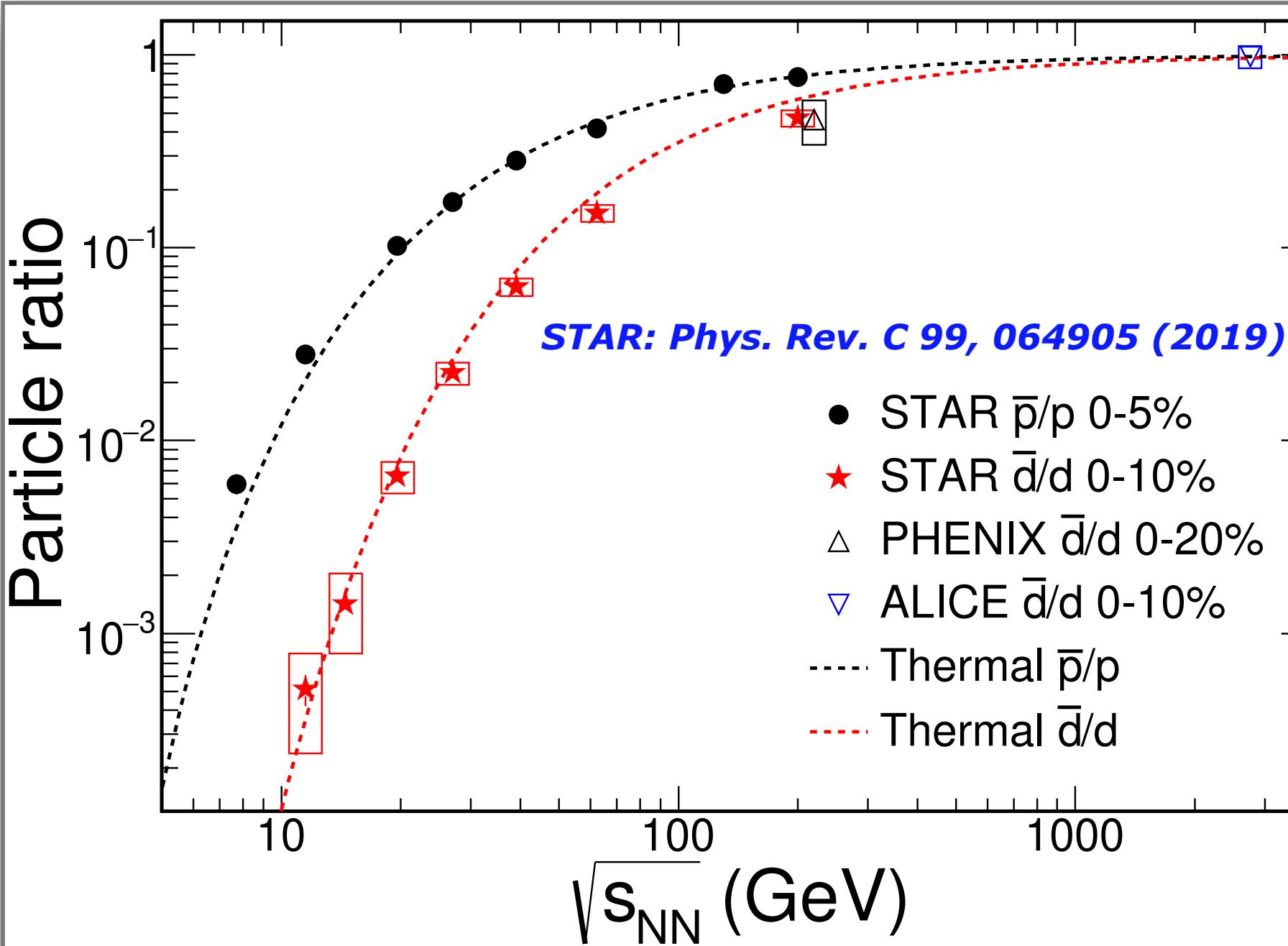


## GCE Thermal Model

**Yield of deuteron:**  $N_d = \frac{g_d V}{\pi^2} m_d^2 T K_2(m_d/T) \exp(\mu_d/T)$

where,  $g_d$ : degeneracy,  $\mu_d$ : chemical potential.

- Deuteron is treated as a free and point particle.
- Degeneracy, mass and baryon number are inputs.



- Anti-particle to particle ratio well explained by the thermal model for a range of  $\sqrt{s_{NN}}$ .

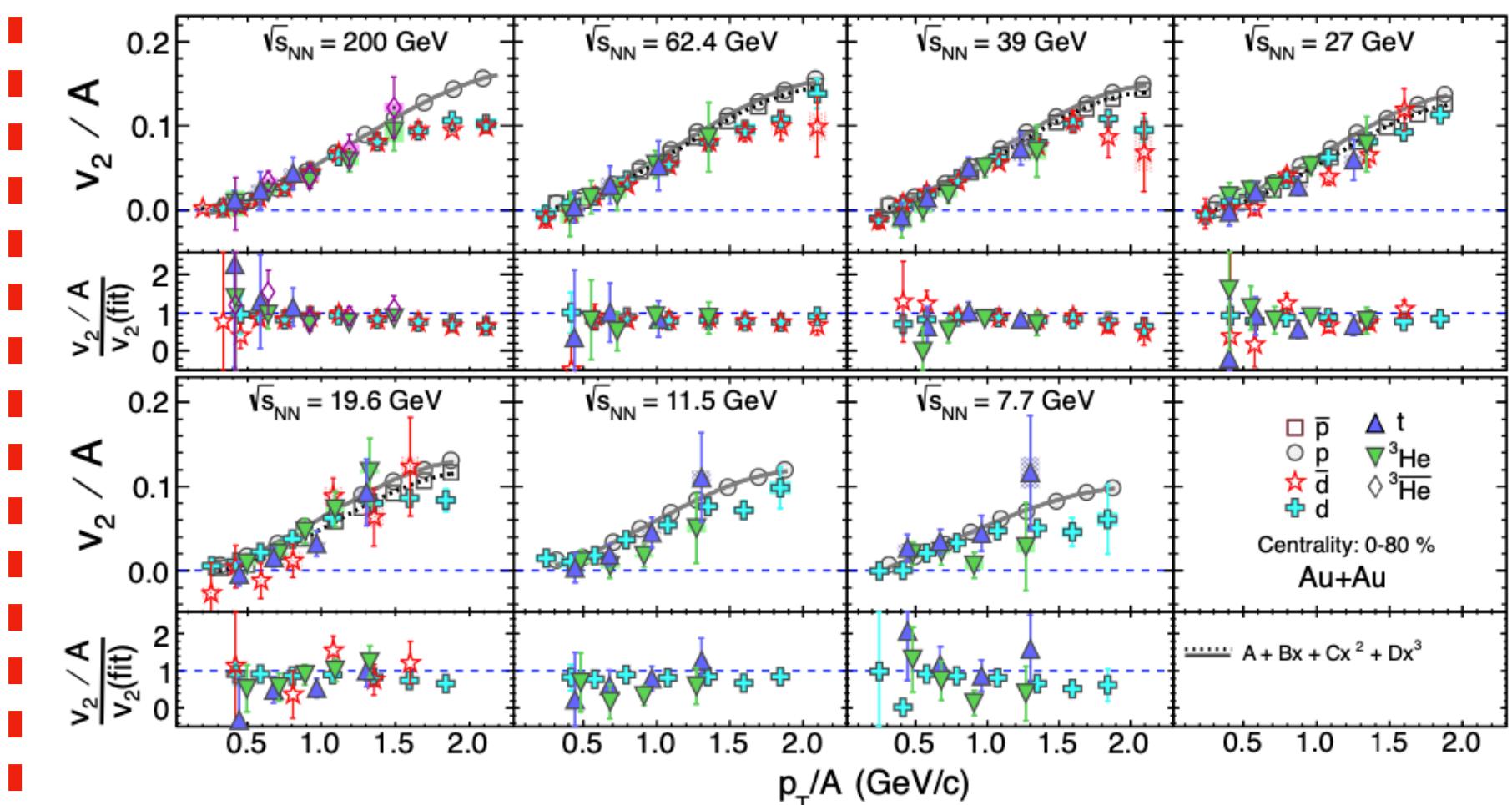
## Coalescence Model

**Invariant Yield:**  $E_d \frac{d^3 N_d}{dp_d^3} = B_2 \left( E_p \frac{d^3 N_p}{dp_p^3} \right) \left( E_n \frac{d^3 N_n}{dp_n^3} \right)$

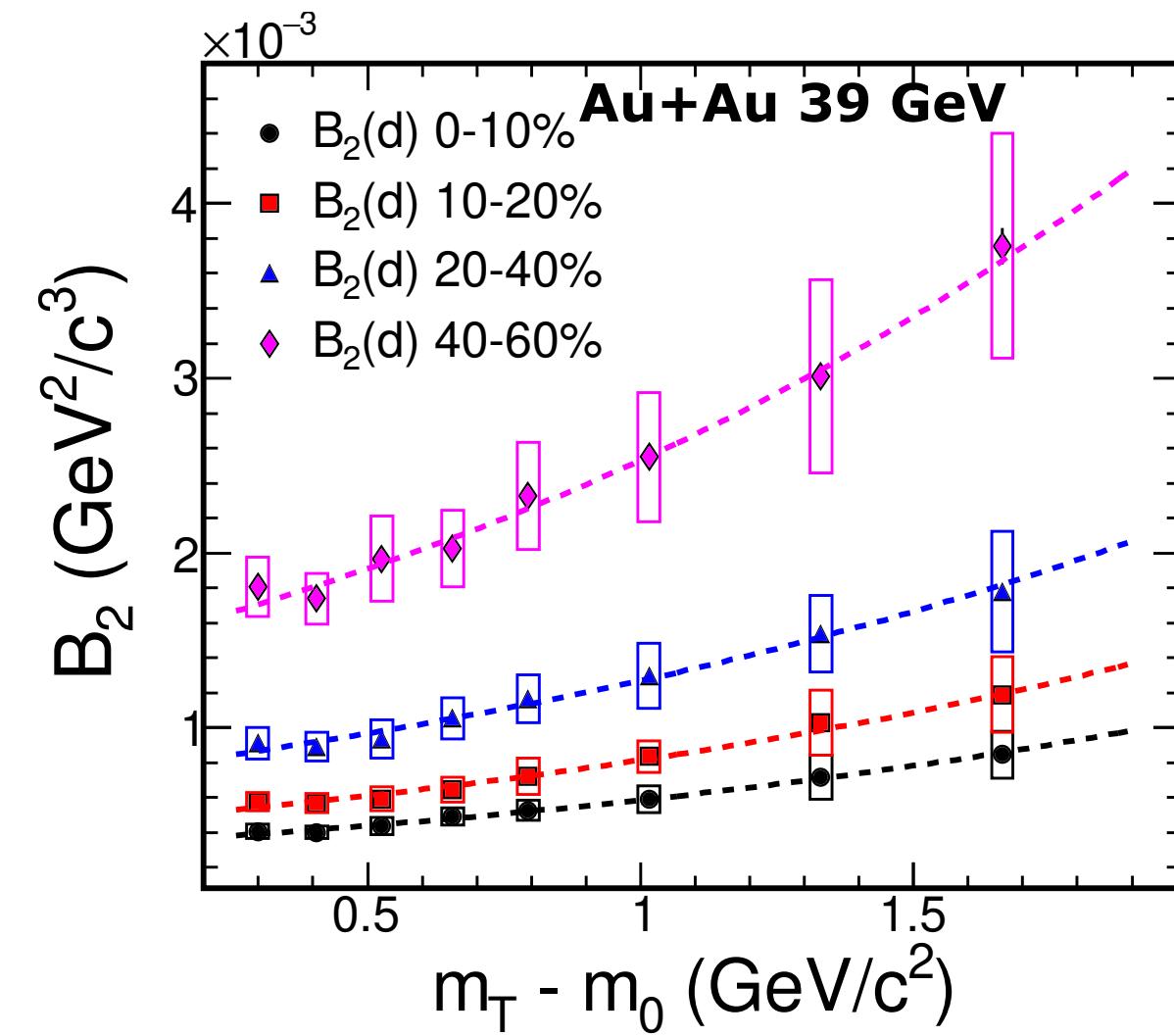
**Elliptic Flow:**  $v_2^d(p_T) \approx 2v_2^p \left( \frac{p_T}{2} \right)$

- Light nuclei created using protons and neutrons.
- $B_2$  extracted as a function of centrality,  $m_T$ , and  $\sqrt{s_{NN}}$ .

STAR: Phys. Rev. C 94 (2016) 3, 034908



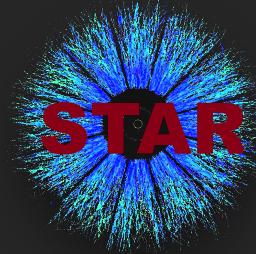
- Nucleon coalescence picture works up to  $p_T/A \leq 1.5$  GeV/c.



STAR: Phys. Rev. C 99, 064905 (2019)

- $B_2 \propto e^{(m_T - m)}$
- $B_2 \propto (4/3)\pi p_0^3$
- $p_0$  is the radius in momentum space.

# Light Nuclei Synthesis

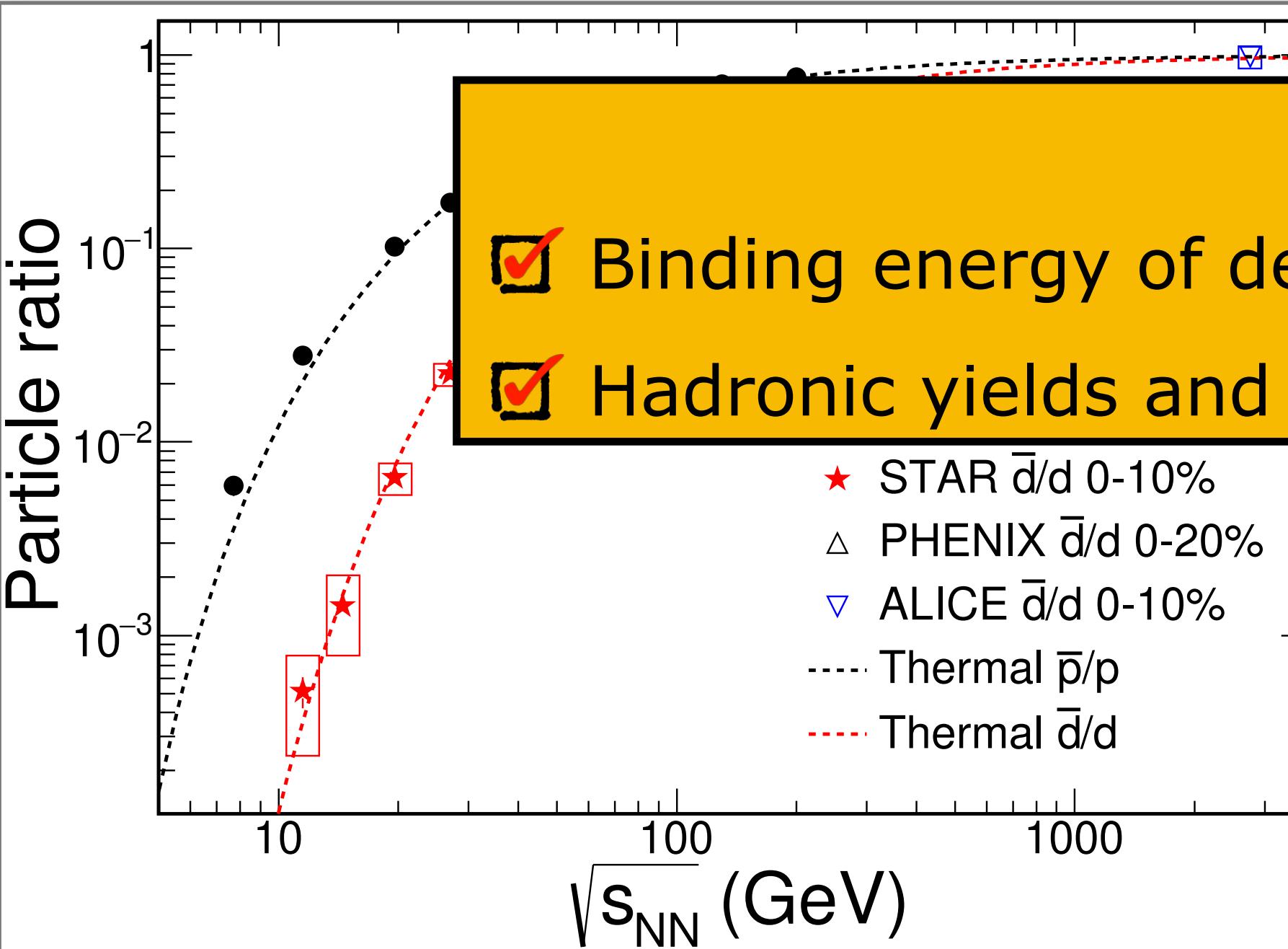


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## Coalescence Model

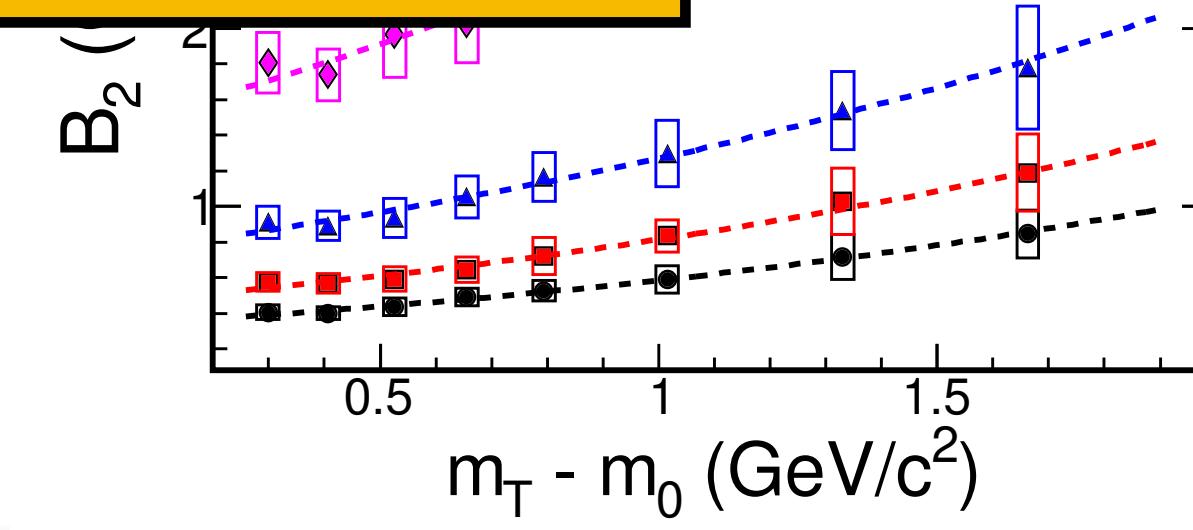
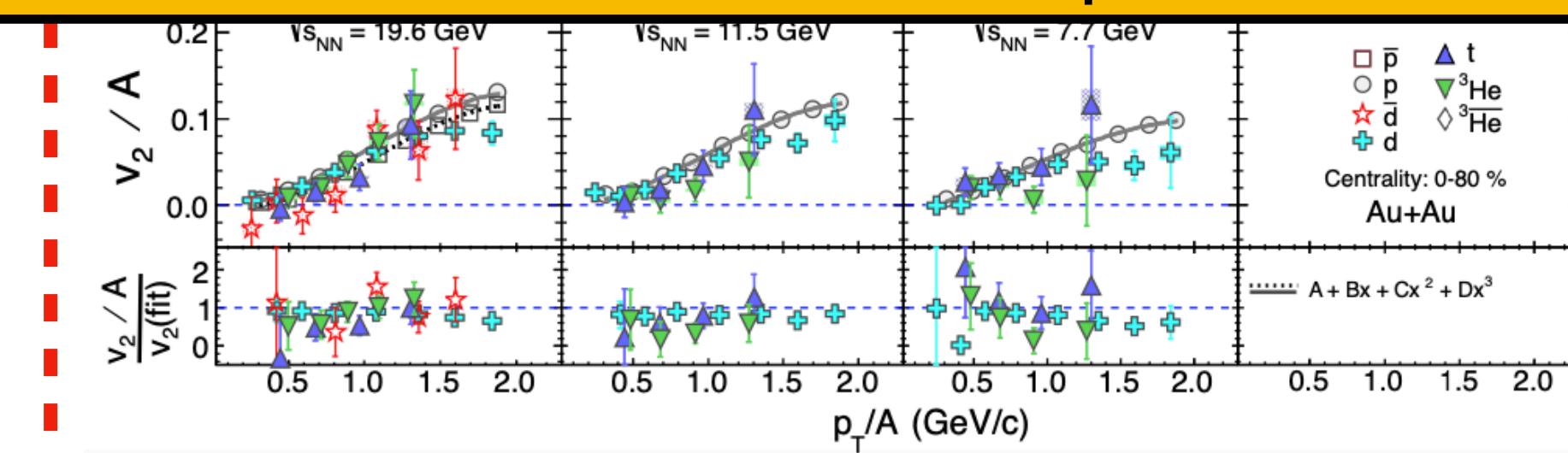
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**Elliptic Flow:**  $v_2^d(p_T) \approx 2v_2^p \left( \frac{p_T}{2} \right)$

- Light nuclei created using protons and neutrons.
- $B_2$  extracted as a function of centrality,  $m_T$ , and  $\sqrt{s_{NN}}$ .

## Typical Scales

- Binding energy of deuteron  $\sim 2.2$  MeV.
- Hadronic yields and spectra are fixed around temperature  $\sim 90 - 160$  MeV.



STAR: Phys. Rev. C 99, 064905 (2019)

- Nucleon coalescence picture works up to  $p_T/A \leq 1.5$  GeV/c.

$B_2 \propto e^{(m_T - m)}$

$B_2 \propto (4/3)\pi p_0^3$

$p_0$  is the radius in momentum space.

# Observables

- Higher-order cumulants characterise the subtle features of a distribution.

$$C_1 = \langle N \rangle$$

$$C_2 = \langle (\delta N)^2 \rangle$$

$$C_3 = \langle (\delta N)^3 \rangle$$

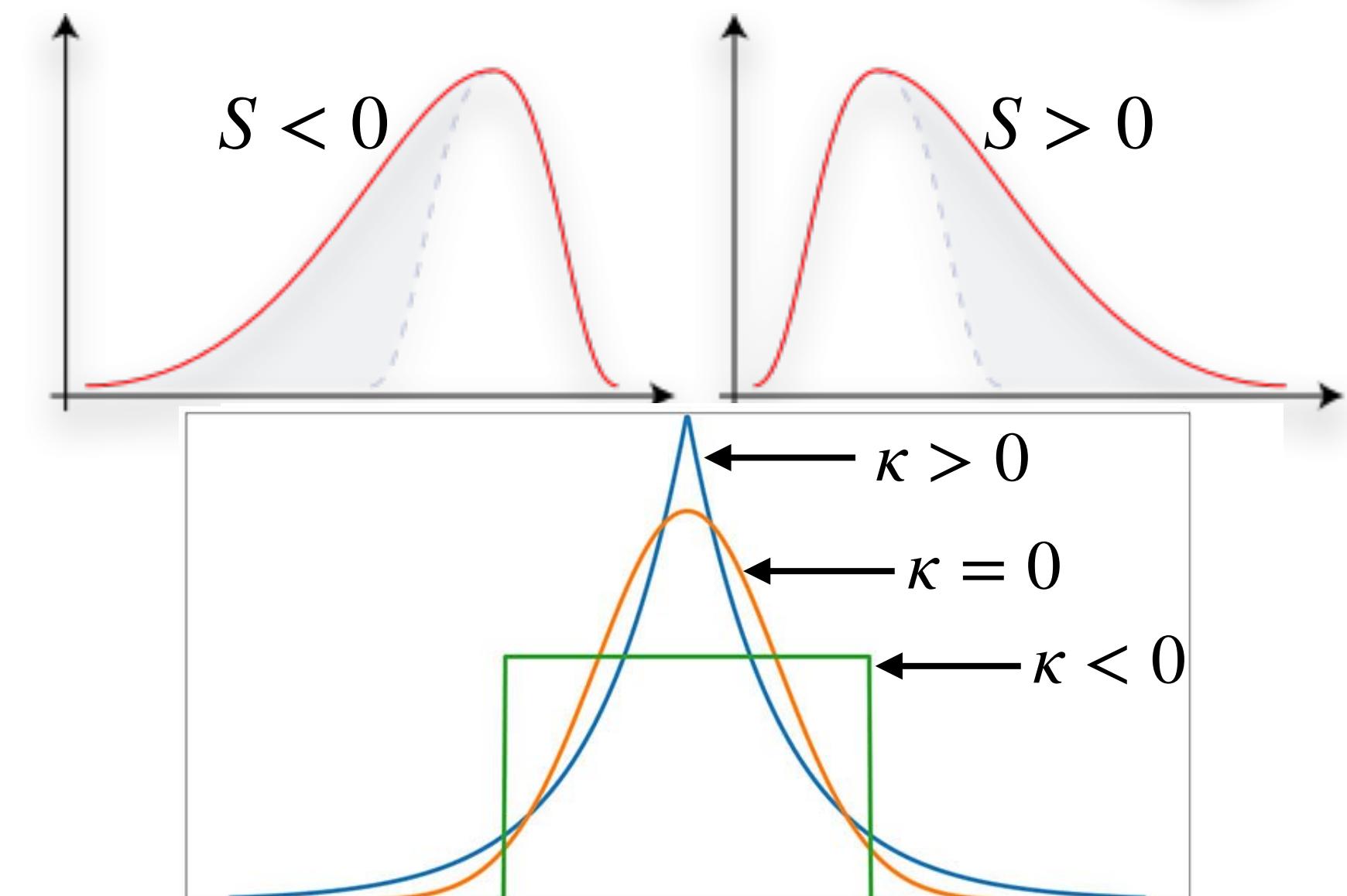
$$C_4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2$$

$$\frac{C_2}{C_1} = \frac{\sigma^2}{M}$$

$$\frac{C_3}{C_2} = S\sigma$$

$$\frac{C_4}{C_2} = \kappa\sigma^2$$

M = Mean  
 $\sigma^2$  = Variance  
 S = Skewness  
 $\kappa$  = Kurtosis



- Higher order cumulants of conserved number distributions are, in general, sensitive observables.
- Related to the correlation length and susceptibilities.
  - Deuteron cumulants add more information on baryon number fluctuation.

$$C_2 \sim \xi^2$$

$$C_4 \sim \xi^7$$

\*Quantitative numbers - Model dependent

$$\frac{\chi_q^{(4)}}{\chi_q^{(2)}} = \kappa\sigma^2 = \frac{C_{4,q}}{C_{2,q}}$$

$$\frac{\chi_q^{(3)}}{\chi_q^{(2)}} = S\sigma = \frac{C_{3,q}}{C_{2,q}}$$

S. Ejiri, F. Karsch, K. Redlich, Phys. Lett. B633 (2006) 275-282

M. A. Stephanov, Phys. Rev. Lett. 102, 032301 (2009)

R.V. Gavai, S. Gupta, Phys. Lett. B696:459-463, 2011

A. Bazavov et. al, Phys. Rev. Lett. 109, 192302 (2012)

A. Bzdak et. al, Physics Reports 853 (2020) pp. 1-87

S. Borsanyi et. al, Phys. Rev. Lett. 111, 062005 (2013)

Pearson correlation coefficient

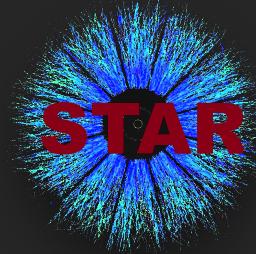
$$\rho(N_x, N_y) = \frac{\langle (\delta N_x \delta N_y) \rangle}{\sigma_x \sigma_y}$$

$\rho$  measures linear correlation between two variables.

$\rho > 0$ : Positive correlation

$\rho < 0$ : Anti-correlation

# Fluctuation as Probe of Synthesis Mechanism



## Coalescence Toy Model

Z. Fecková, J. Steinheimer, B. Tomášik and M. Bleicher: Phys. Rev. C 93, 054906 (2016)

Probability of deuteron formation,  $\lambda_d = B_2 n_p n_n$

Assume, proton ( $n_p$ ) and neutron ( $n_n$ ) follow Poisson distributions,

- At low  $\sqrt{s_{NN}}$ ,  $B_2$  increases. STAR: Phys. Rev. C 99, 064905 (2019)
- Larger value of  $n_p$  and  $n_n$  at low  $\sqrt{s_{NN}}$ .
- Results in rise of scaled moments of deuteron number.

**Scaled Moments:**  $\sigma^2/M = C_2/C_1$  ,  $S\sigma = C_3/C_2$  ,  $\kappa\sigma^2 = C_4/C_2$

### Two assumptions in the model:

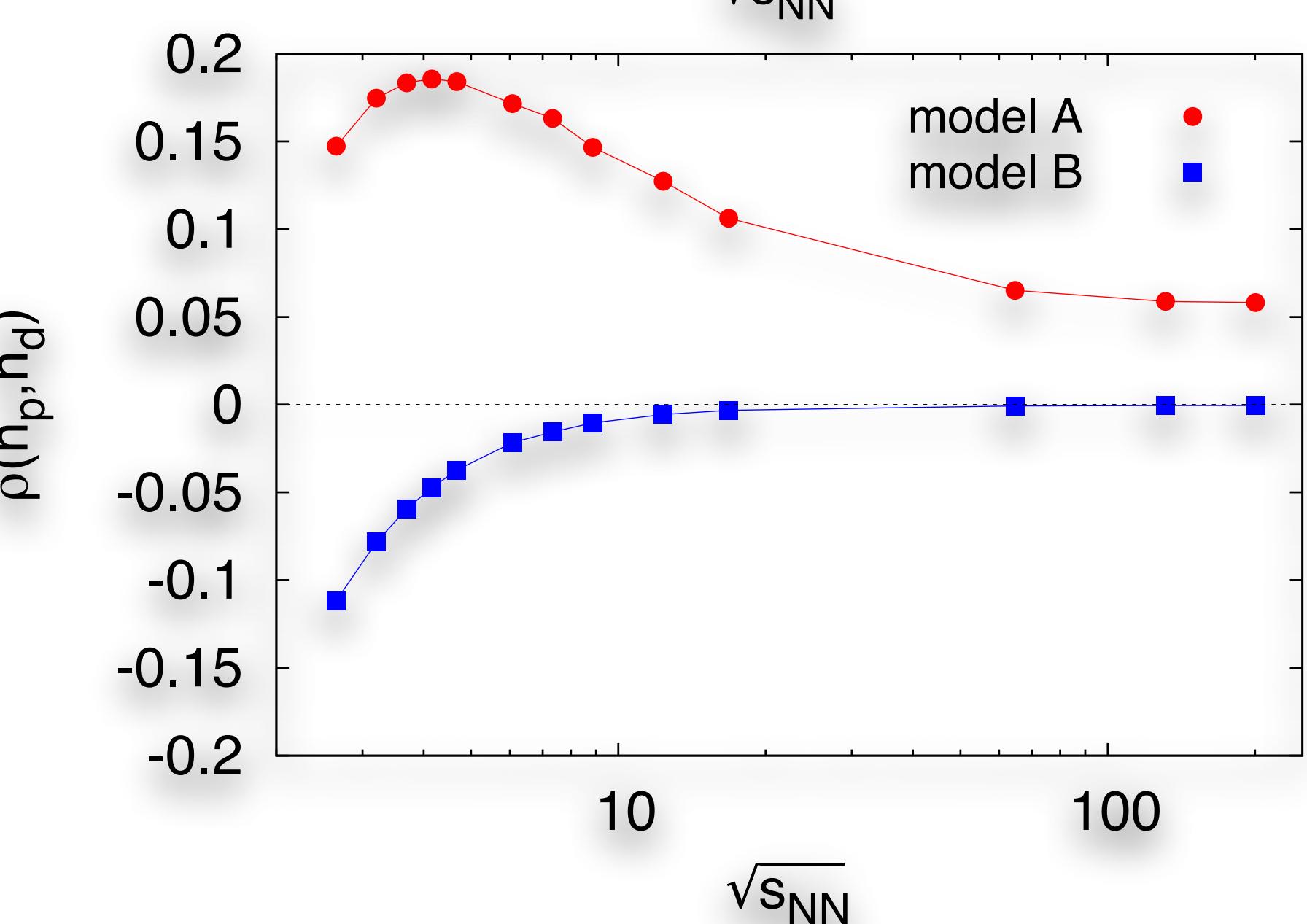
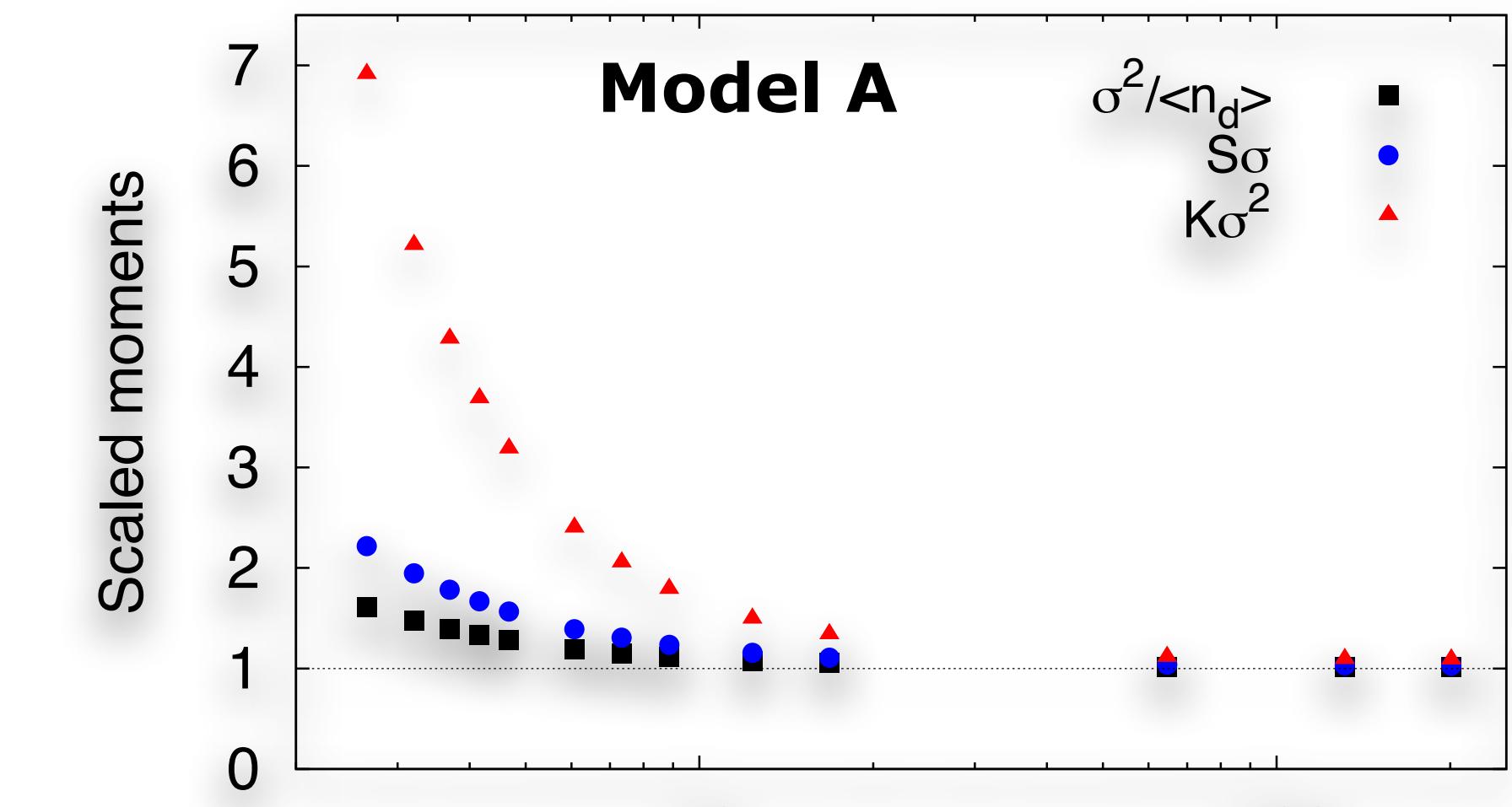
Model A: Correlated p and n ( $n_p=n_n$ ).      Model B: Independent p and n.

$$\lambda_d = B_2 n_p^2$$

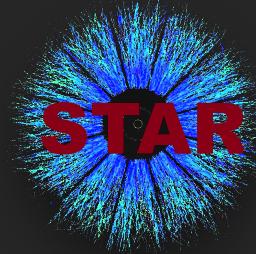
$$\lambda_d = B_2 n_p n_n$$

$$\rho(n_p, n_d) = \frac{\langle (n_p - \langle n_p \rangle)(n_d - \langle n_d \rangle) \rangle}{\sigma_p \sigma_d}$$

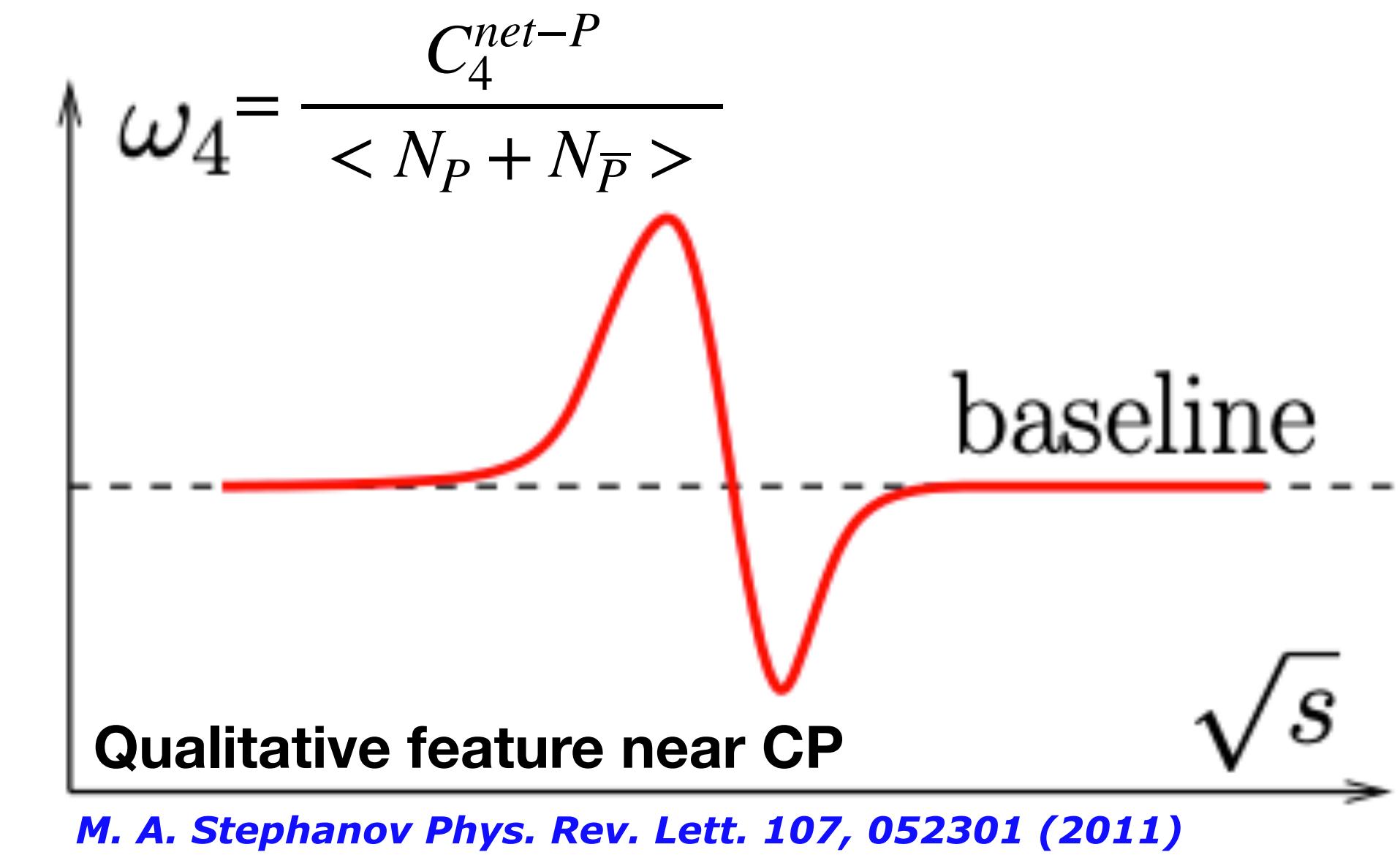
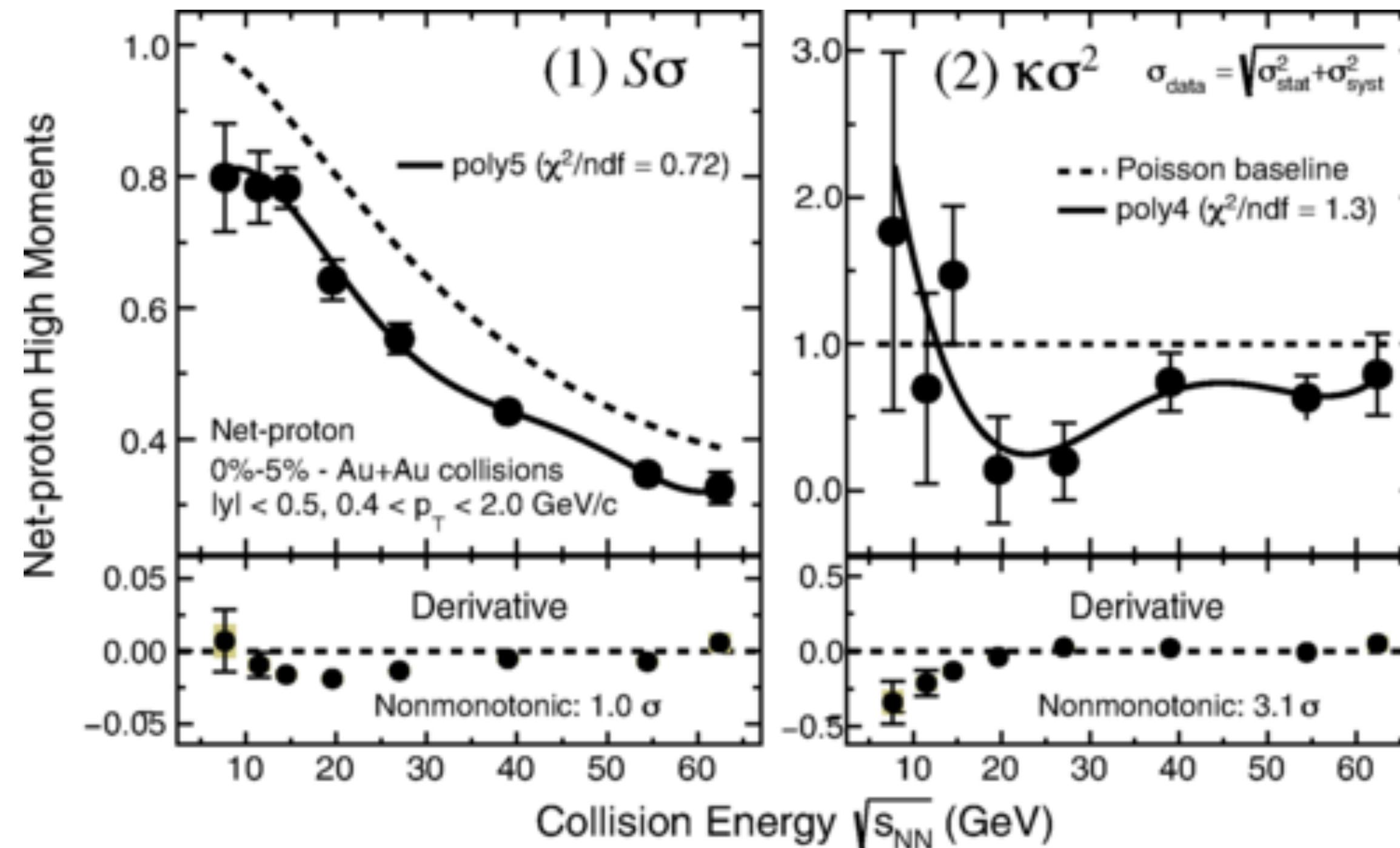
- Model A:  $\rho > 0$
- Model B:  $\rho < 0$



# Baryon Number Fluctuation



*STAR: Phys. Rev. Lett. 126 (2021) 092301*



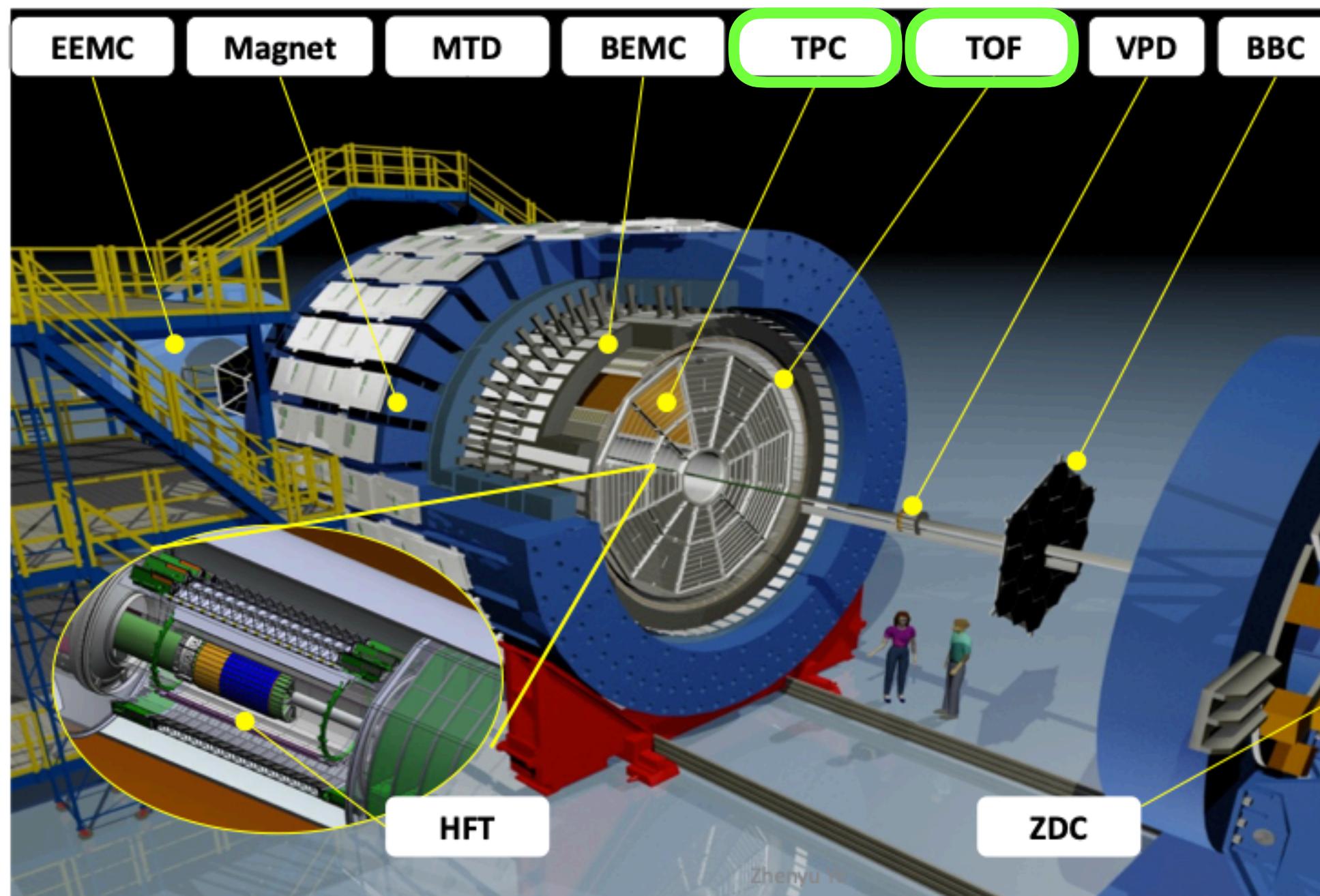
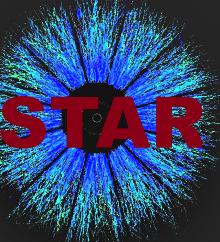
*M. A. Stephanov Phys. Rev. Lett. 107, 052301 (2011)*

*M. A. Stephanov 2011 J. Phys. G: Nucl. Part. Phys. 38 124147*

- Cumulants of deuteron number distribution and proton-deuteron correlation are sensitive to production mechanism.
- Until now studies have been done only with baryons of  $|B|=1$ .
- QCD critical point leads to large density fluctuation within certain correlation length.  
Deuteron production might be affected by local density fluctuations.

*Ed. Shuryak et. al, Phys. Rev. C 101 (2020) 3, 034914*  
*K.J. Sun et. al, Phys. Lett. B 774 (2017) 103-107*

# STAR Detector



*STAR: Nucl.Instrum.Meth.A 499 (2003) 624-632*

**PID and Centrality:** Using both Time Projection Chamber (TPC) and Time-of-Flight (ToF) detectors.

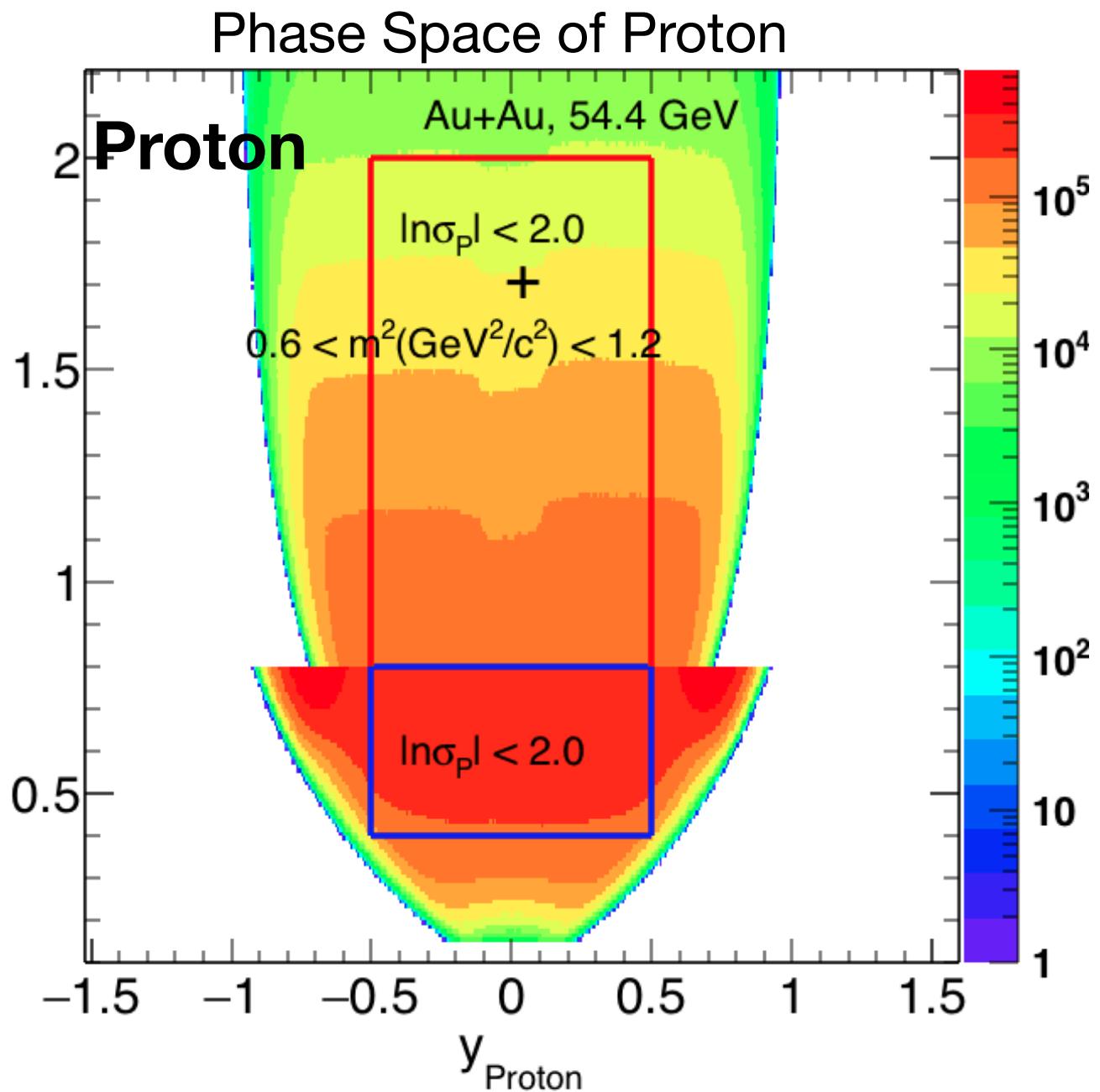
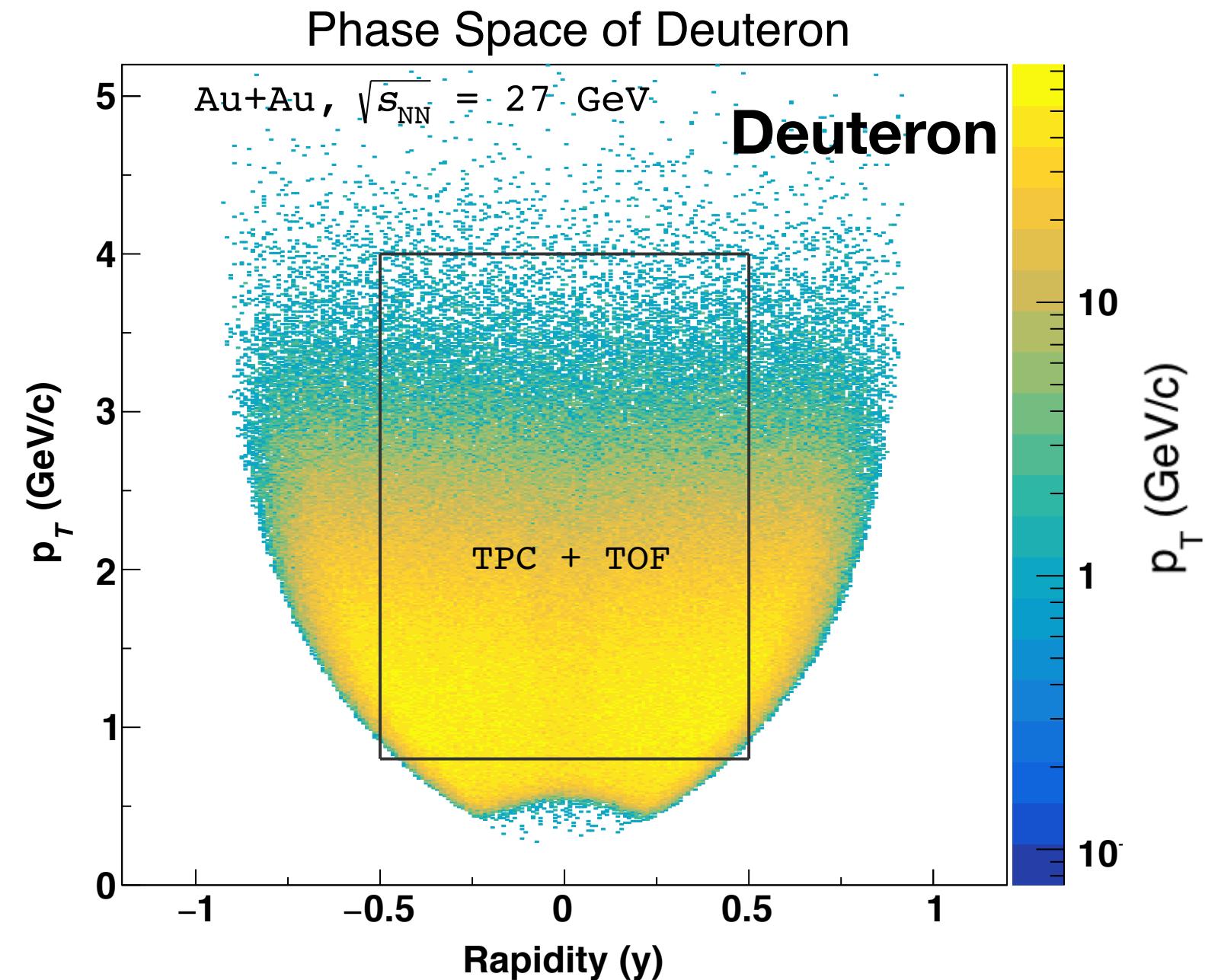
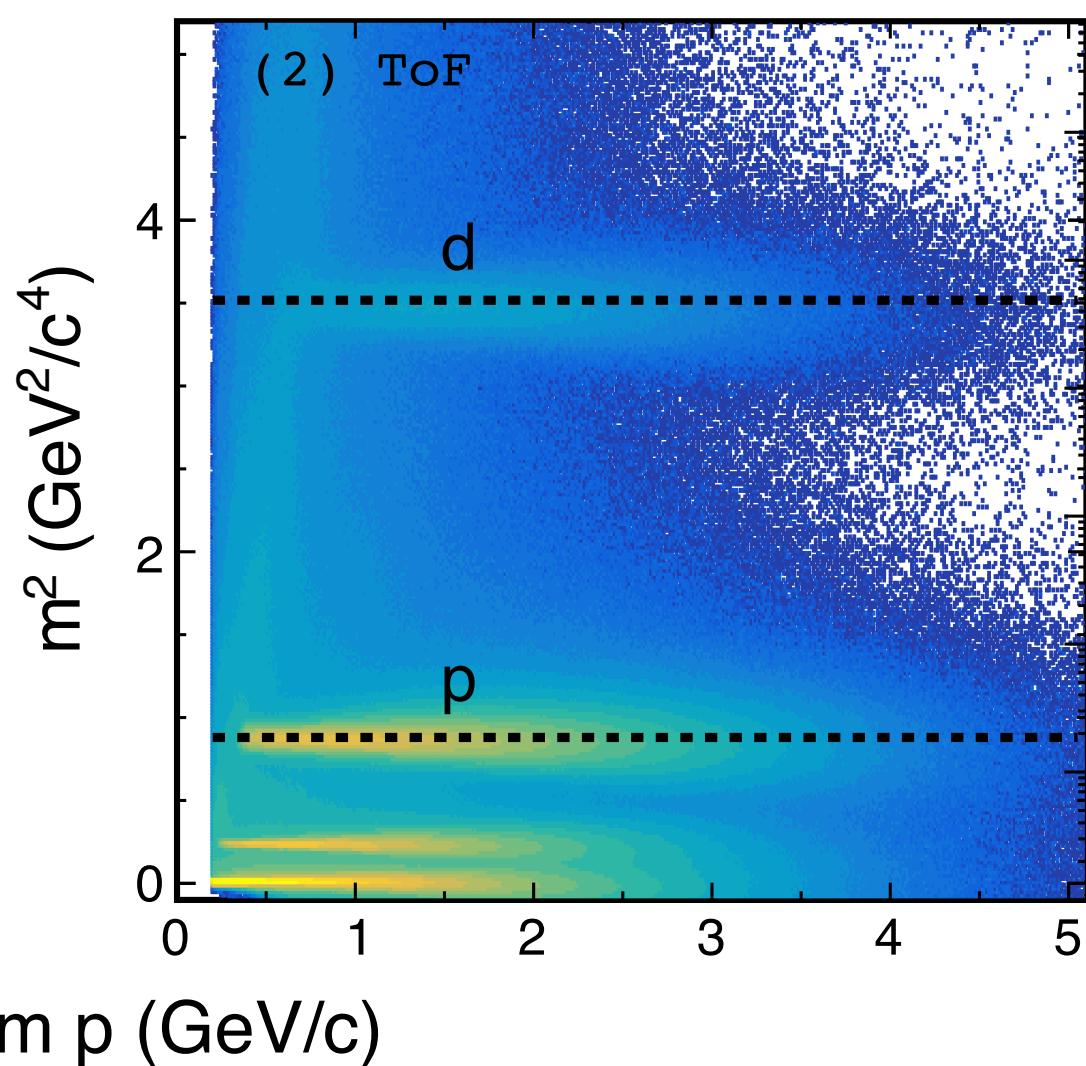
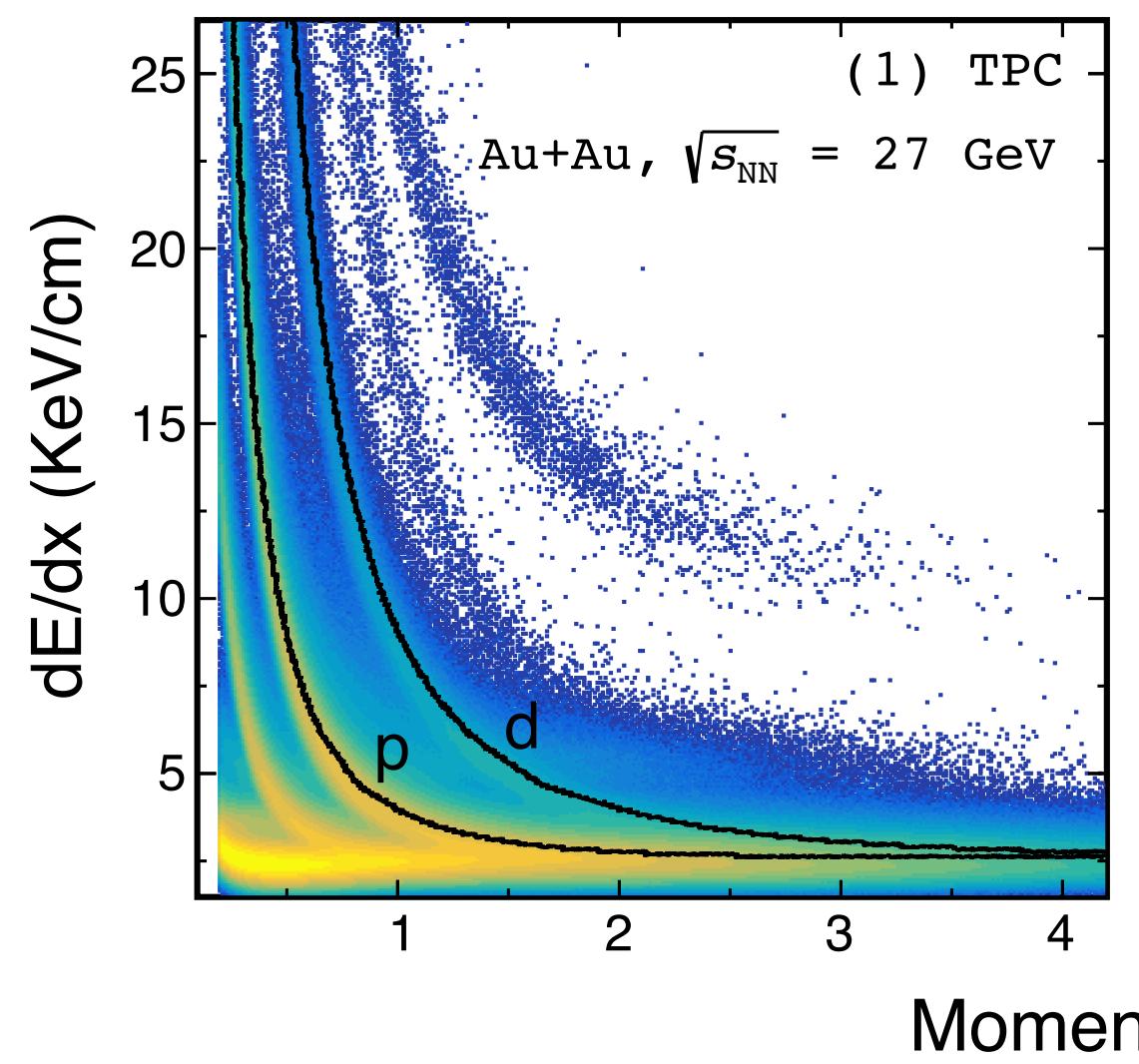
**Uniform coverage for full azimuth and  $|\eta| < 1$ .**  
**Excellent PID capability.**

## Dataset: BES-I

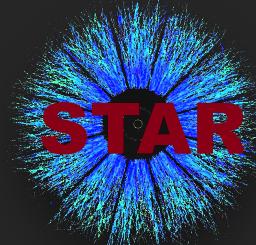
Collision system: Au+Au collision (centrality: 0-5% , 70-80%)

CoM energy: 7.7, 11.5, 14.5, 19.6, 27, 39, 54.4, 62.4, 200 GeV

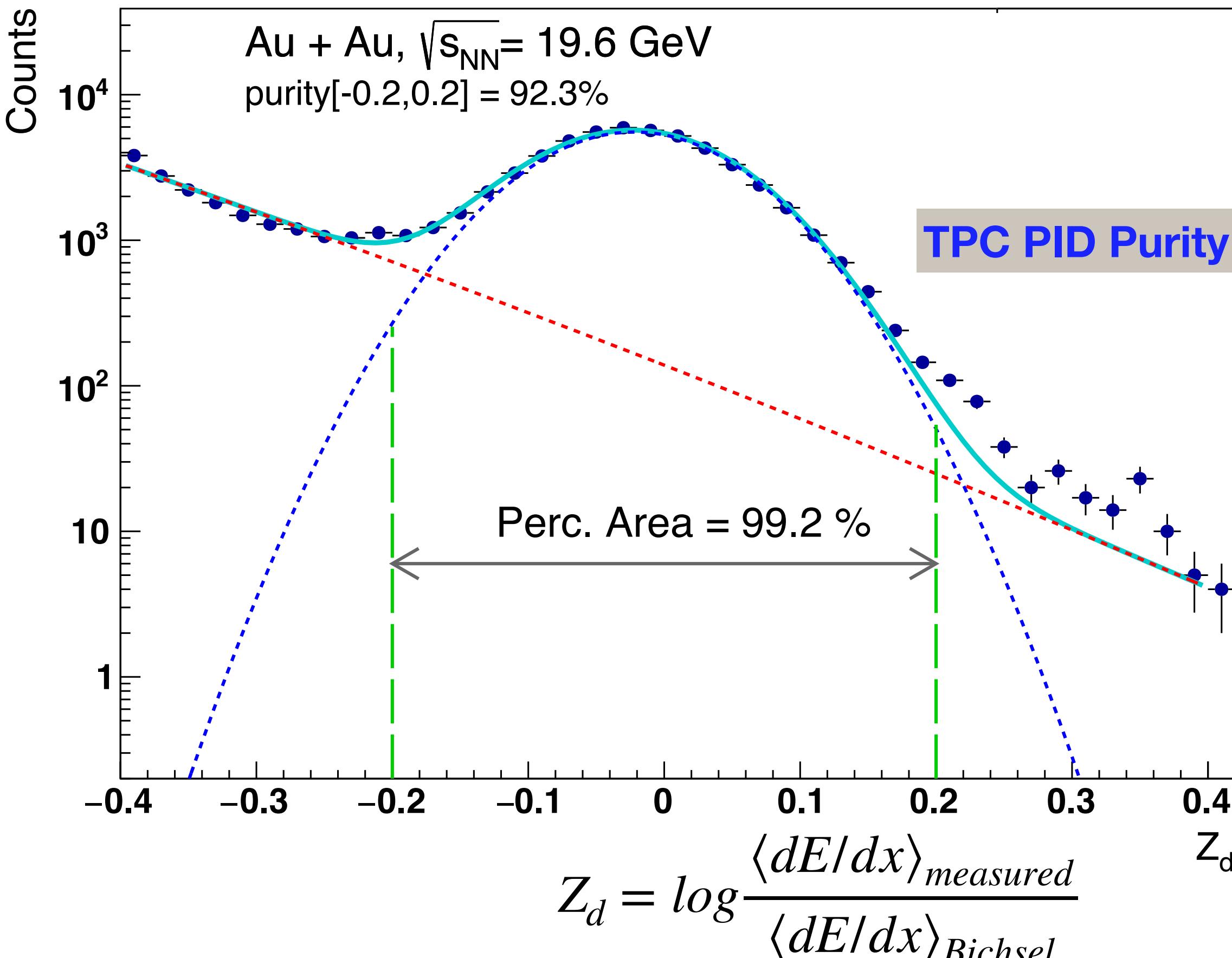
Year : 2010 – 2017



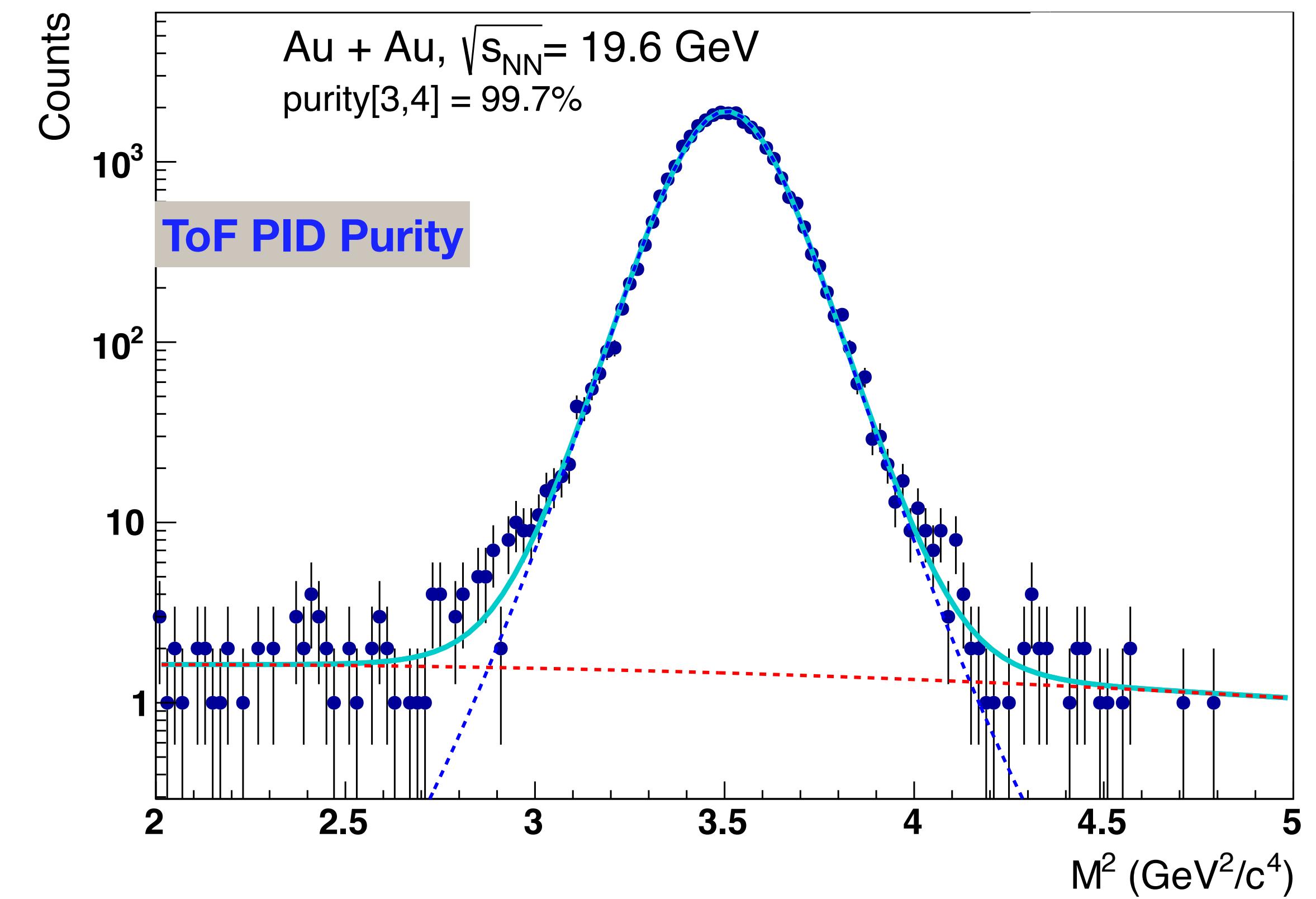
# Purity



$Z_d$  Distribution, 0-10% , $0.8 < p_T < 1.0 \text{ GeV}/c$ ,  $|y| < 0.5$



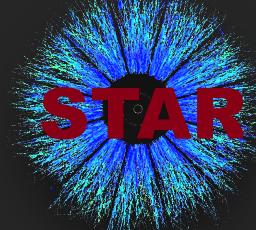
$m^2$  Distribution, 0-10% , $0.8 < p_T < 1.0 \text{ GeV}/c$ ,  $|y| < 0.5$



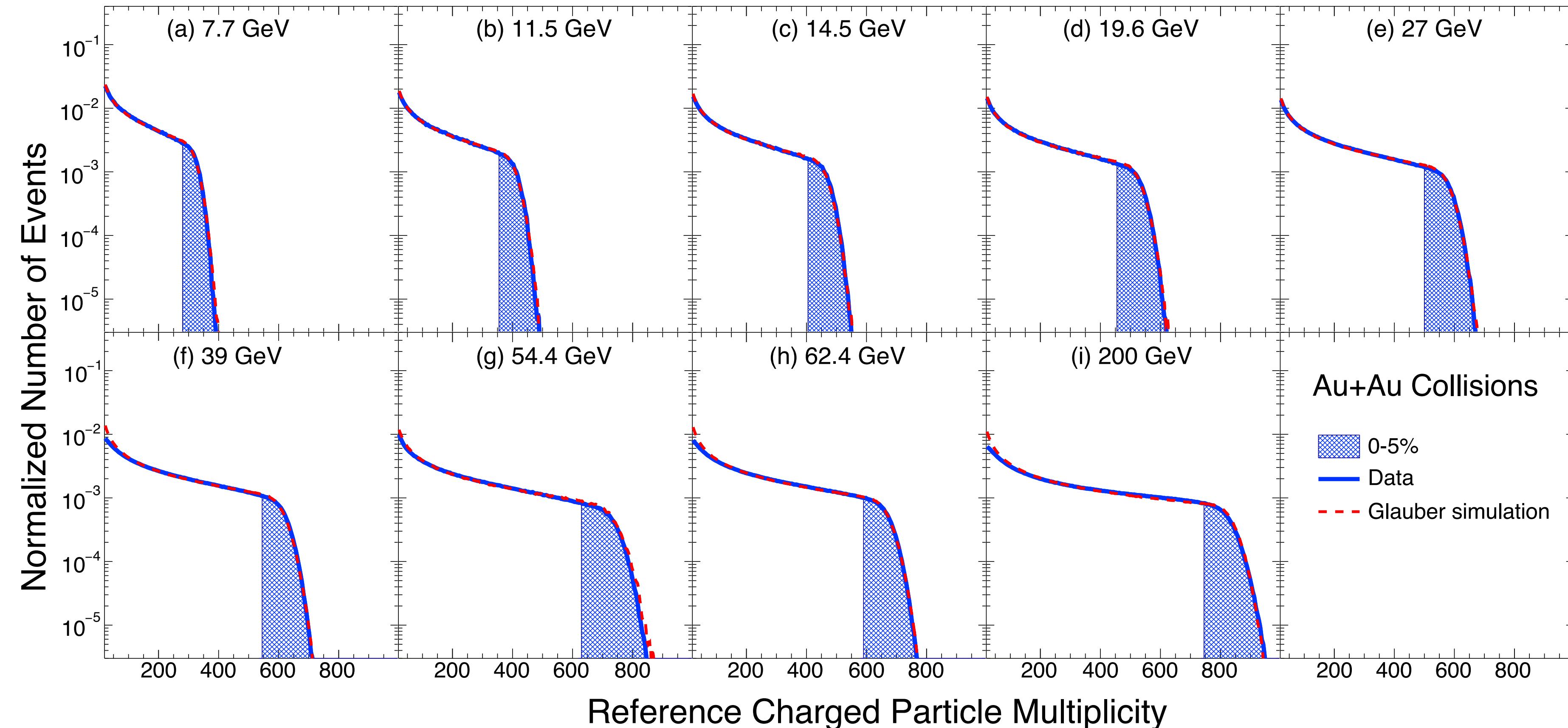
To achieve better PID purity, deuterons are identified using always both TPC and ToF detector.

Distance of Closest Approach (DCA) is kept as DCA<1cm to reduce the background contribution.

# Centrality Definition



Centrality using charged particles within  $|\eta| < 1.0$ , **excluding protons and deuterons**



Au+Au Collisions

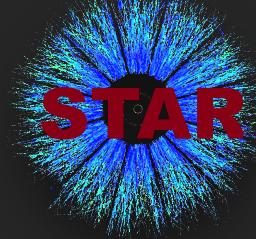
- 0-5%
- Data
- - Glauber simulation

Charged particle multiplicity is corrected for  
(a) Collision vertex and  
(b) Beam luminosity dependencies

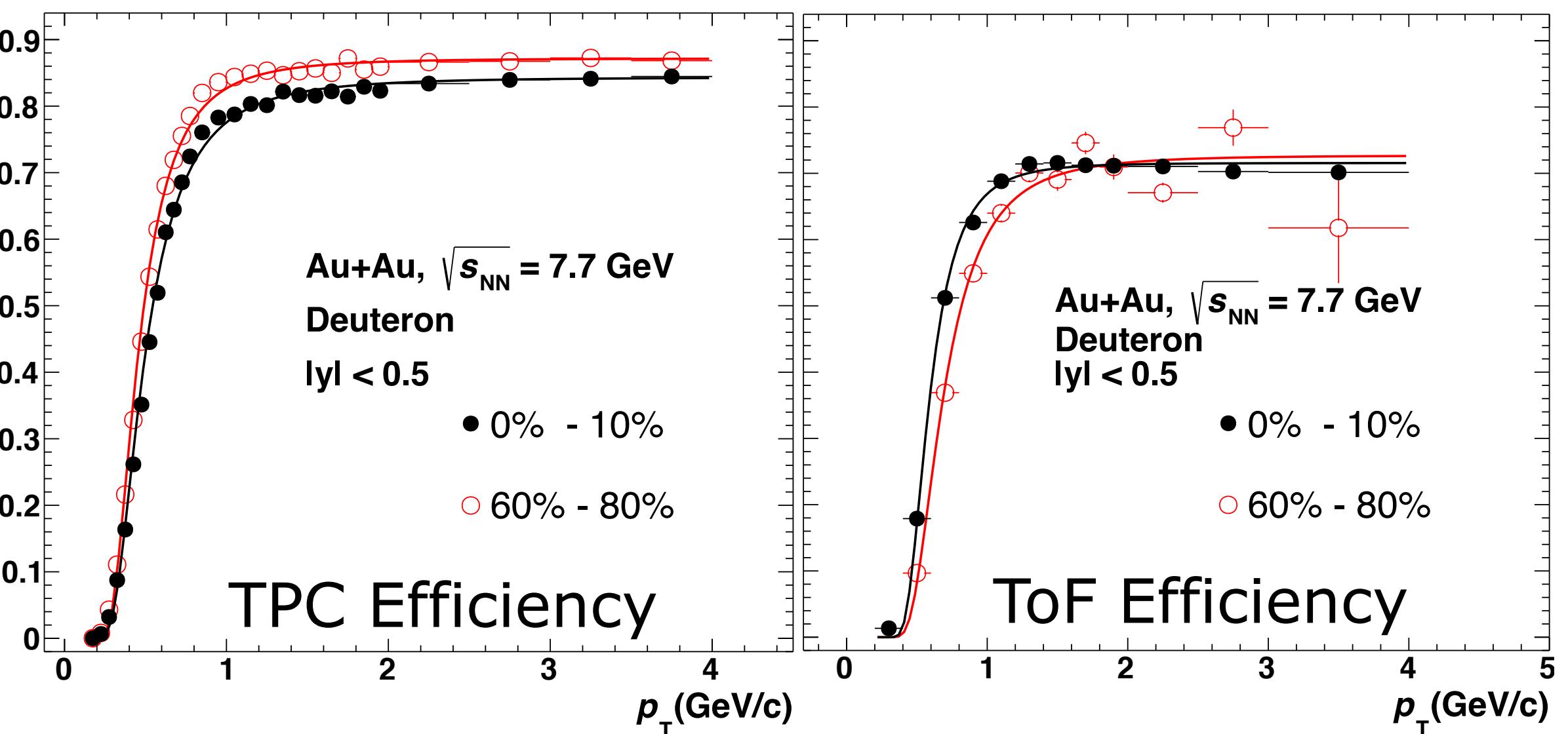
STAR: Phys. Rev. C 104, 024902 (2021)

This definition excludes self/auto -correlations between centrality and particle of interest.

# Analysis Methods



## 1) Detection efficiency correction - binomial model



## 2) Centrality bin-width (CBW) correction:

❖ Effect arises from the dependence of  $C_n$  on multiplicity.

$$C_n = \sum_r \omega_r C_{n,r}, \quad \omega_r = \frac{n_r}{\sum_r n_r}.$$

$n_r$  is number of events in r-th multiplicity bin.

## 3) Statistical uncertainty:

Using re-sampling technique called Bootstrap method.

For a statistic  $X$ ,  $\text{Var}(X) = \frac{1}{S-1} \sum_{s=1}^S (X_s^* - \bar{X})^2$ .

$S$  is the number of samples.

$X_s^*$  is " $X$ " measured from s-th sample.

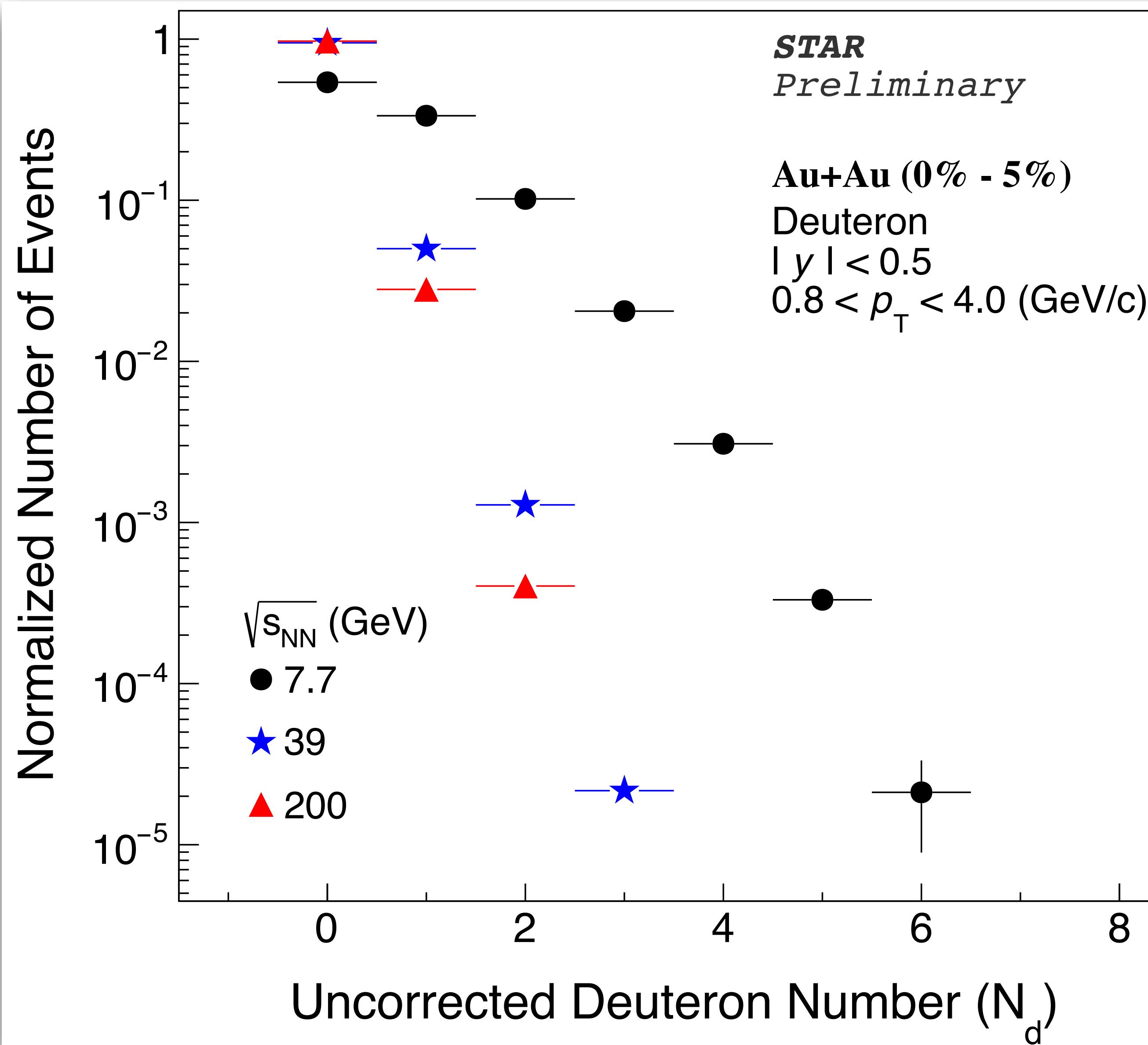
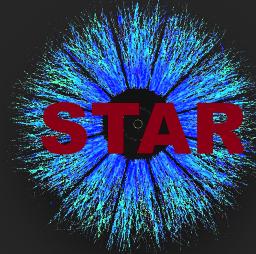
## 4) Systematic uncertainty:

Sources:

- Particle identification from TPC and ToF
- Background/decay estimates (DCA)
- Quality cuts for track reconstruction
- Uncertainty in detection efficiency estimation

**STAR:** Phys. Rev. C 104, 024902 (2021)  
 X. Luo , Phys. Rev. C 91, (2015) 034907  
 T. Nonaka et al, Phys. Rev. C 95, (2017) 064912  
 X. Luo et al, J.Phys. G 40, 105104 (2013)  
 X. Luo, J. Phys. G 39, 025008 (2012)  
 X.Luo et al, Phys.Rev. C99 (2019) no.4, 044917  
 A.Pandav et al, Nucl. Phys. A 991, (2019)121608

# Raw Deuteron Number Distribution

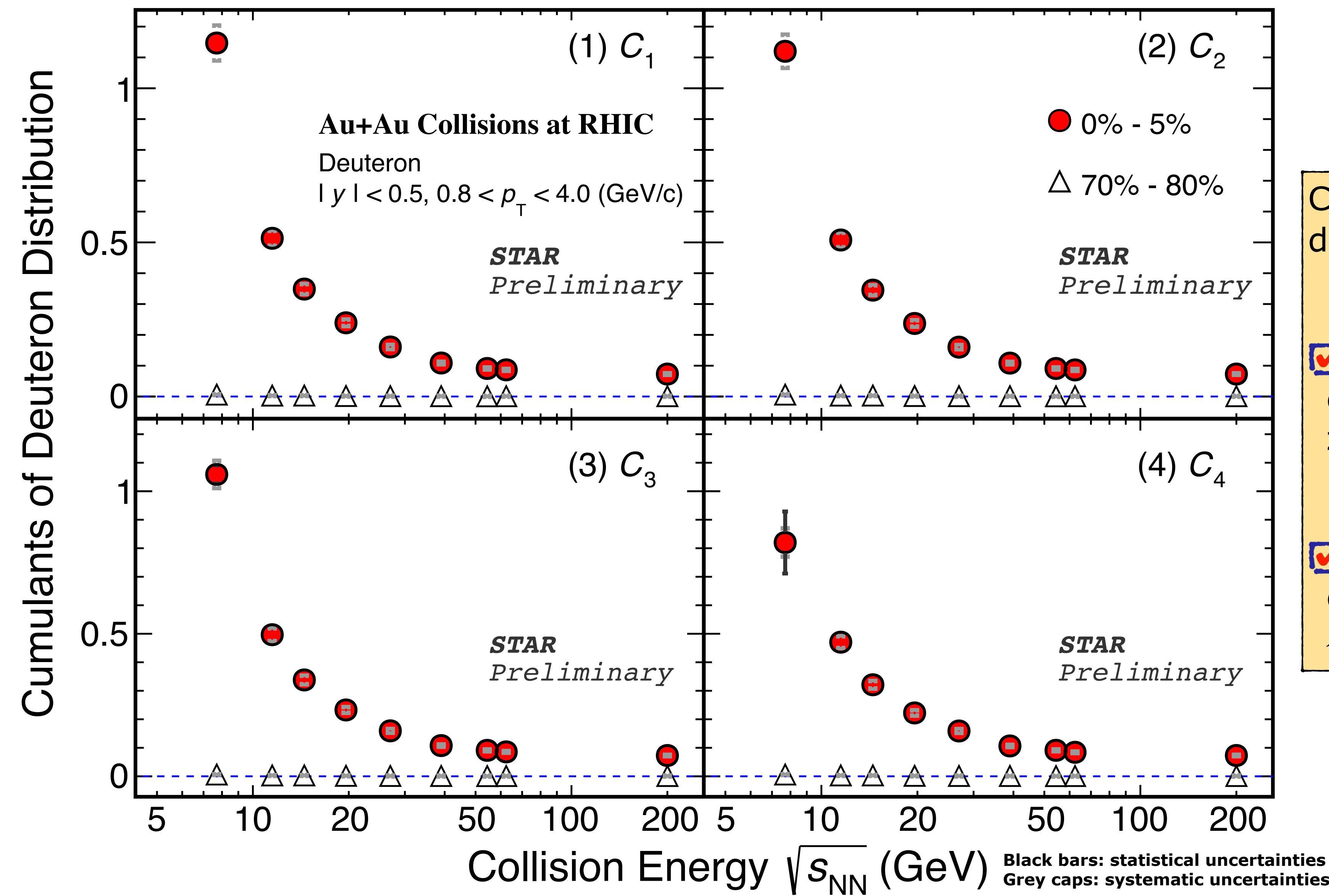
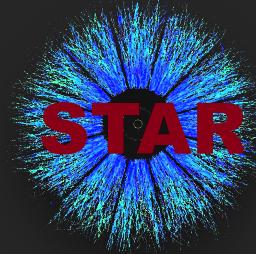


Uncorrected for efficiency and CBW effect.

Deuteron production increases towards low  $\sqrt{s_{NN}}$ .

Events with zero  $N_d$  value are most probable.

# Cumulants of Deuteron Distribution

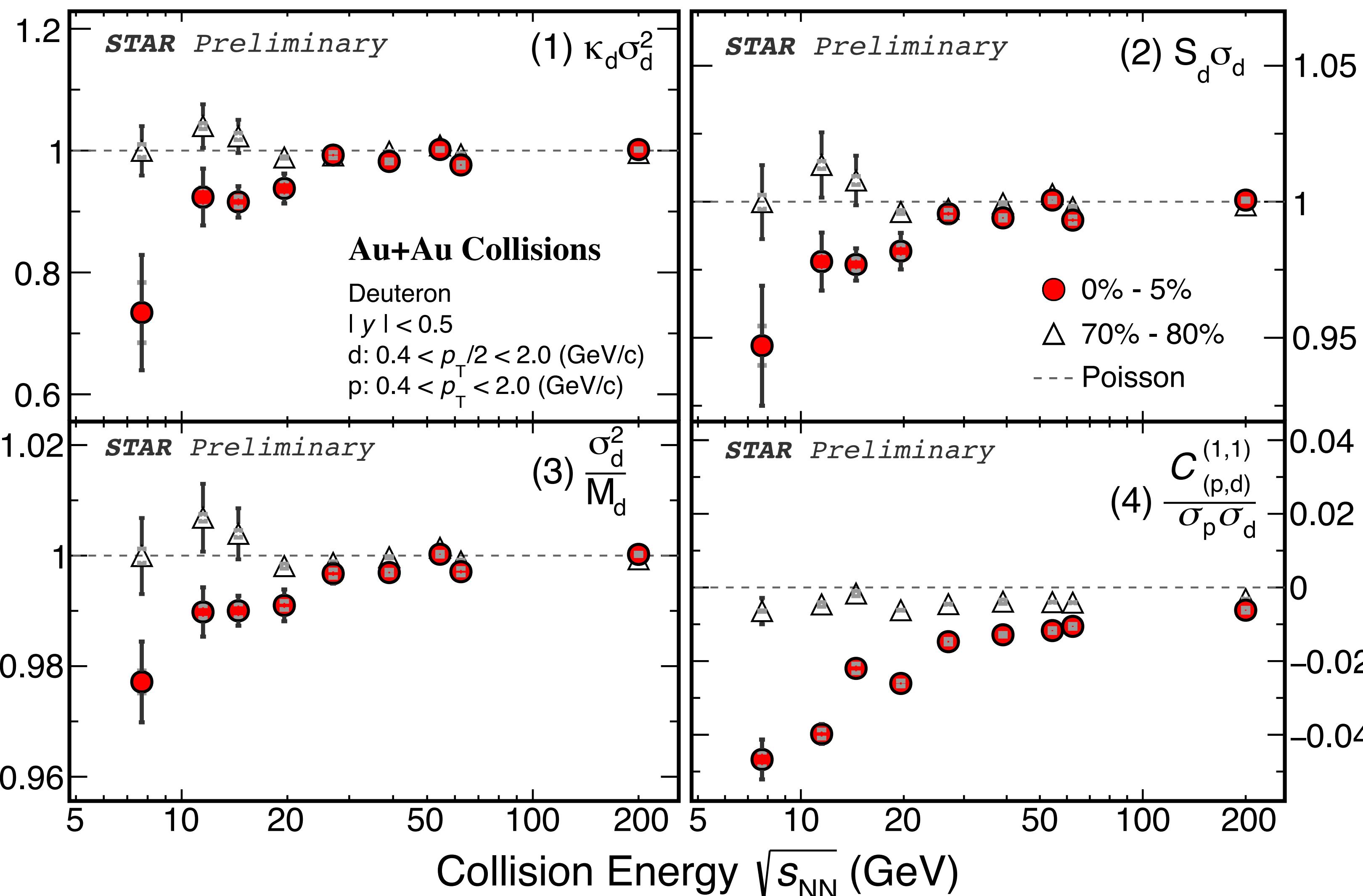
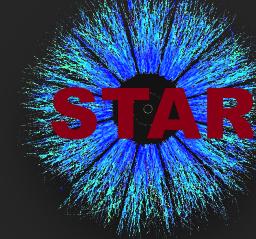


Cumulants ( $C_n$ ) of the deuteron distributions.

For peripheral (70%-80%) Au+Au collisions, cumulants are close to zero.

In most central (0-5%) collisions, cumulants increase as the collision  $\sqrt{s_{NN}}$  decreases.

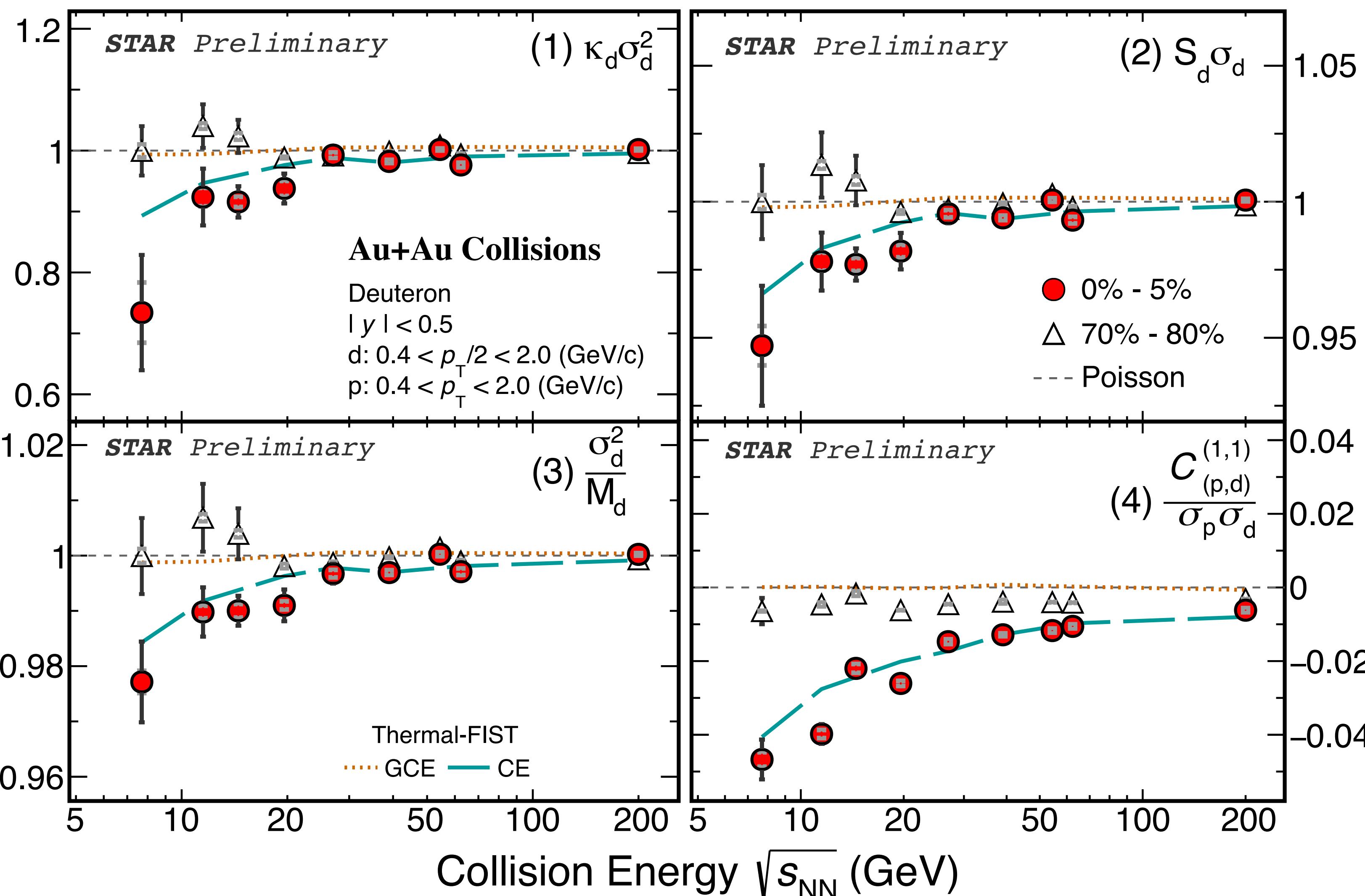
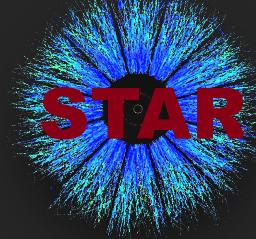
# Cumulant Ratios and p-d Correlation



- Cumulant ratios in 0-5% centrality, show monotonic dependence on  $\sqrt{s_{NN}}$ .
- Ratios in 70-80% centrality show weak  $\sqrt{s_{NN}}$  dependence and are close to 1.
- In panel(4), negative value of correlation suggests, proton and deuteron number are anti-correlated across all collision energy and centrality.
- With lowering the  $\sqrt{s_{NN}}$ , anti-correlation becomes stronger.

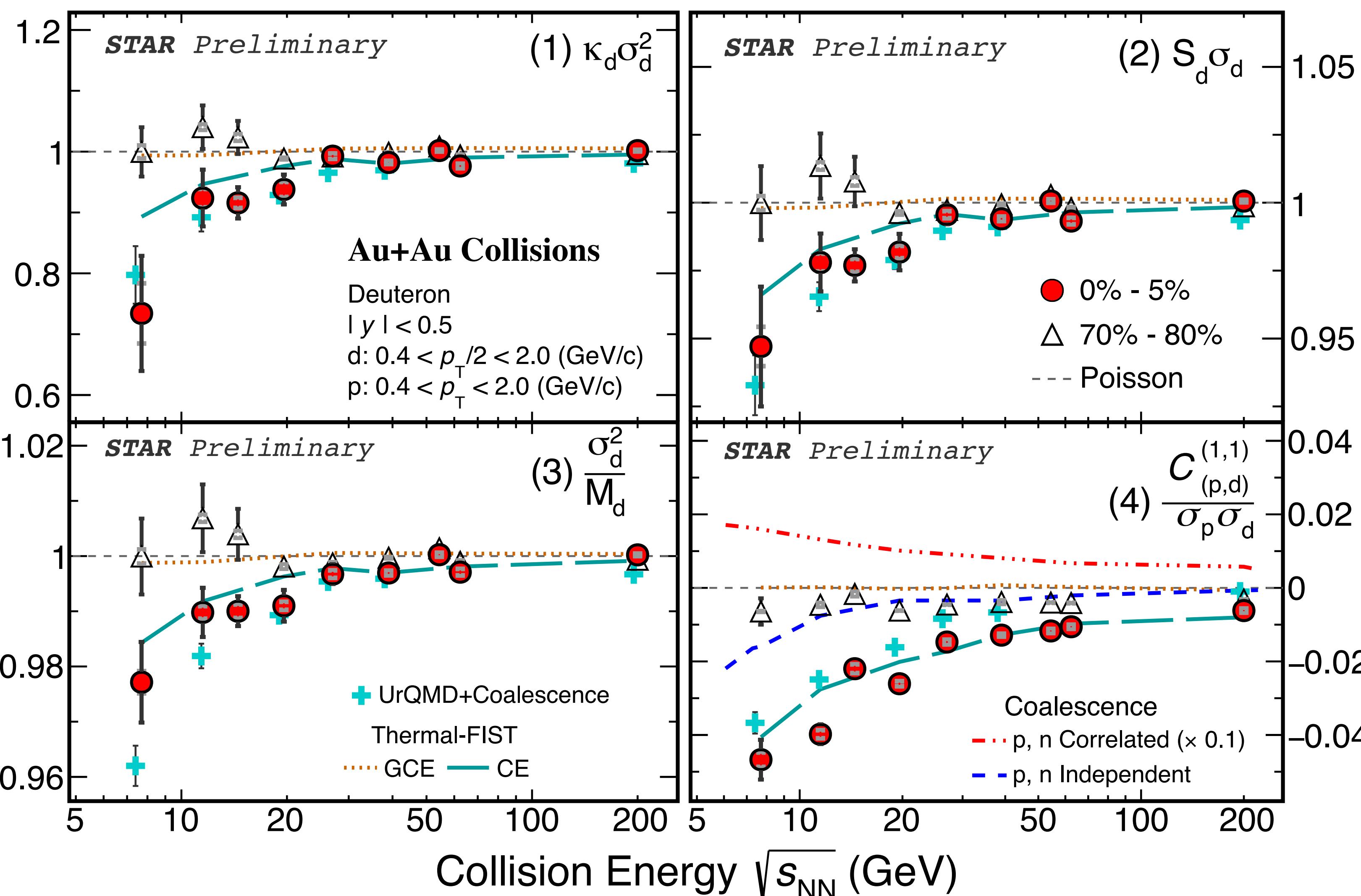
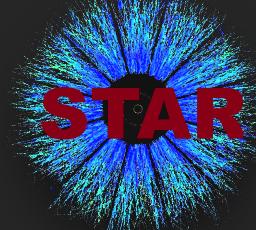
Black bars: statistical uncertainties \* Model-A: Correlated Proton and Neutron numbers ( $n_p = n_n$ )  
Grey caps: systematic uncertainties Model-B: Independent Proton and Neutron numbers.

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- GCE thermal model seems to fail to describe the cumulant ratios for lower  $\sqrt{s_{NN}}$ .

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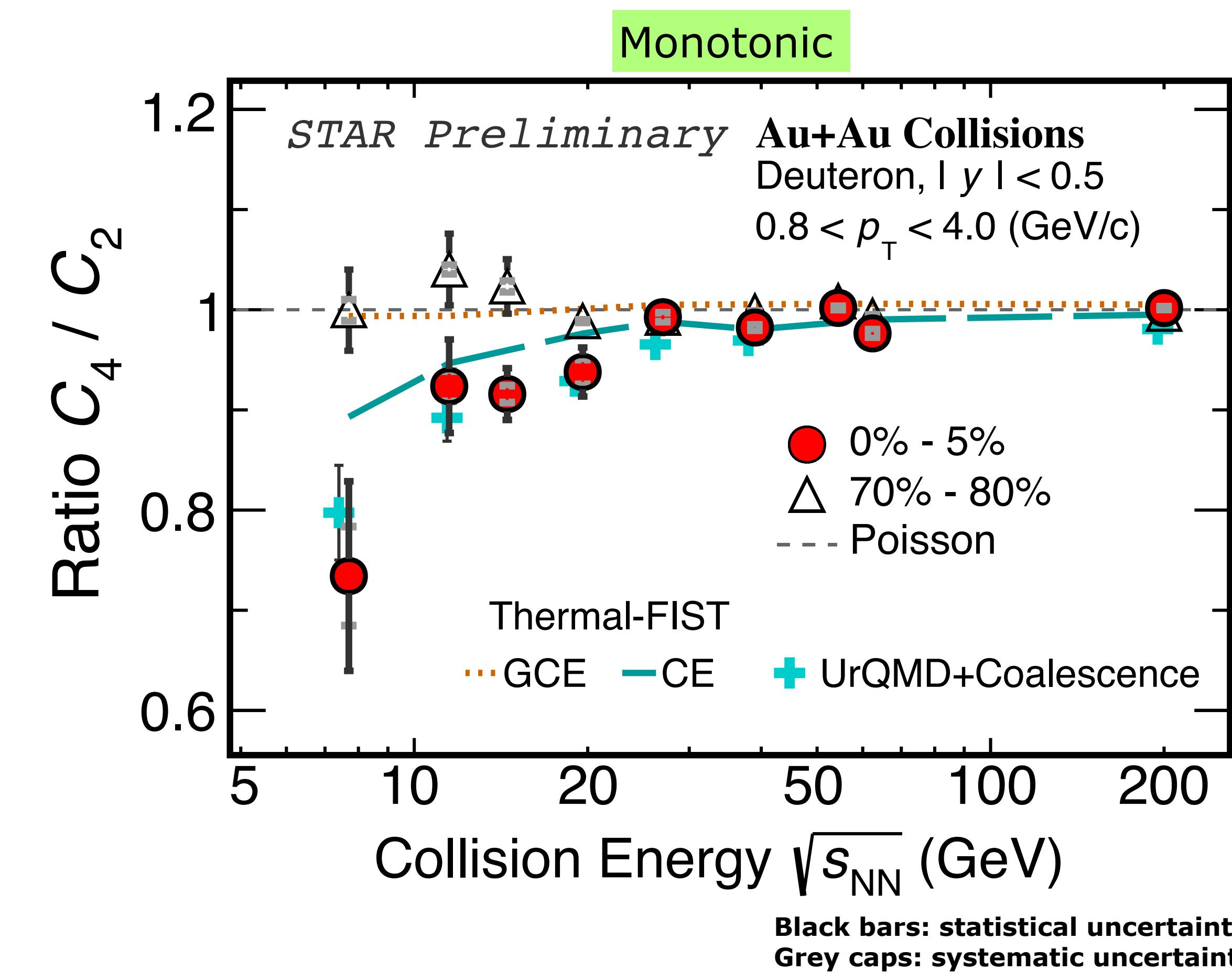
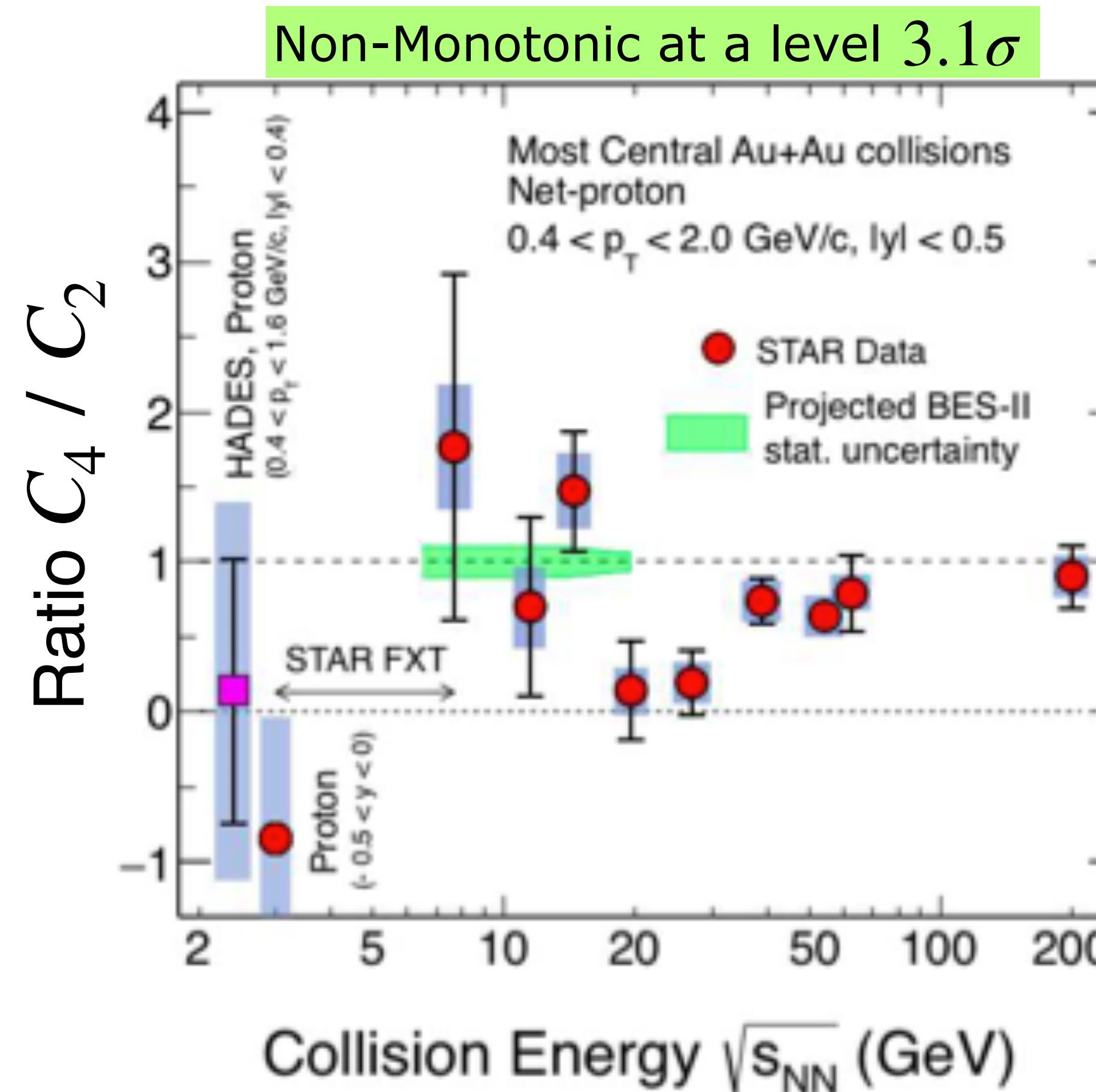
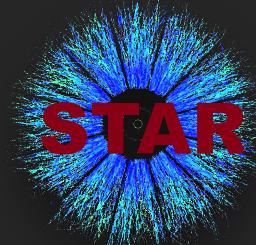


Black bars: statistical uncertainties  
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**Model-A: Correlated Proton and Neutron numbers ( $n_p = n_n$ )**  
**Model-B: Independent Proton and Neutron numbers.**

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- In panel(4), negative value of correlation suggests, proton and deuteron number are anti-correlated across all collision energy and centrality.
- With lowering the  $\sqrt{s_{NN}}$ , anti-correlation becomes stronger.
- GCE thermal model seems to fail to describe the cumulant ratios for lower  $\sqrt{s_{NN}}$ .
- UrQMD+Coalescence model qualitatively reproduces collision energy dependence.
- Neither correlated nor independent assumption for proton and neutron in the toy model from [Z. Fecková et. al., Phys. Rev. C 93, 054906 \(2016\)](#) reproduce the data.

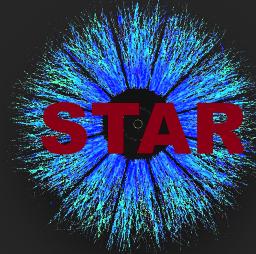
# Comparison with Net-proton



Deuteron number  $\kappa\sigma^2$  in 0-5% centrality show monotonic energy dependence in contrast to protons.

## Possibilities:

- **Low yield** of deuteron affecting sensitivity to critical point physics ?
- Probing **different freeze-out** surfaces ? More Investigation ongoing. Theoretical inputs are also needed.



## Summary:

- We reported the first measurements of cumulants of deuteron distribution, their ratios and proton-deuteron correlation in 0-5% and 70-80% central Au+Au collisions for  $\sqrt{s_{NN}} = 7.7 - 200$  GeV.
- UrQMD + phase-space coalescence model fairly describes the cumulant ratios and correlation for 0-5% centrality.
- For all  $\sqrt{s_{NN}}$ , proton and deuteron numbers are anti-correlated. With lowering  $\sqrt{s_{NN}}$ , anti-correlation in 0-5% centrality becomes stronger.
- HRG GCE thermal model fails to describe cumulant ratios at collision energies  $\sqrt{s_{NN}} \leq 19.6$  GeV. HRG CE and UrQMD show suppression below unity for lower  $\sqrt{s_{NN}}$ , as seen in the data, could arise from the effect of global baryon number conservation. Sensitive to the choice of ensembles.
- Kurtosis  $\times$  Variance of deuteron number in 0-5% centrality shows monotonic energy dependence in contrast to proton fluctuations. [STAR: Phys. Rev.Lett. 126 \(2021\) 092301](#)

## Outlook:

### Using BES-II data,

- Study contribution of  $p$ ,  $d$ ,  $t$ ,  $He^3$  etc. together to understand net-baryon fluctuation in low  $\sqrt{s_{NN}}$ .
- Understand the production mechanism and freeze-out properties of light nuclei.