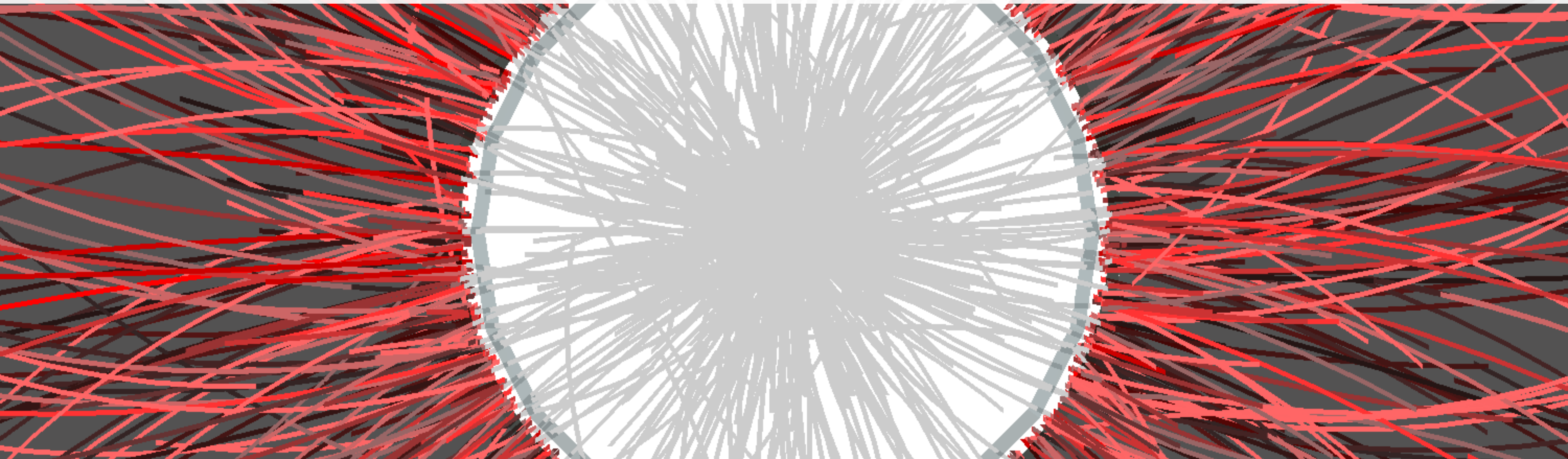


# ALICE determines the scattering parameters of D mesons with light-flavour hadrons



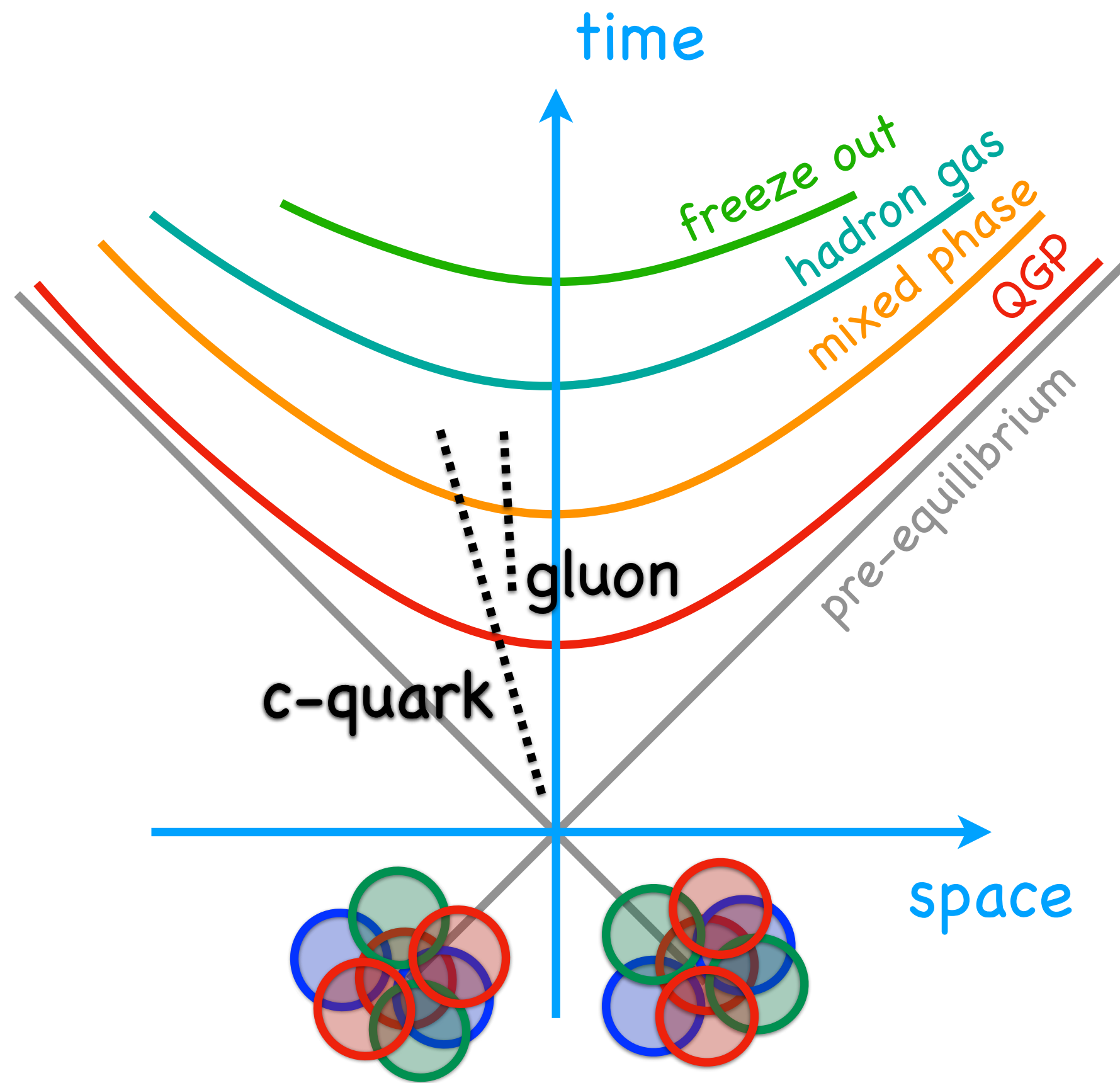
ALICE



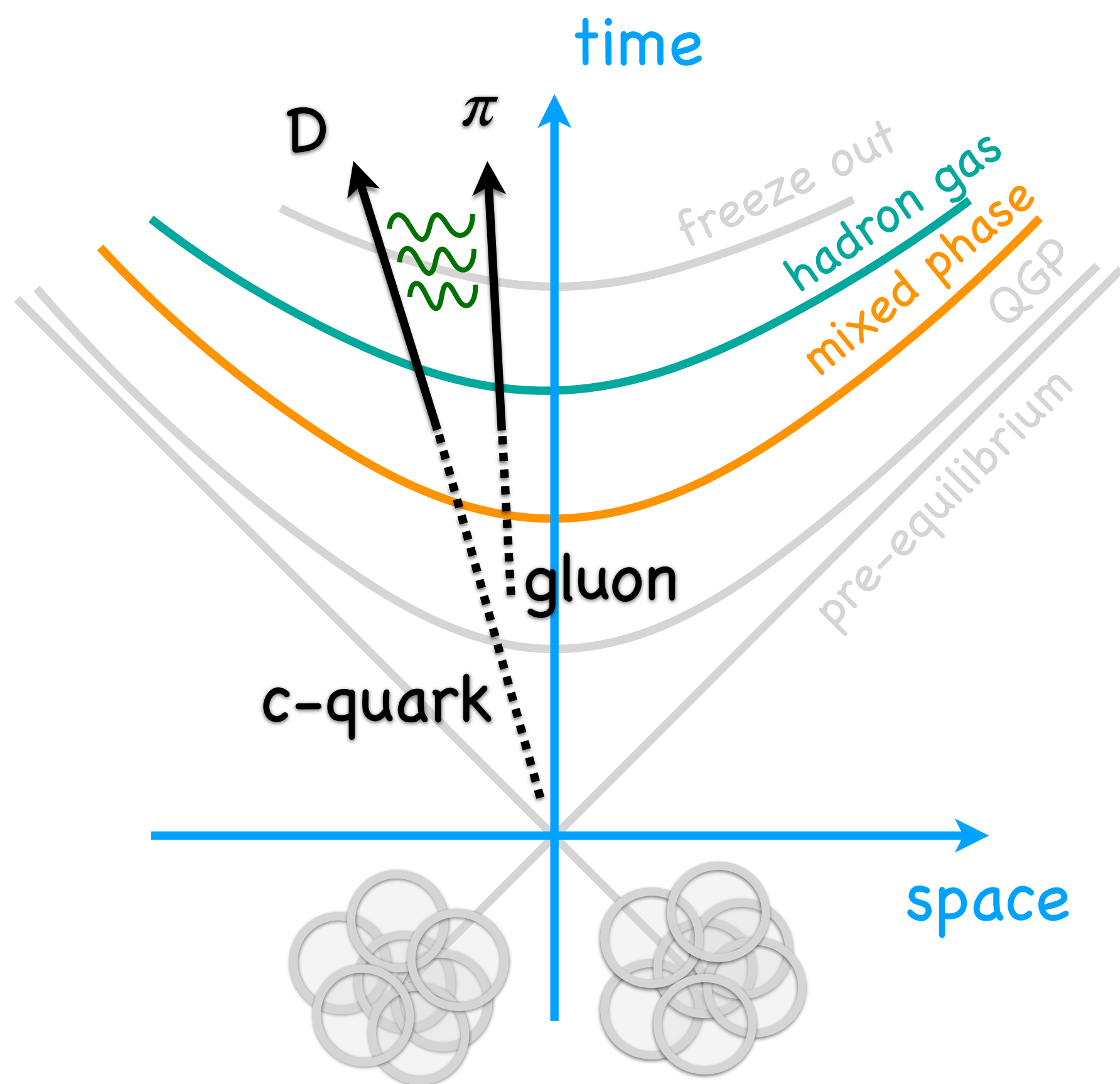
Fabrizio Grosa on behalf of the *ALICE Collaboration*

CERN





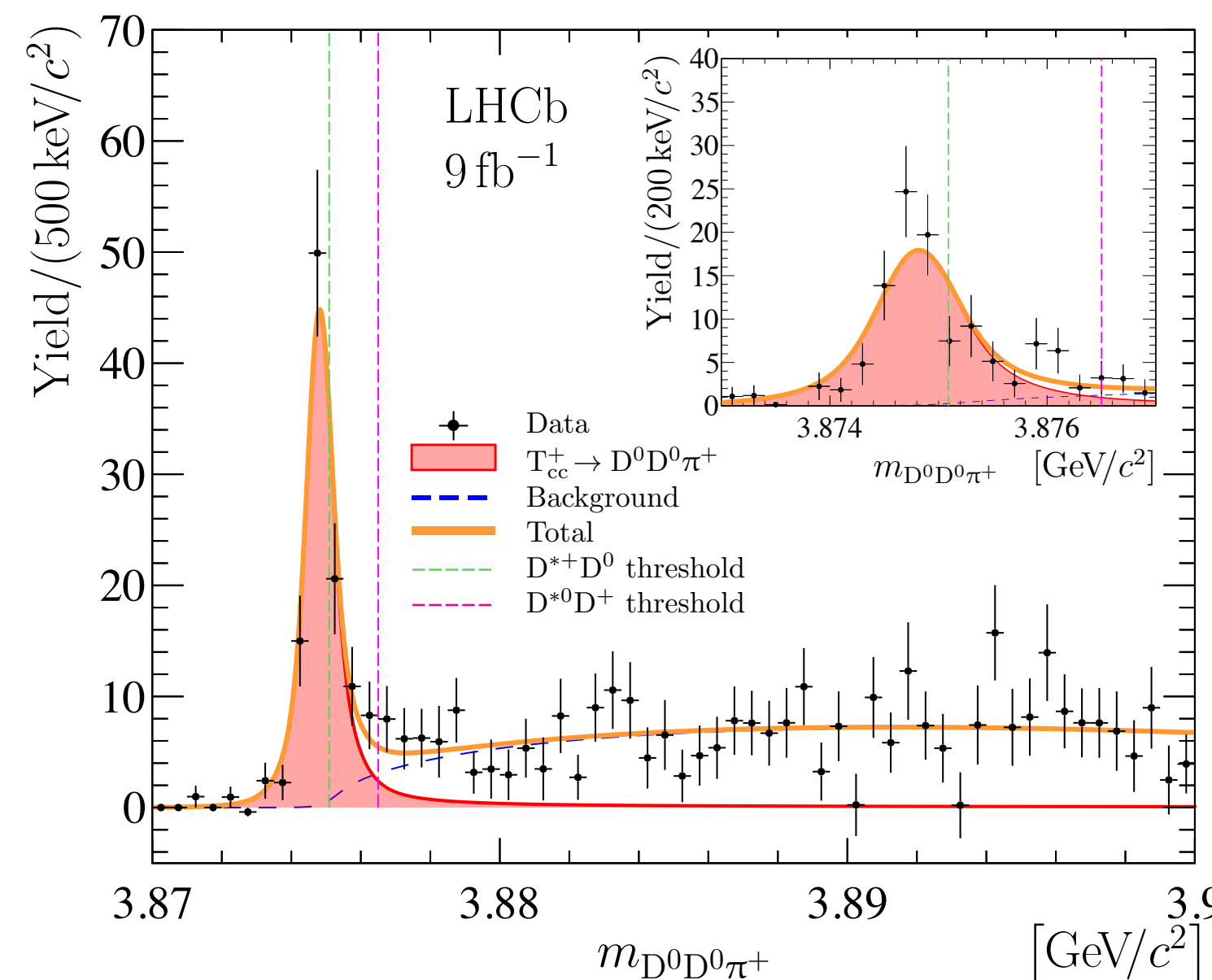
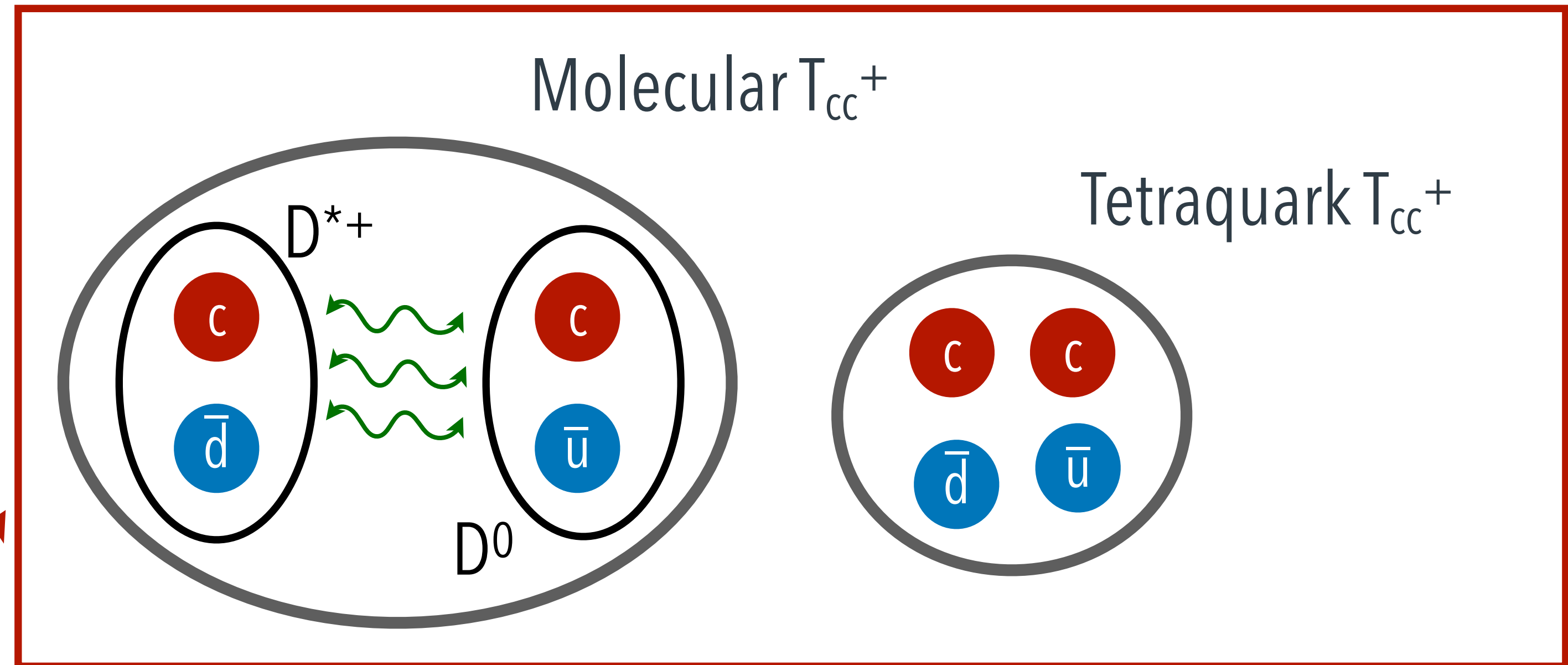
- Charm quarks: produced via hard scattering processes before the formation of the quark–gluon plasma (QGP), subsequently interact with the medium constituents  
→ Ideal probes of the QGP



- Charm quarks: produced via hard scattering processes before the formation of the quark–gluon plasma (QGP), subsequently interact with the medium constituents
  - Ideal probes of the QGP
- After the hadronisation, charm hadrons might still interact with the light hadrons produced
  - How much hadronic rescatterings influence our observables?
- In the TAMU model the scattering lengths used for  $\pi D$  and  $\bar{K} D$  are:
  - $a_{\pi D}(l=3/2) = -0.10$  fm
  - $a_{\bar{K} D}(l=1) = -0.22$  fm
  - No experimental constraints

● Charm molecules?

System	$I (J^{P(C)})$	Candidate
np	$0 (1^+)$	deuteron
ND	$0 (1/2^-)$	$\Lambda_c(2765)$
ND*	$0 (3/2^-)$	$\Lambda_c(2940)$
ND	$0 (1/2^-)$	$\Sigma_c(2800)$
$D^*\bar{D}$	$0 (1^{++})$	X(3872)
$D^*D$	$0 (1^+)$	$T_{cc}$
$D_1\bar{D}$	$0 (1^{--})$	Y(4260)
$D_1\bar{D}^*$	$0 (1^{--})$	Y(4360)
$\Sigma\bar{D}$	$1/2 (1/2^-)$	$P_c(4312)$
$\Sigma\bar{D}^*$	$1/2 (1/2^-)$	$P_c(4457)$
$\Sigma\bar{D}^*$	$1/2 (3/2^-)$	$P_c(4440)$



- Just below  $DD^*$  threshold
- ➔ ideal candidate to be a molecular state

- Femtoscscopy technique: based on the *correlation function (CF)*

$$C(\vec{k}^*) = \underbrace{\mathcal{N} \frac{N_{\text{same}}^{\text{pairs}}(k^*)}{N_{\text{mixed}}^{\text{pairs}}(k^*)}}_{\text{Experiment}} = \underbrace{\int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*}_{\text{Theory}}$$

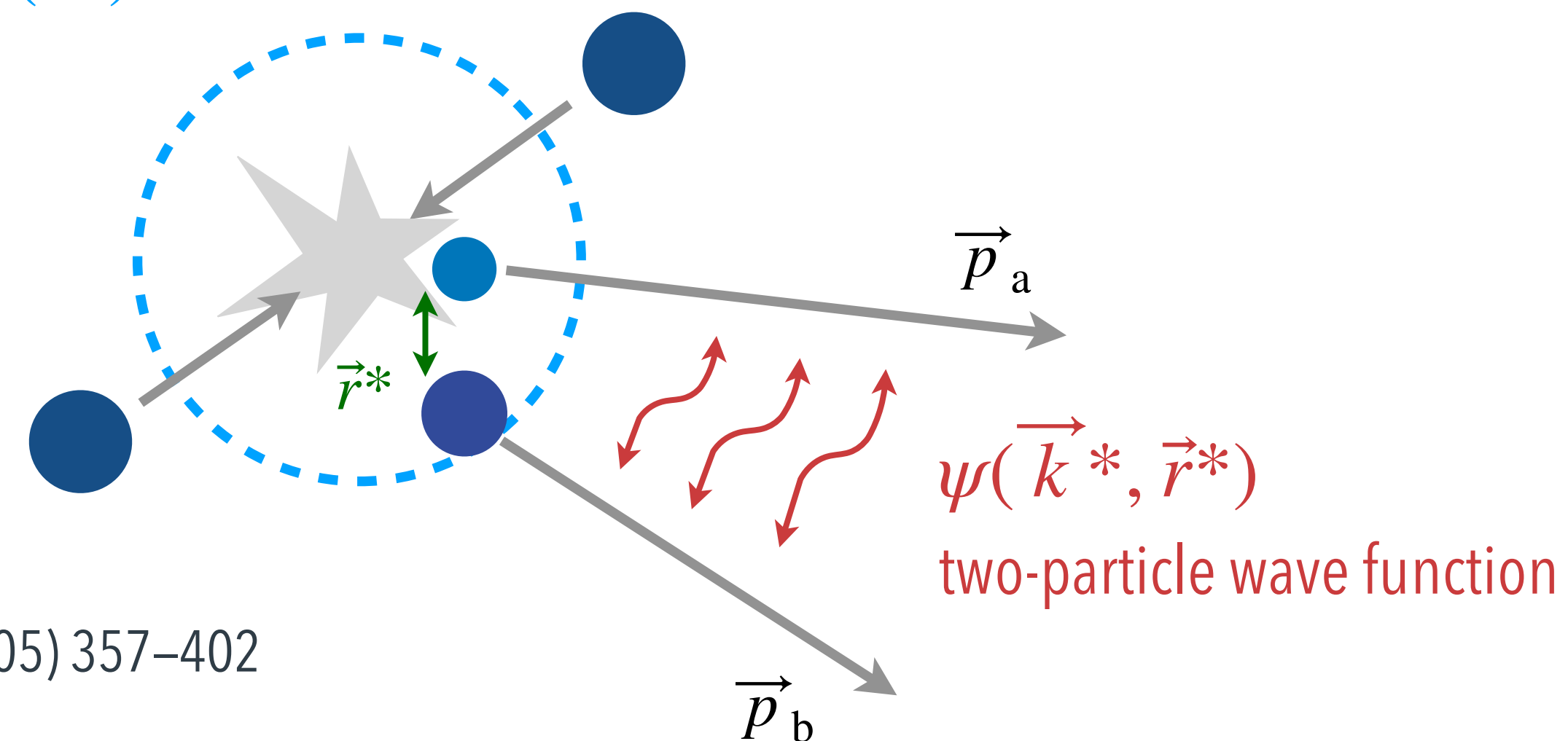
Koonin-Pratt equation

 M.Lisa, S. Pratt et al, Ann.Rev.Nucl.Part.Sci. 55 (2005) 357–402

where  $\vec{k}^* = \frac{\vec{p}_a^* - \vec{p}_b^*}{2}$  is in the rest frame of the particle pair

- **Relative wave function** sensitive to interaction potential
- **Emitting source**: hypersurface at kinematic freeze out of final-state particles
- CF sensitive to strong interaction when the source size  $\sim 1$  fm

$S(\vec{r}^*)$  source function



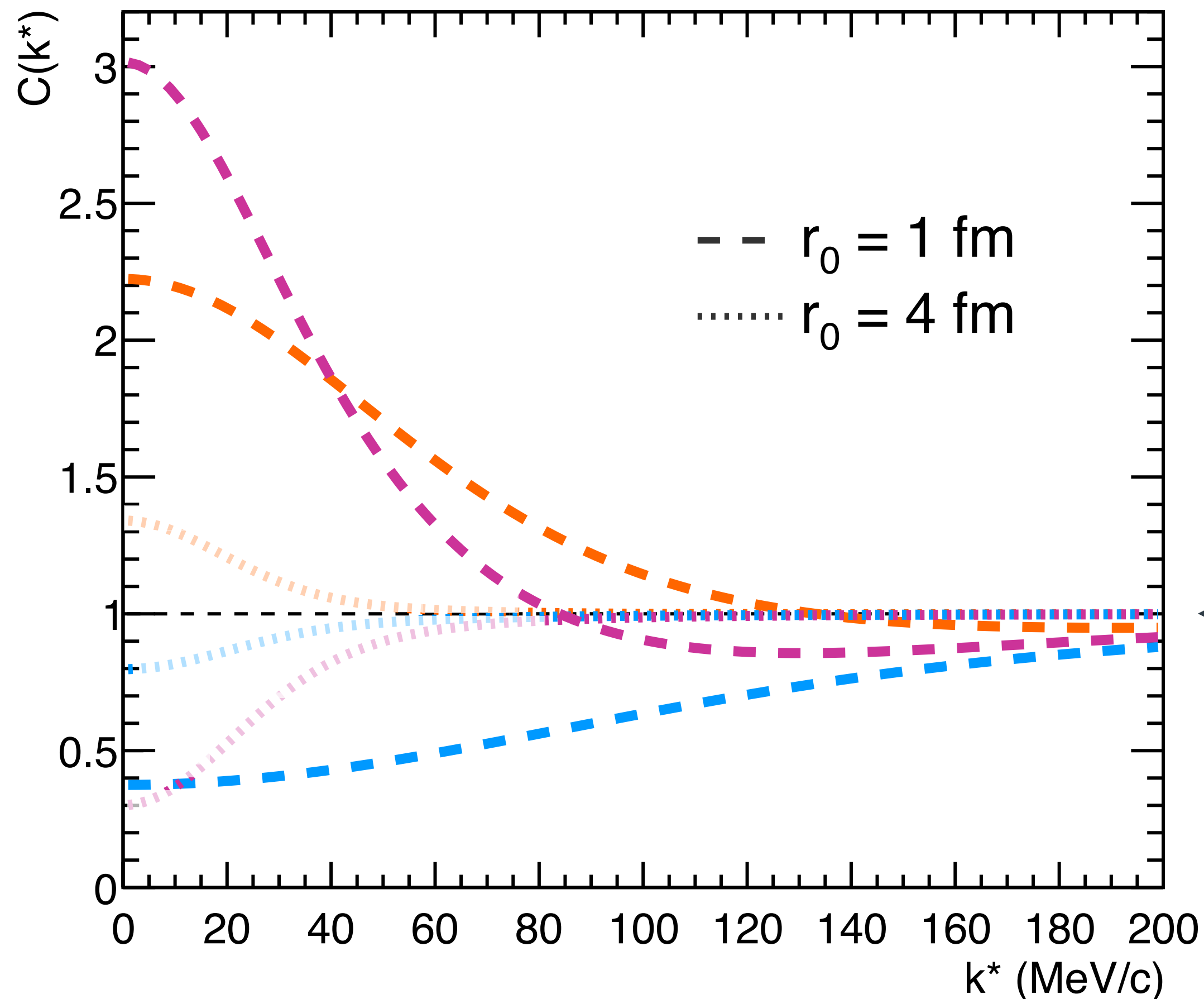
CF computed in ALICE using **CATS** (Correlation Analysis Tool using the Schrödinger equation)

- Developed at Technische Universität München
- Provides exact solution of Schrödinger equation for wave function

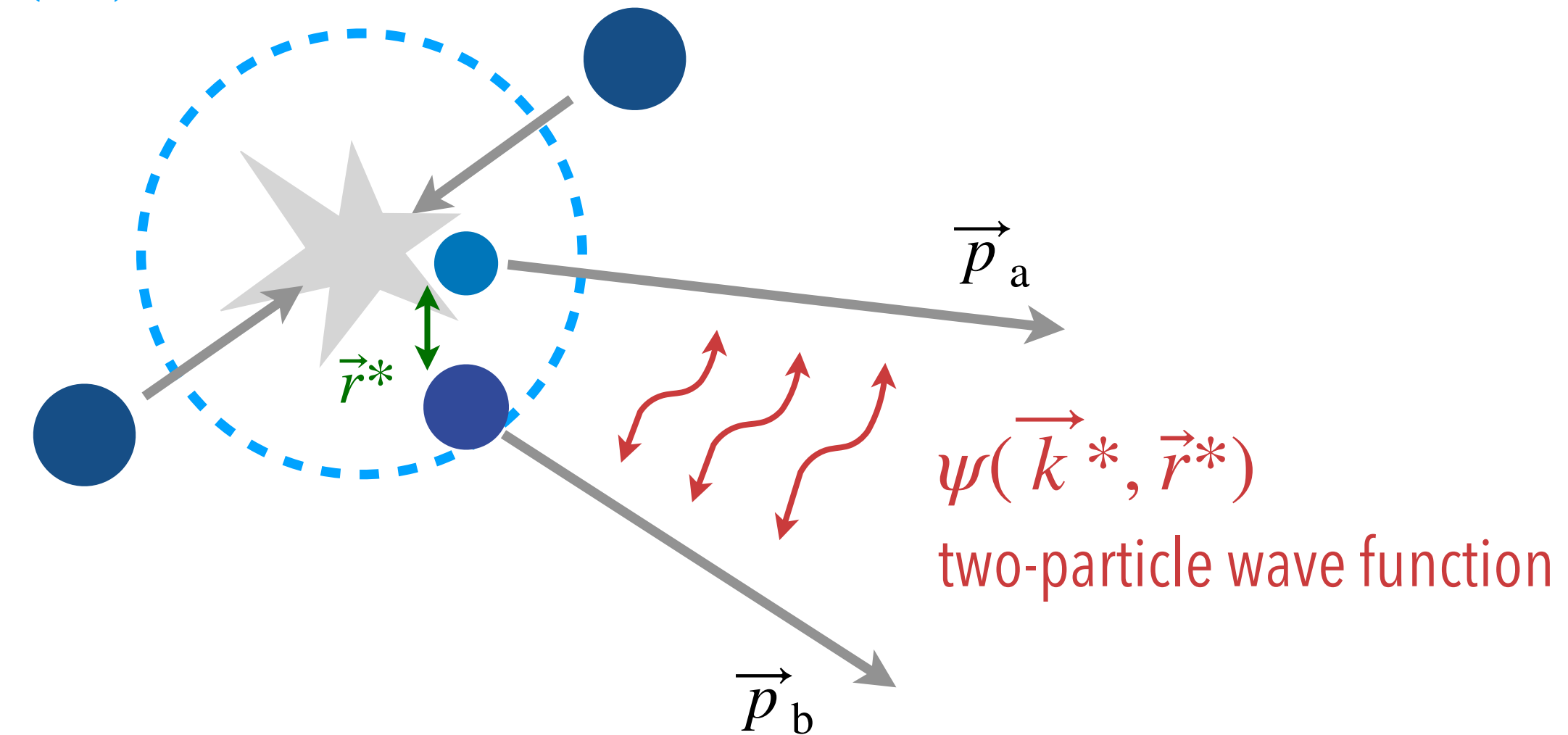
 D. L. Mihaylov et al,  
Eur. Phys. Journal C 78 (2018) 394

$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3r^*$$

→ Relative wave function sensitive to interaction potential



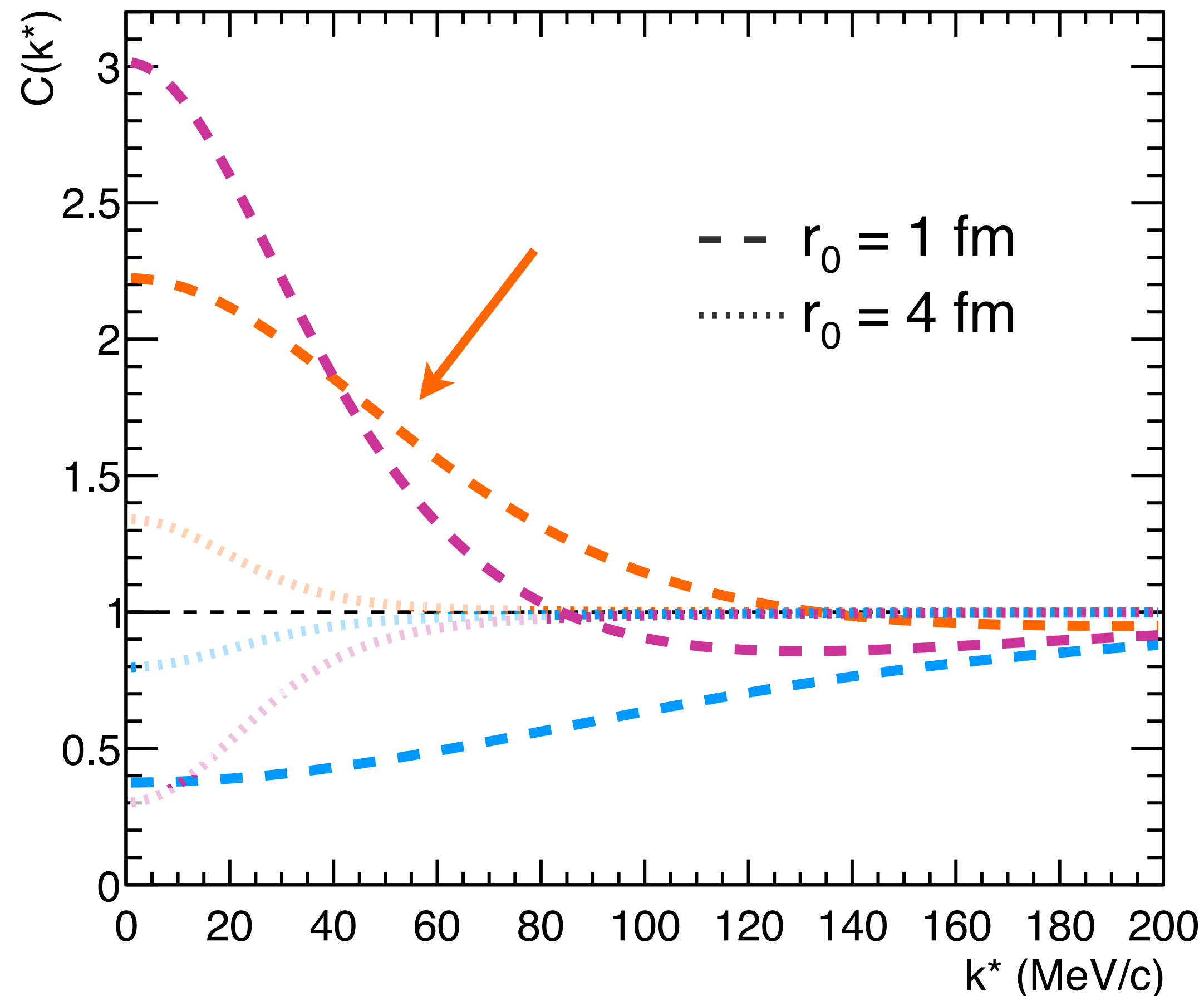
$S(\vec{r}^*)$  source function



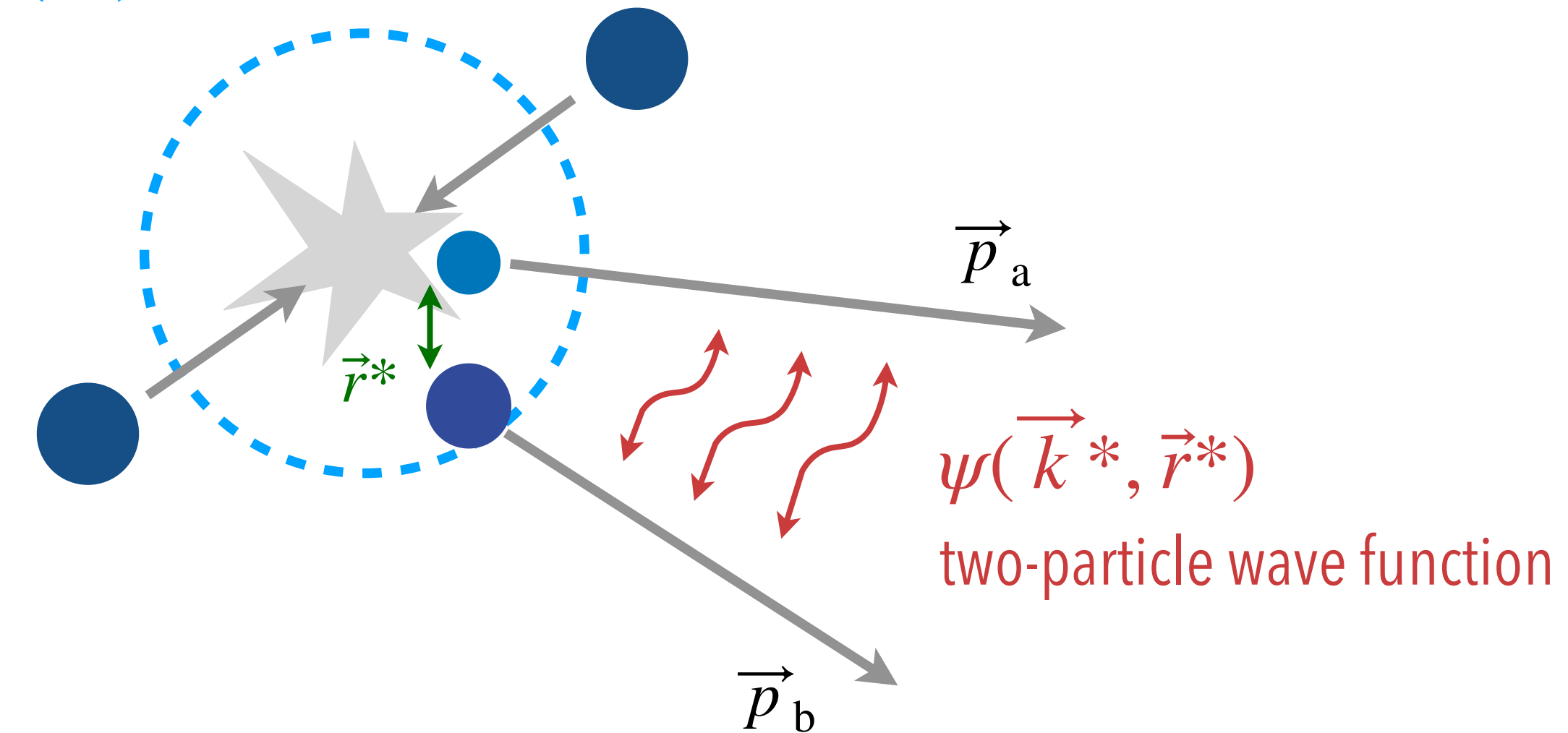
→ Absence of interaction  $C(k^*) = 1$

$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3r^*$$

→ Relative wave function sensitive to interaction potential



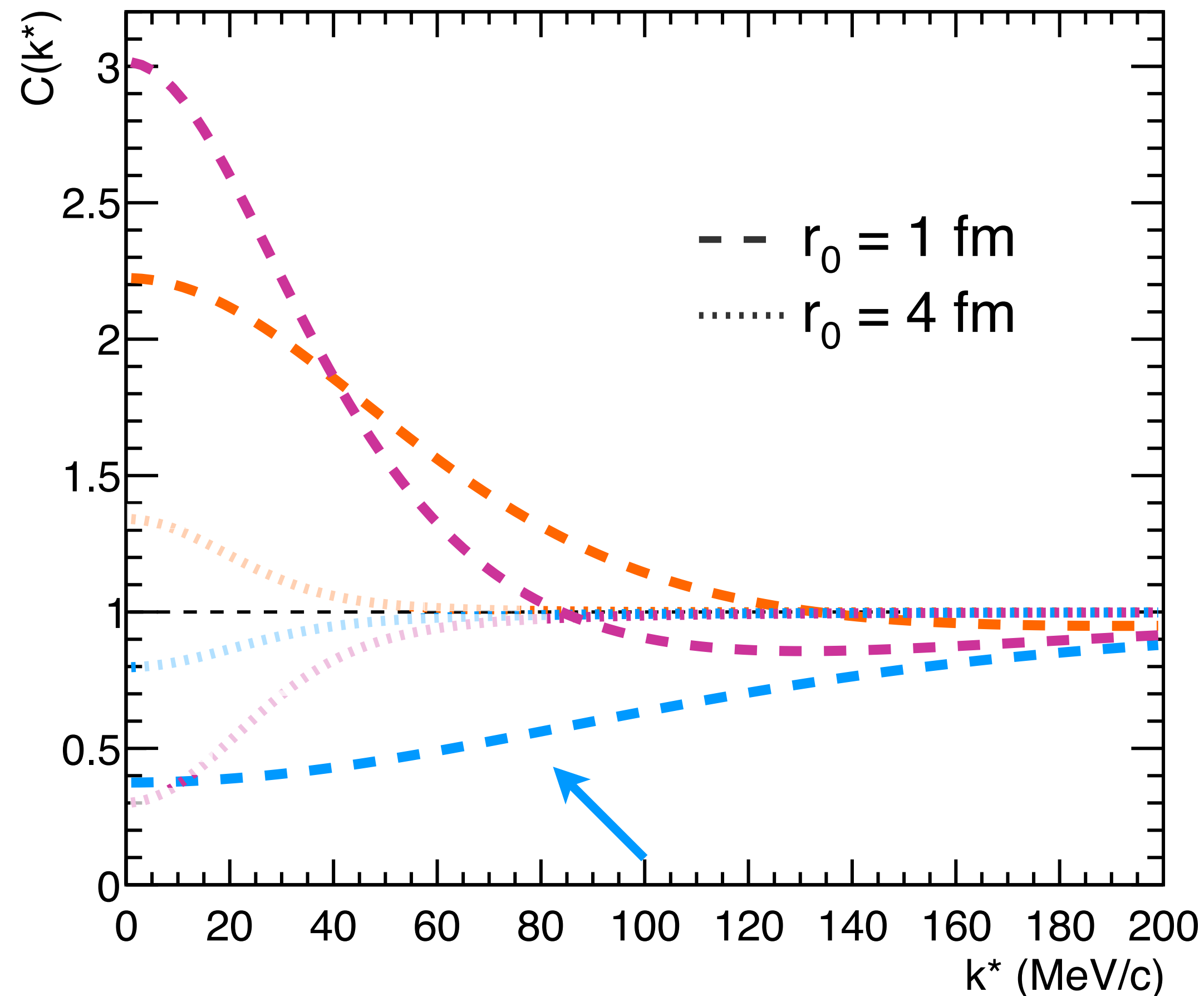
$S(\vec{r}^*)$  source function



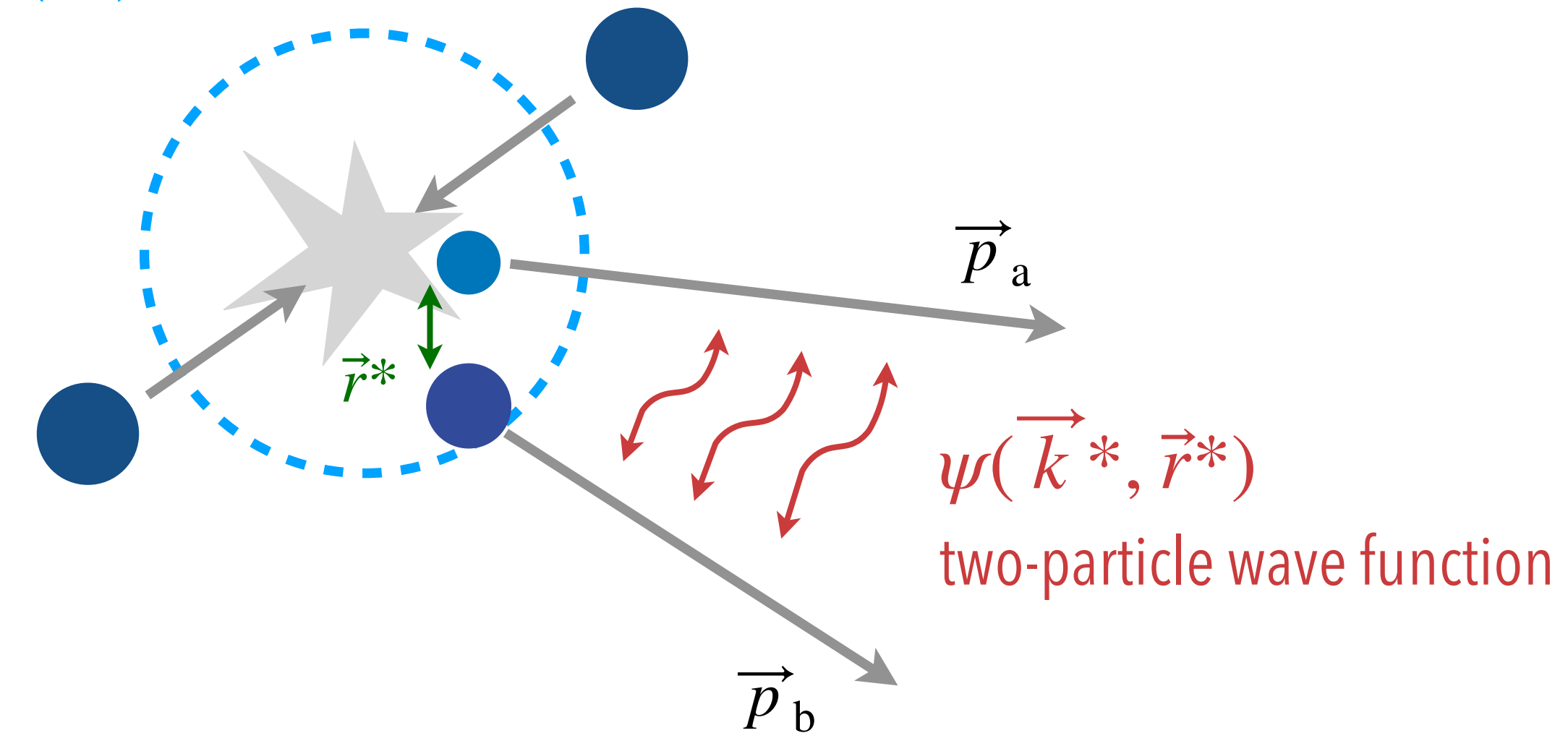
- Absence of interaction  $C(k^*) = 1$
- Attractive potential  $C(k^*) > 1$

$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3r^*$$

→ Relative wave function sensitive to interaction potential



$S(\vec{r}^*)$  source function

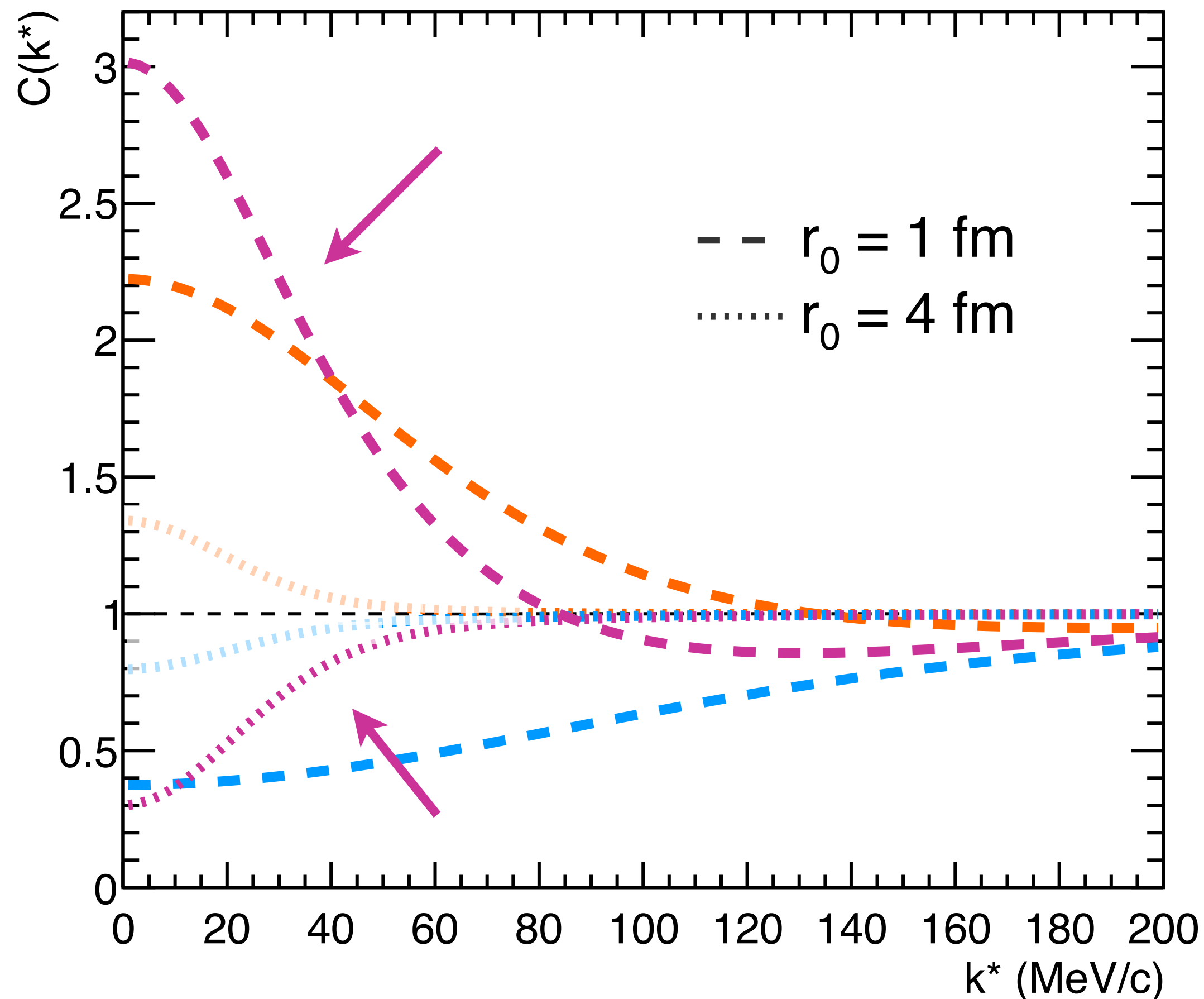


- Absence of interaction  $C(k^*) = 1$
- Attractive potential  $C(k^*) > 1$
- Repulsive potential  $C(k^*) < 1$

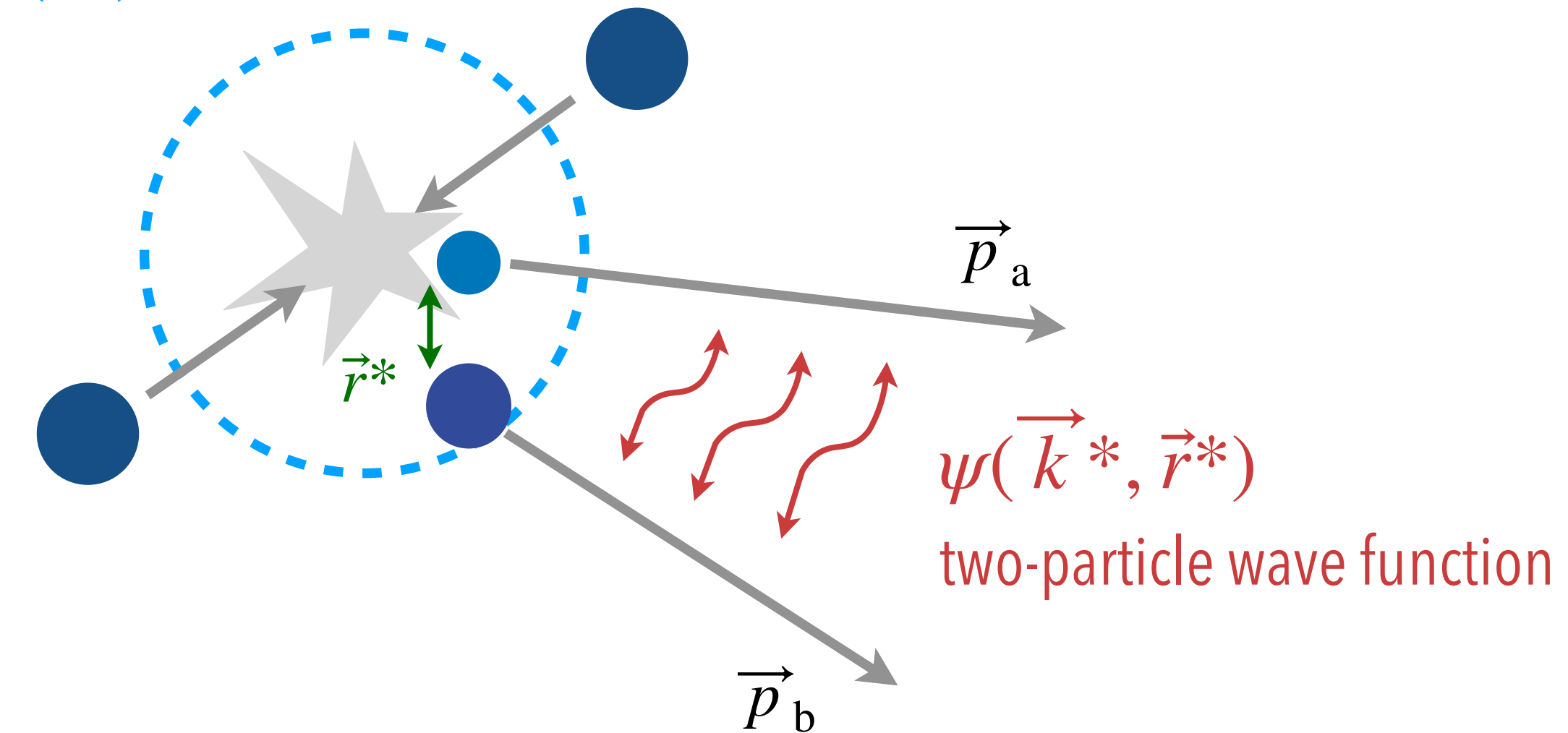


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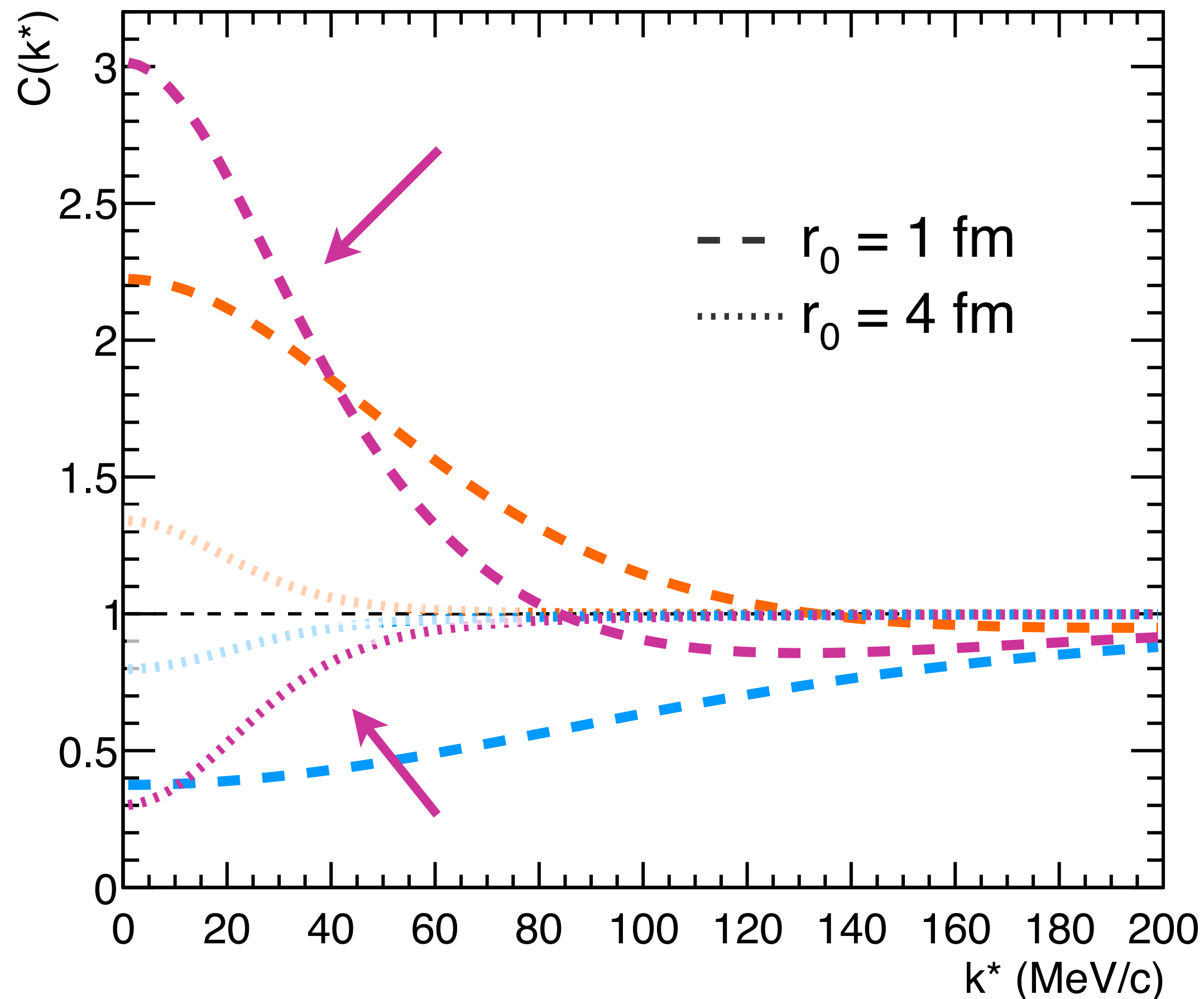
$S(\vec{r}^*)$  source function



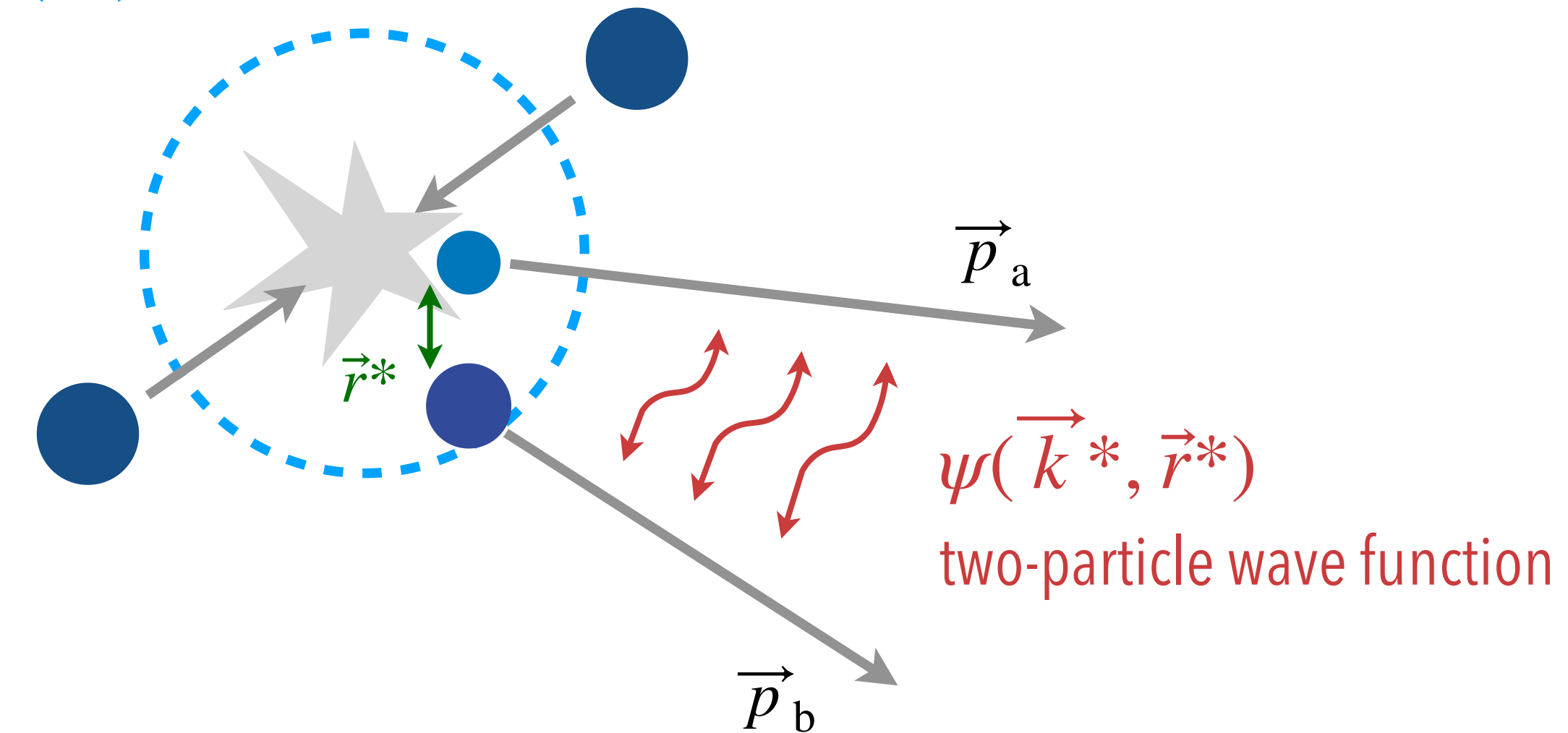
- Absence of interaction  $C(k^*) = 1$
- Attractive potential  $C(k^*) > 1$
- Repulsive potential  $C(k^*) < 1$
- Bound-state formation  $C(k^*) \ll 1$

$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3r^*$$

→ Relative wave function sensitive to interaction potential



$S(\vec{r}^*)$  source function



- Absence of interaction  $C(k^*) = 1$
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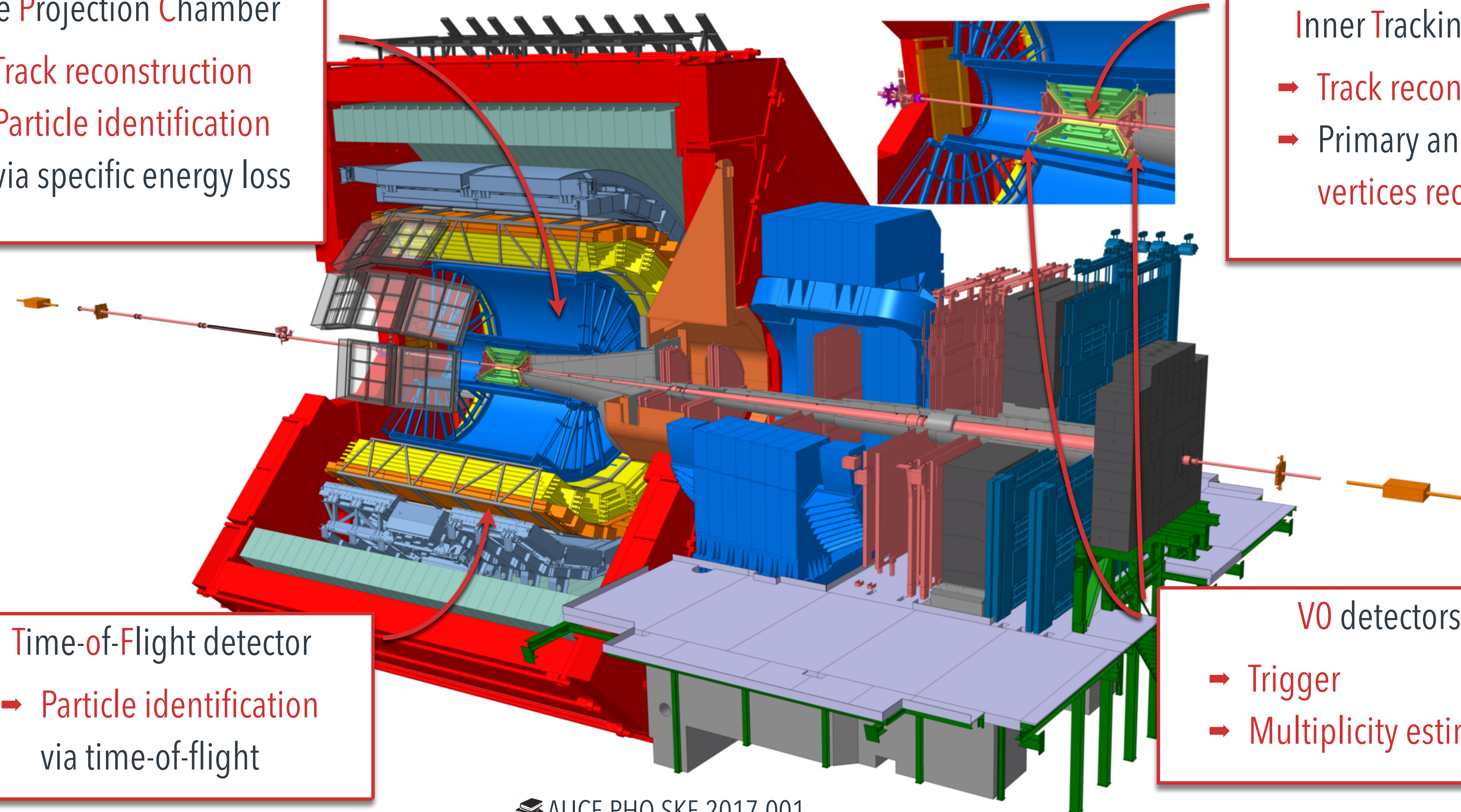
Strange-particle femto, V. Mantovani Sarti [talk](#) T08 06/04 10:00  
 Three-body interactions, R. Del Grande [talk](#) T08 06/04 15:00

## Time Projection Chamber

- Track reconstruction
- Particle identification via specific energy loss

## Inner Tracking System

- Track reconstruction
- Primary and decay vertices reconstruction



## Time-of-Flight detector

- Particle identification via time-of-flight

## V0 detectors

- Trigger
- Multiplicity estimation

## Time Projection Chamber

- Track reconstruction
- Particle identification via specific energy loss

## Inner Tracking System

- Track reconstruction
- Primary and decay vertices reconstruction

## Dataset

High-multiplicity pp at  
 $\sqrt{s} = 13 \text{ TeV} \sim 10^9$  events

## Time-of-Flight detector

- Particle identification via time-of-flight

## V0 detectors

- Trigger
- Multiplicity estimation

- Decay channel:



- BR =  $(9.38 \pm 0.16)\%$

- 📖 PDG, Prog. Theor. Exp. Phys. 2020 083C01

- Fully reconstruct decay topologies ( $c\tau \approx 312 \mu\text{m}$ )

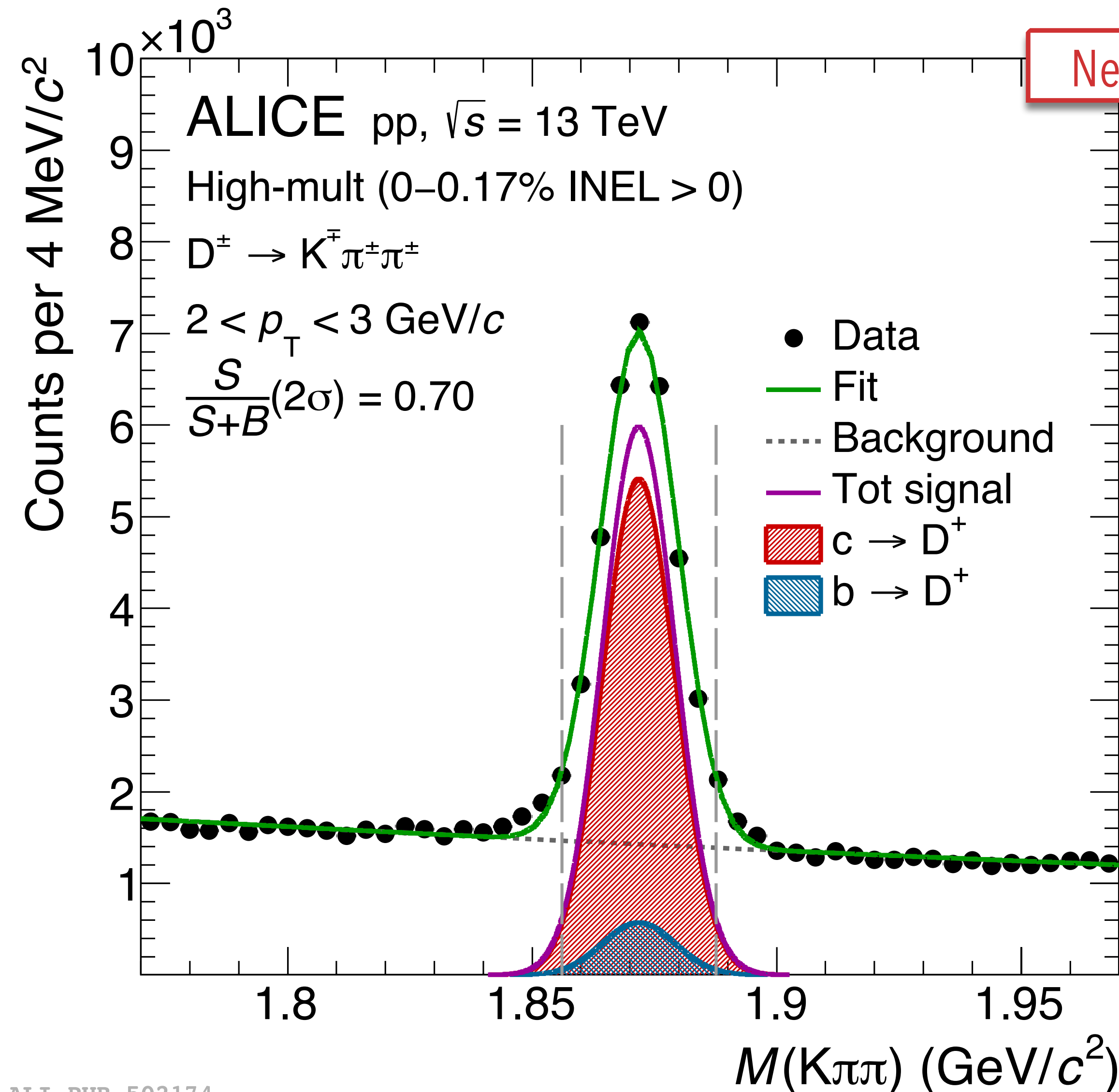
- Purity of about 70%, non-prompt contribution about 7%

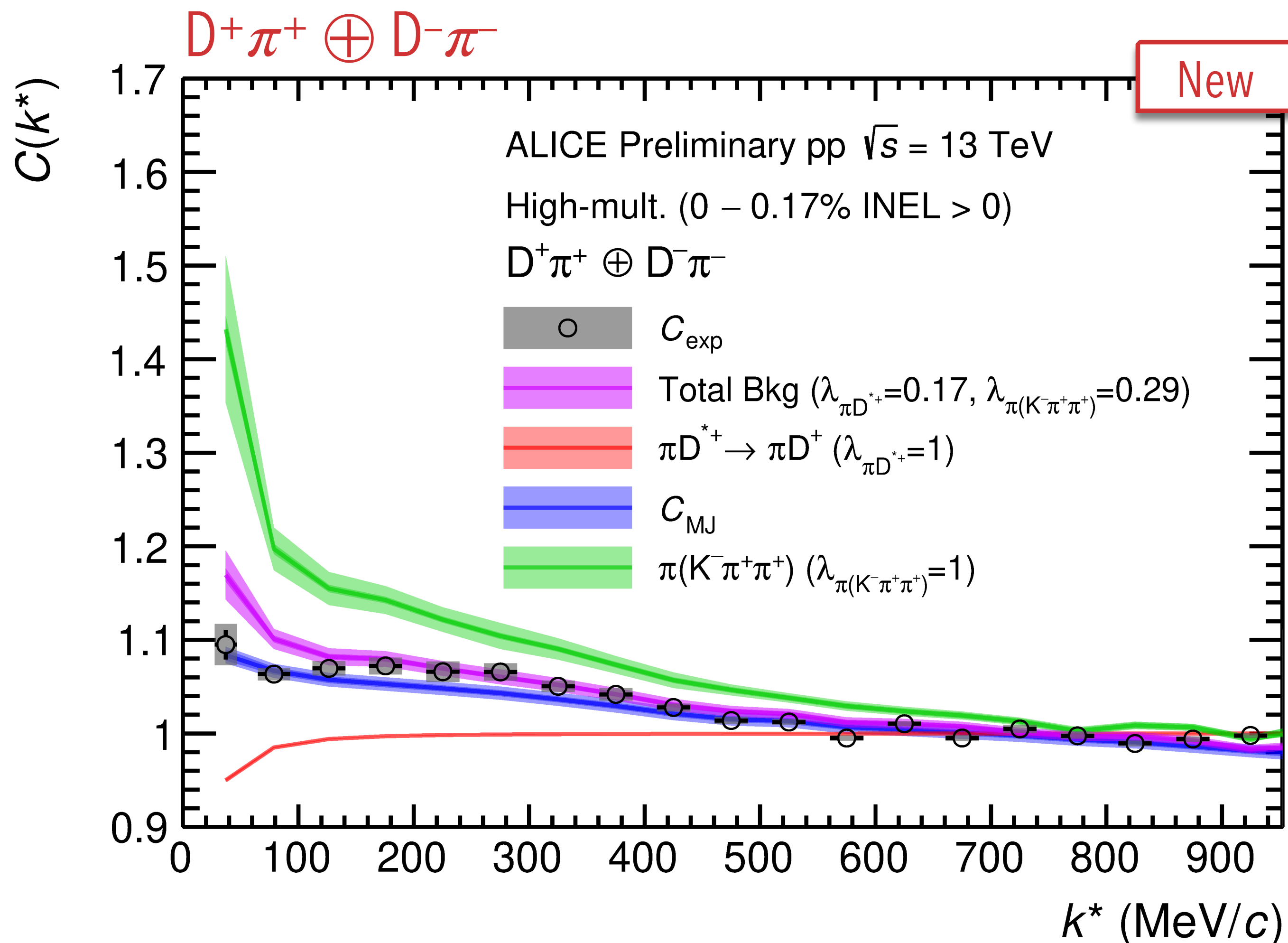
- Contributions:

- **Prompt:**

- ▶ from c hadronisation
    - ▶ from D\* decays

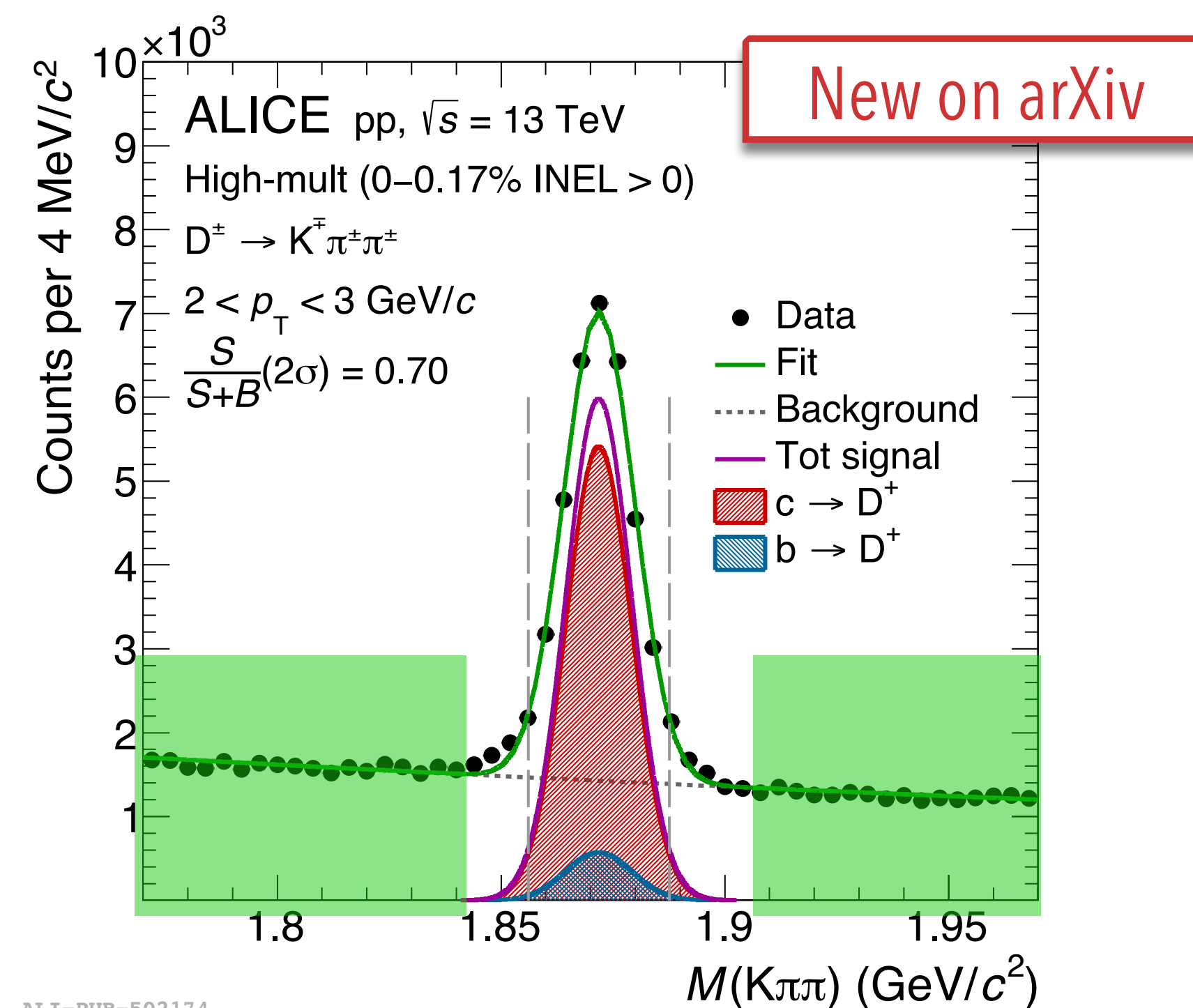
- **Non-prompt:** from b decays

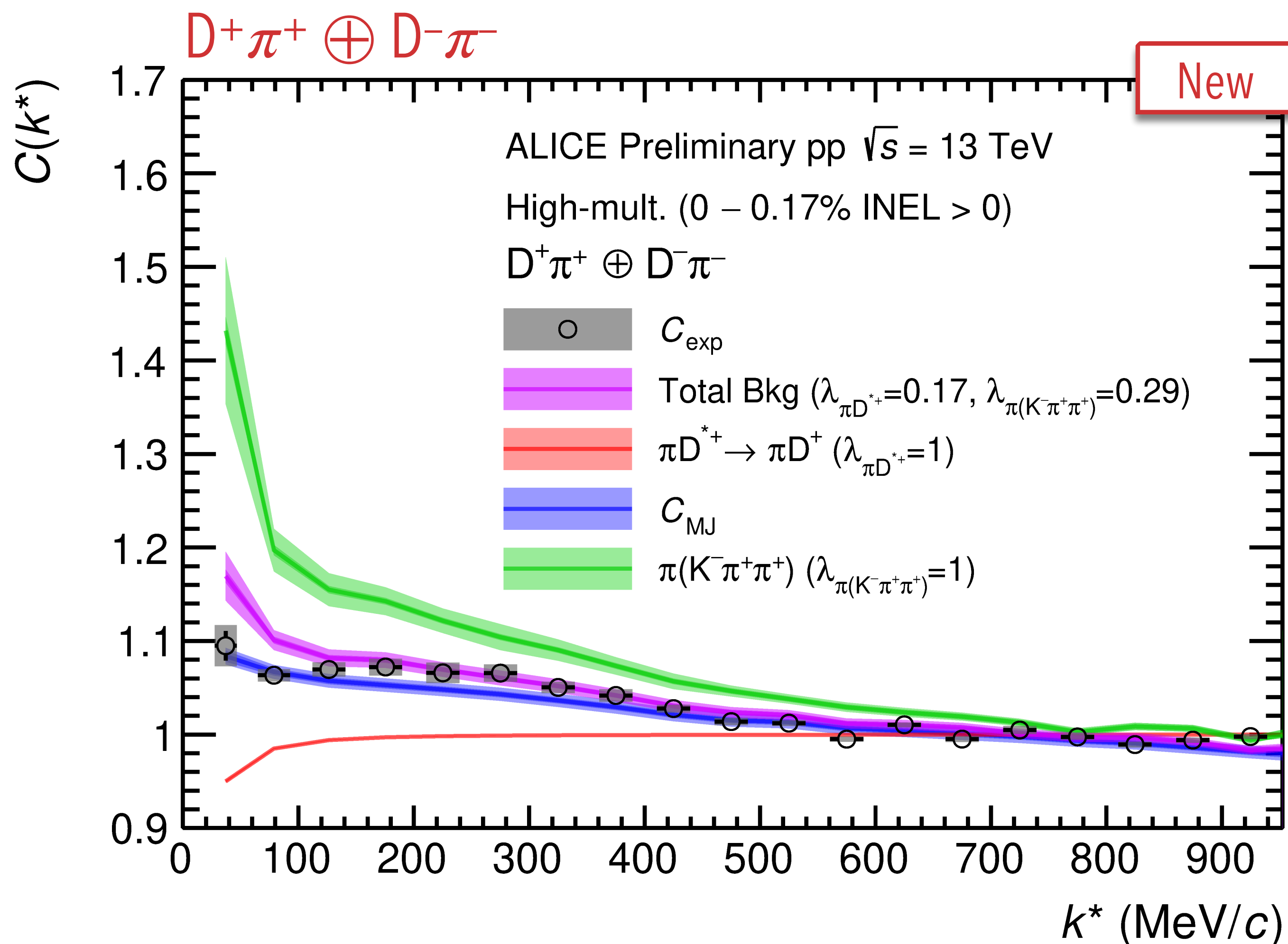




- Raw correlation function includes different sources of backgrounds

- Combinatorial background estimated from D-meson sidebands

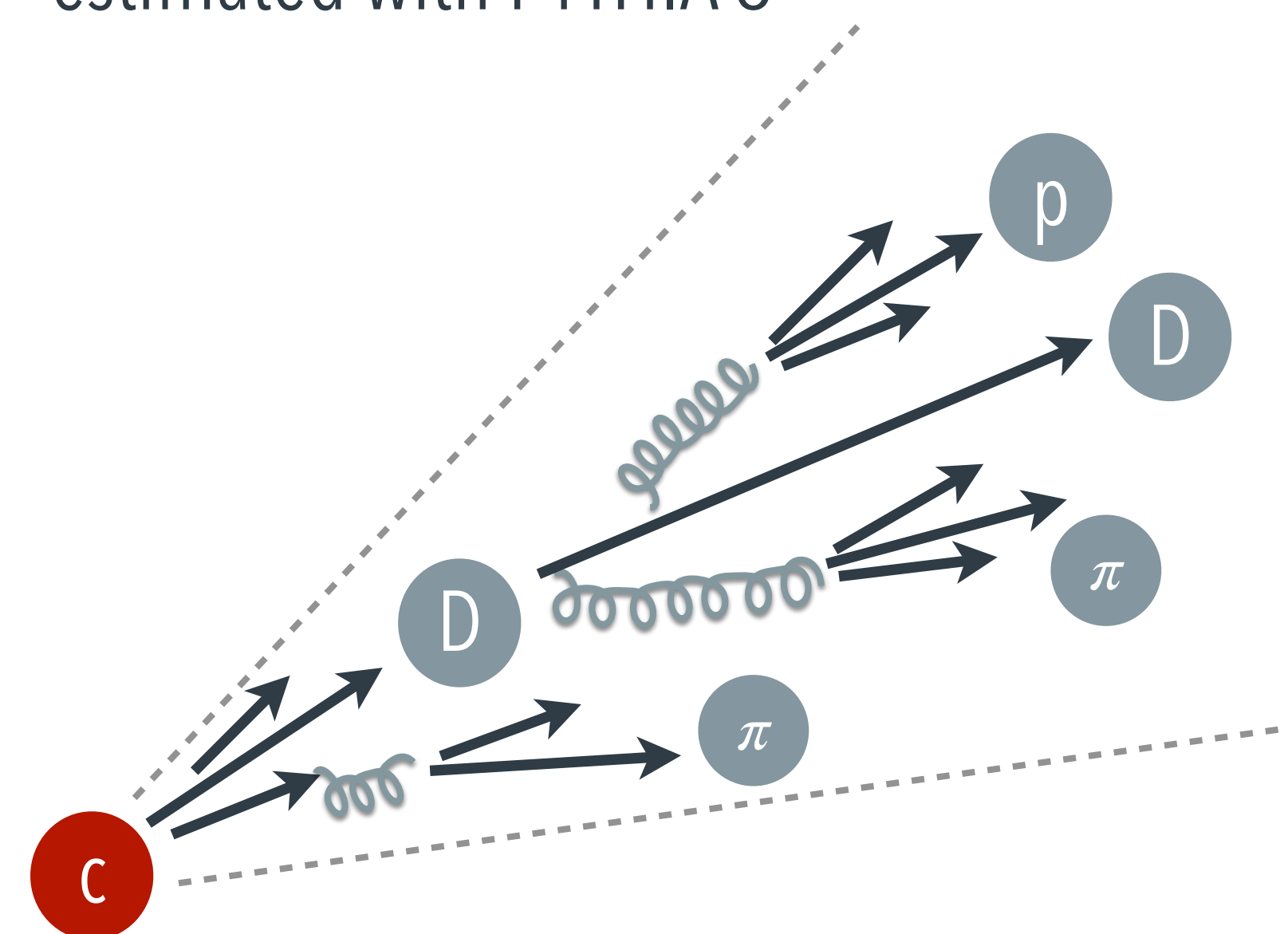


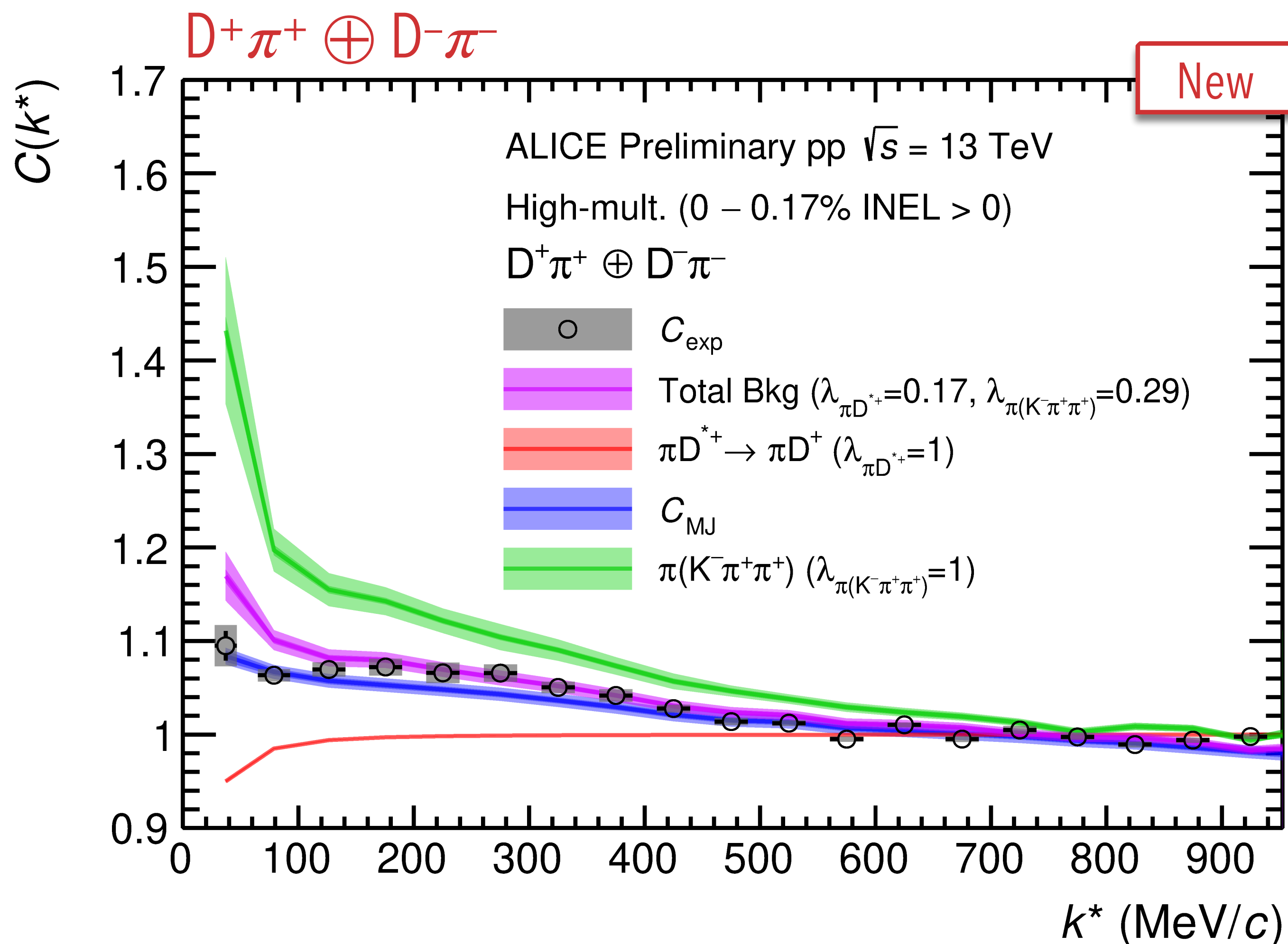


- Raw correlation function includes different sources of backgrounds

i. **Combinatorial background**  
estimated from D-meson sidebands

ii. **Jet-like correlations**  
estimated with PYTHIA 8





- Raw correlation function includes different sources of backgrounds

i. **Combinatorial background**

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ii. **Jet-like correlations**

estimated with PYTHIA 8

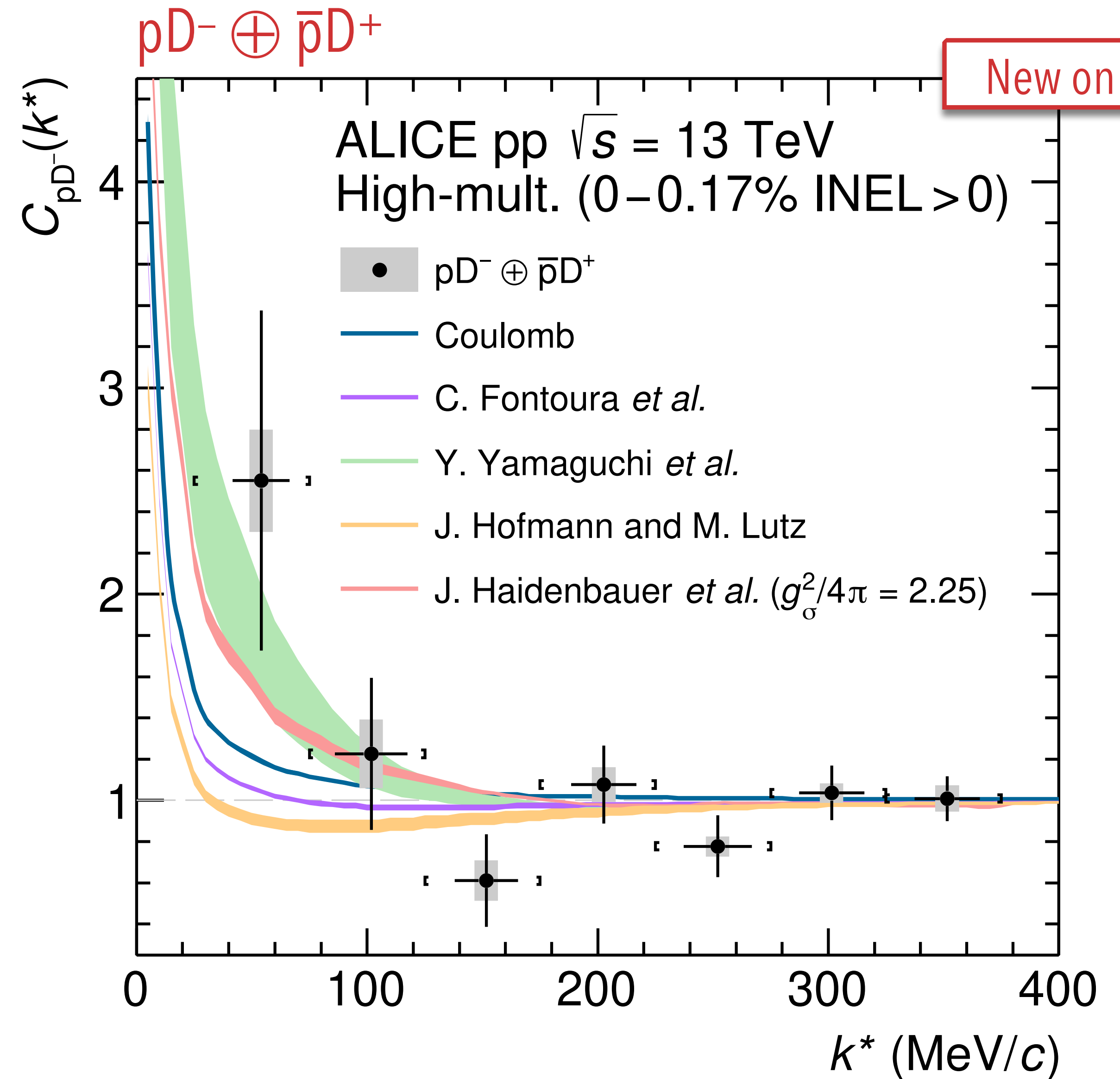
iii.  **$D^{*\pm} \rightarrow D^\pm + X$**

modelled with Coulomb-only interactions

- **Total background** well describes CF for large  $k^*$



- $pD^-$ 
  - ➔ Typically very small compared to other interactions (light-light  $\sim 7-8$  fm, light-strange  $\sim 1.5$  fm)
  - ➔ Most of the models predict repulsive interaction
  - ➔ Possible bound state formation (Yamaguchi et al)
- Data compatible with Coulomb only interaction, but comparison slightly improved when also attractive strong interaction is considered



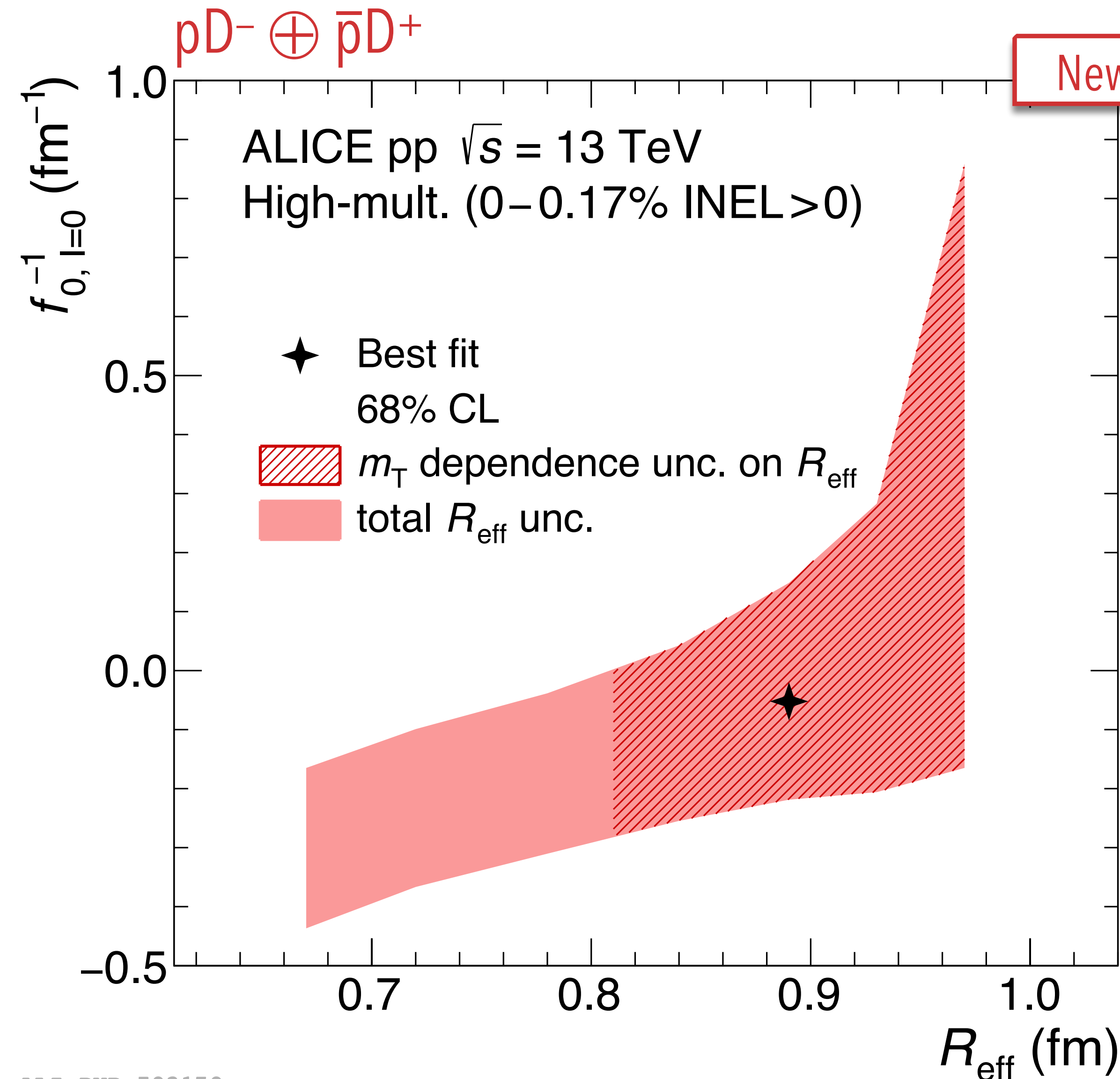
J. Haidenbauer et al, Eur. Phys. J. A33 (2007) 107–117

J. Hofmann and M. Lutz, Nucl. Phys. A 763 (2005) 90–139

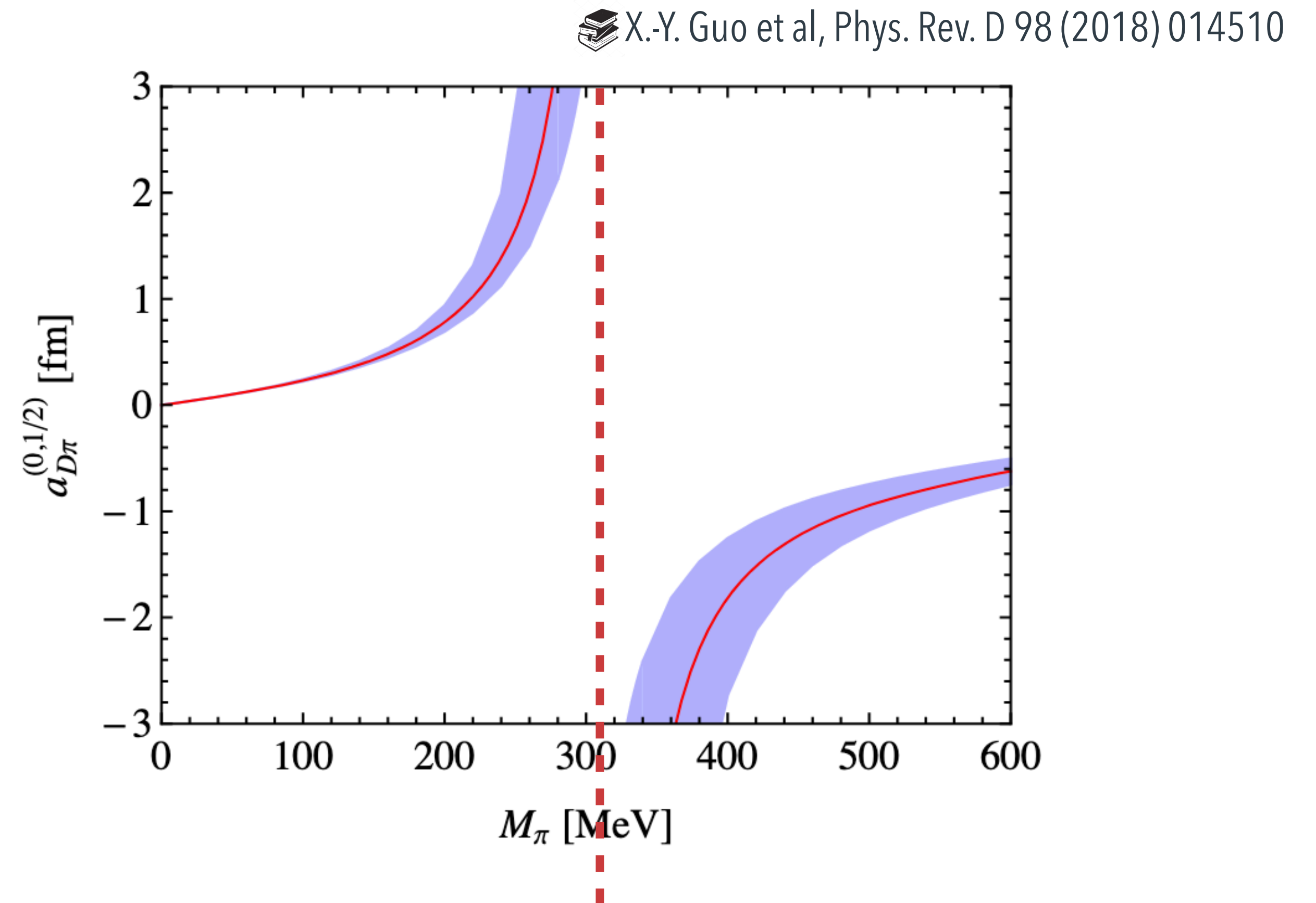
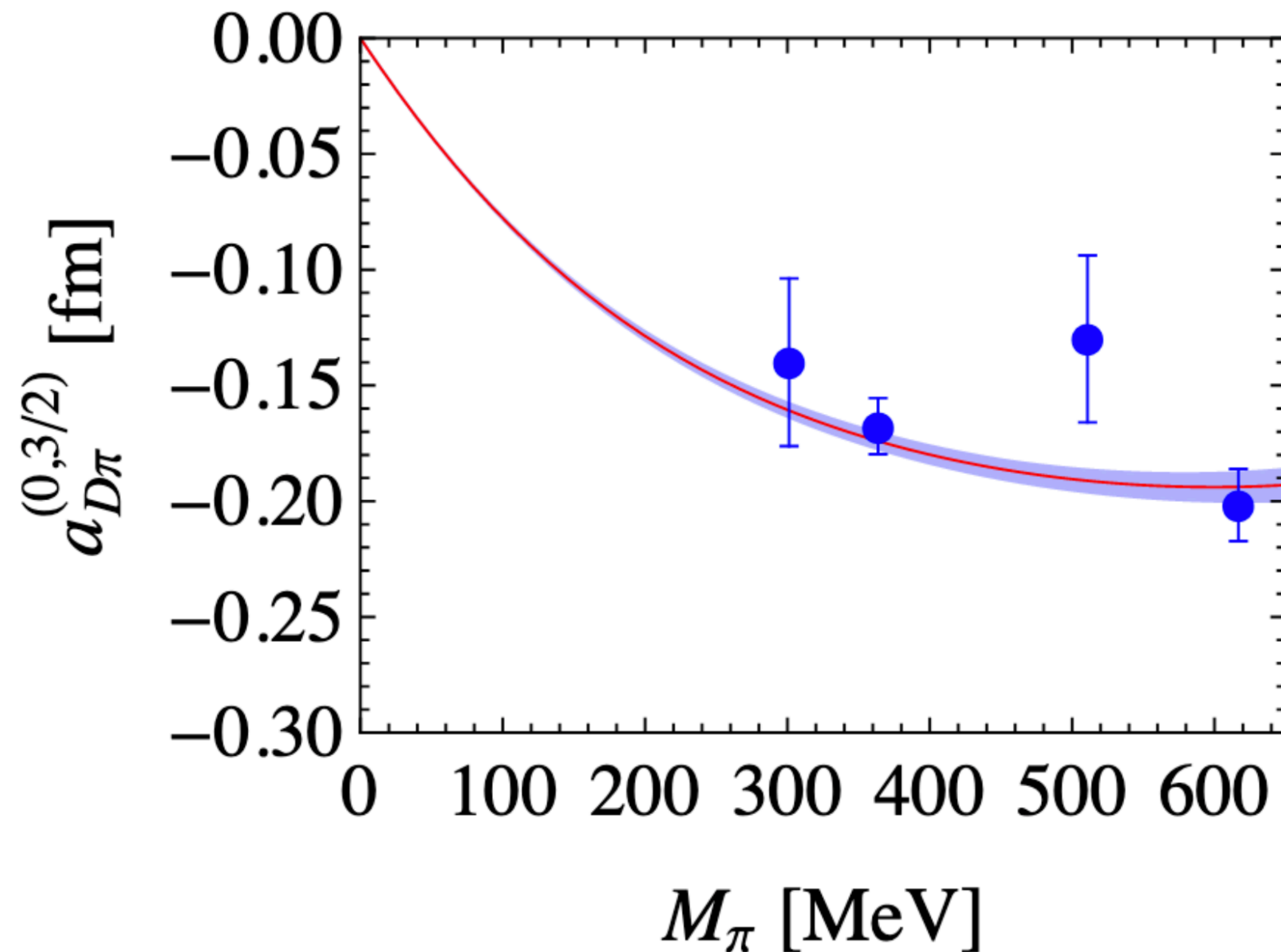
Fontoura et al, Phys. Rev. C 87 (2013) 025206

Yamaguchi et al, Phys. Rev. D84 (2011) 014032

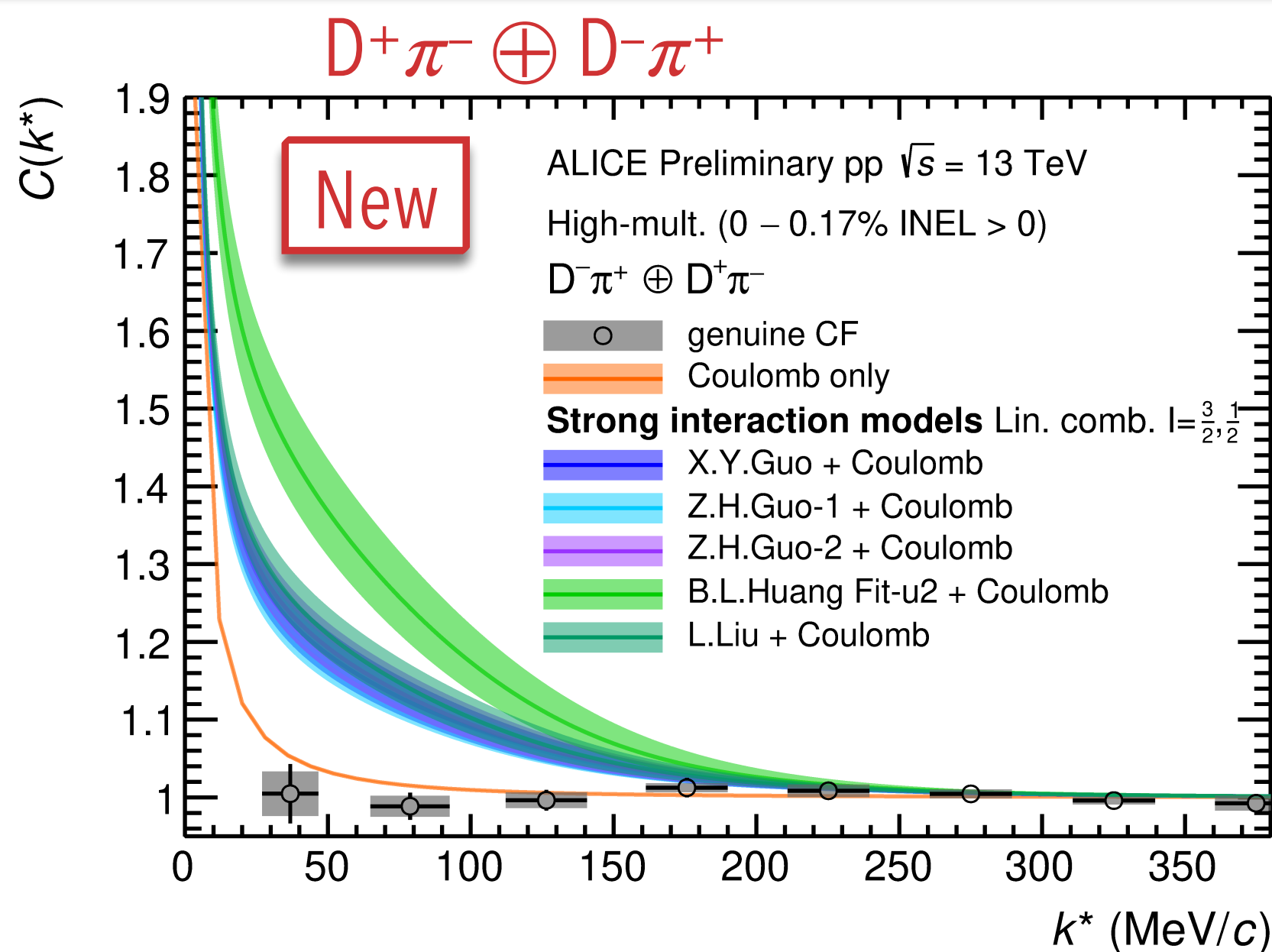
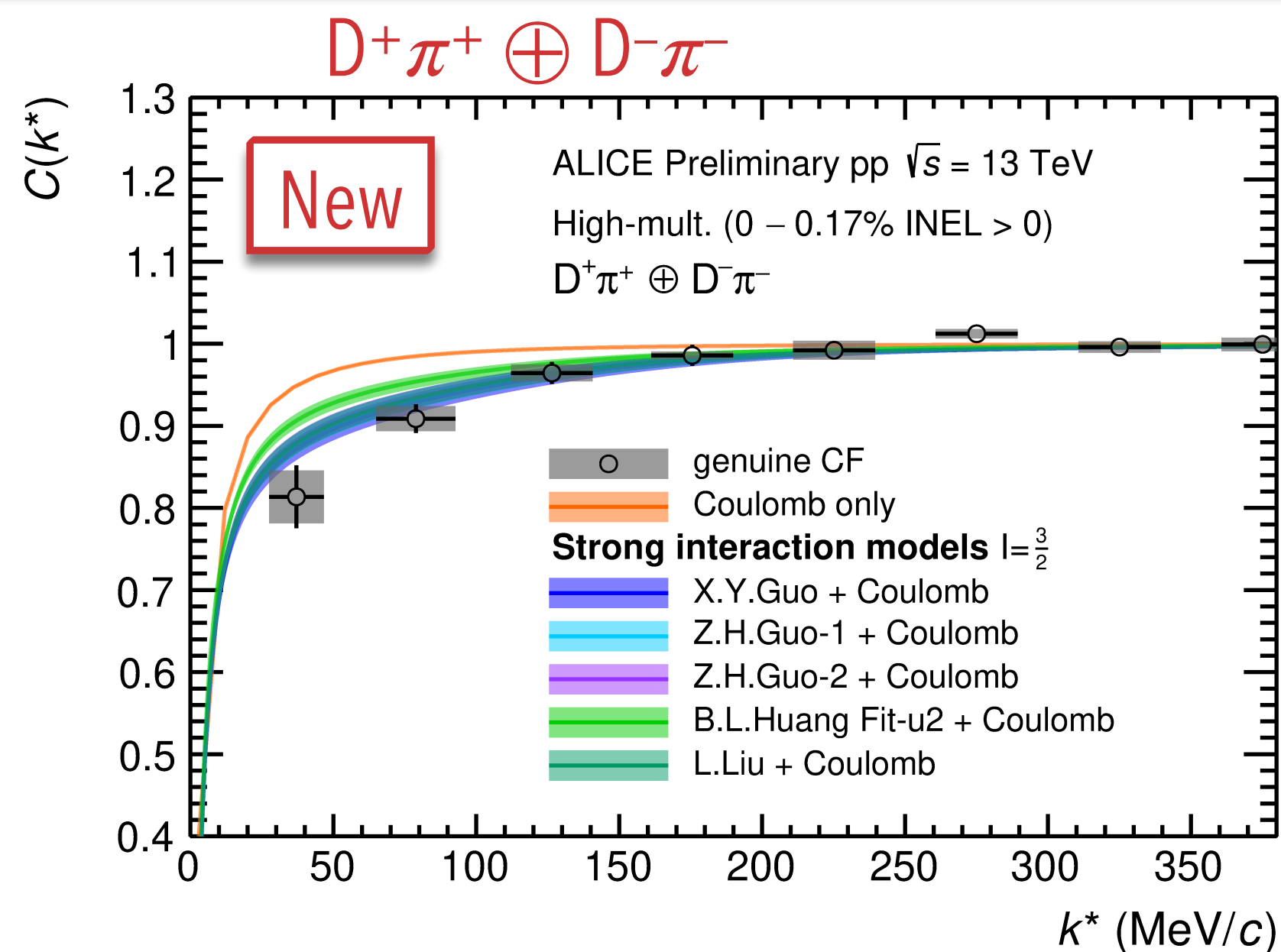
- Confidence interval of scattering length for isospin  $I=0$  channel evaluated by comparing data with CF computed using a Gaussian potential modelled with  $\rho$ -meson exchange
  - ➔ Assuming  $I=1$  negligible
  - ➔ First constrain to scattering length for  $I=0$
  - ➔ Indicates either rather shallow attractive interaction or strong attractive interaction with formation of a bound state



- Predictions of scattering lengths derived from lattice QCD calculations
  - ➔  $\sim 0.1-0.5$  fm: very small compared to other interactions (light-light  $\sim 7-8$  fm, light-strange  $\sim 1.5$  fm)
  - ➔ No constraints from data
  - ➔ For pions  $I=3/2$  channel more constrained than  $I=1/2$  channel

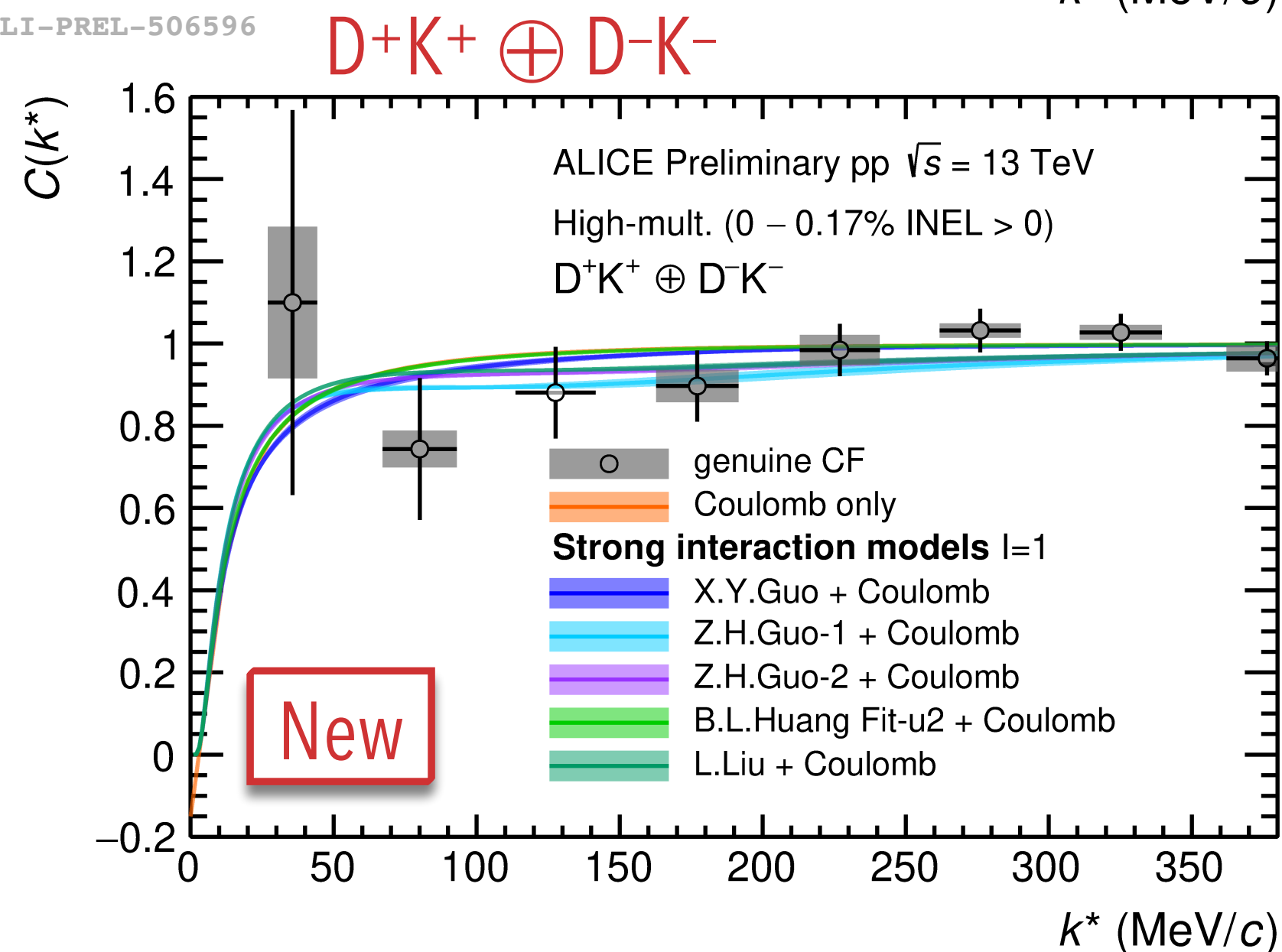


Bound-state pole formation corresponding to  $D_{s0}^*(2317)$



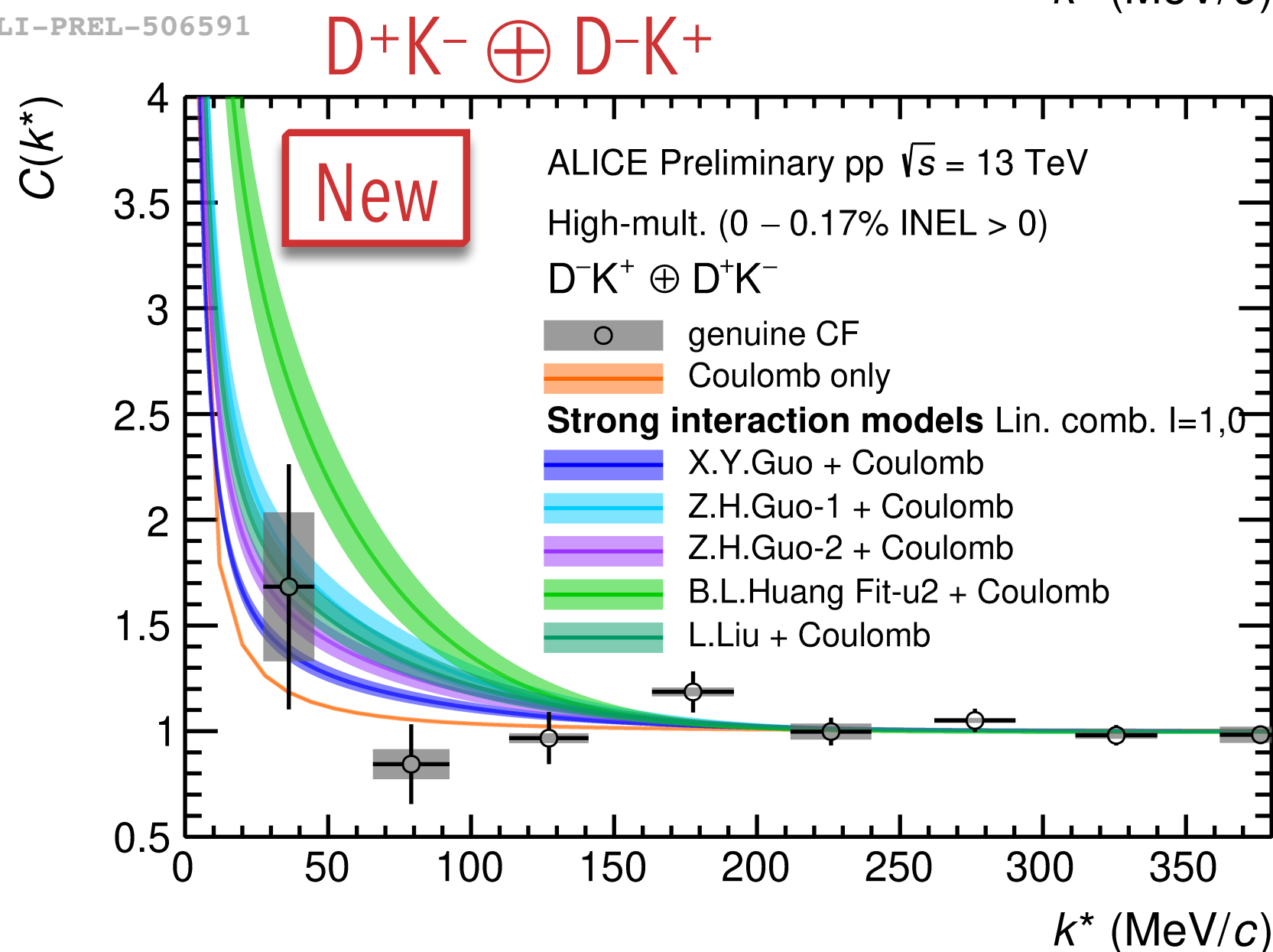
- Models agree with data in case of same-charge CF
- Models overestimate data in case of opposite-charge CF

ALI-PREL-506596



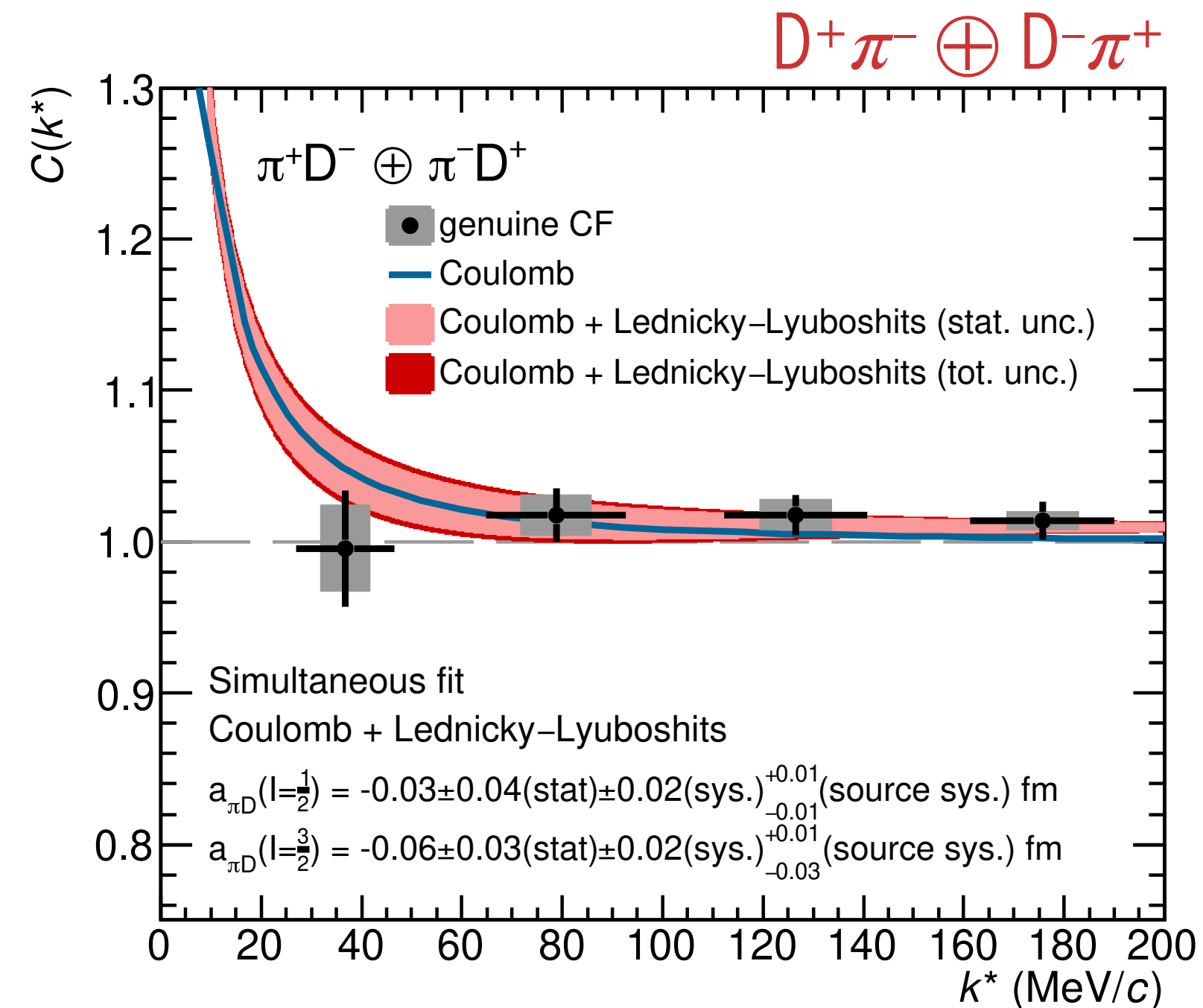
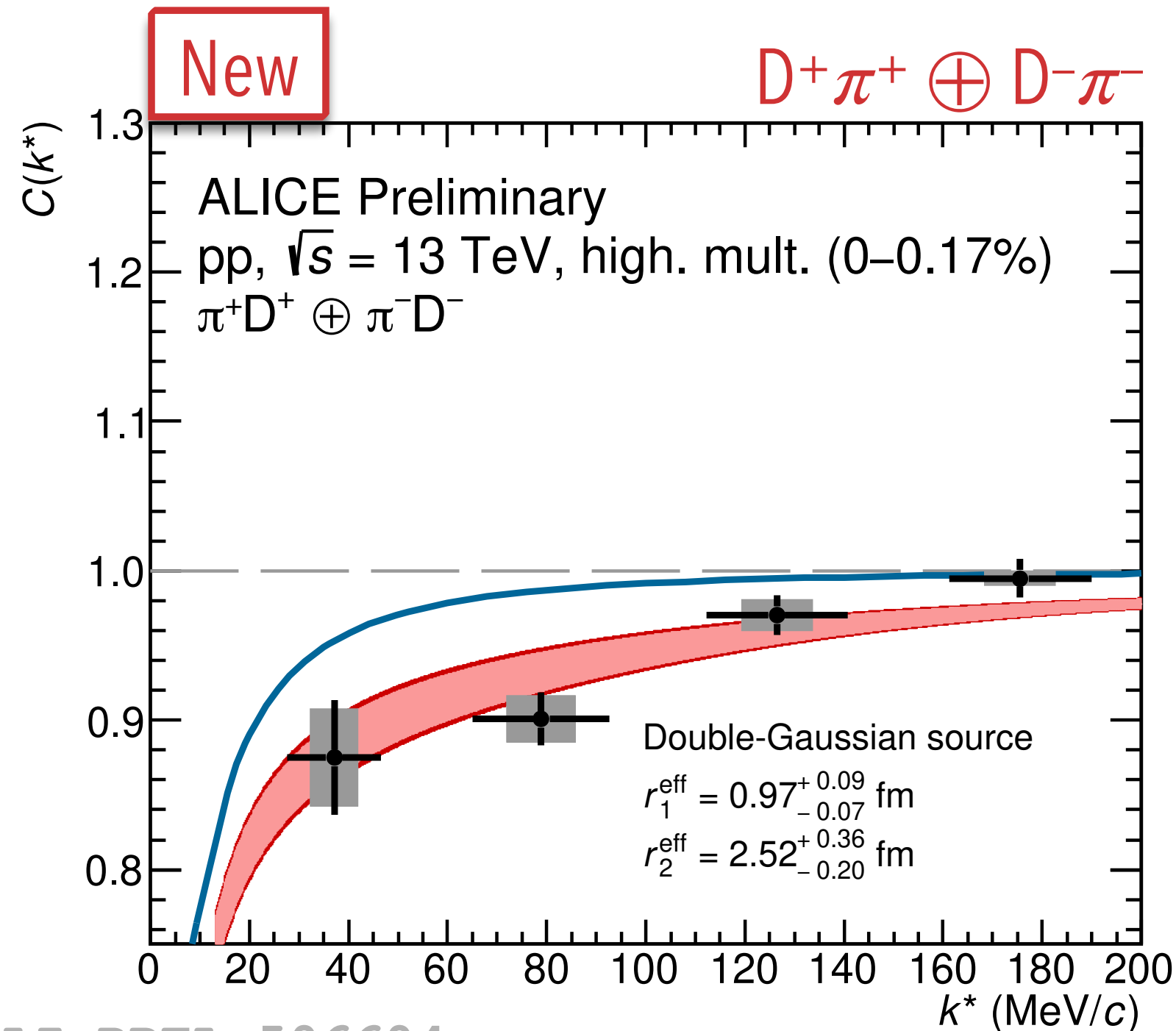
ALI-PREL-506586

ALI-PREL-506591

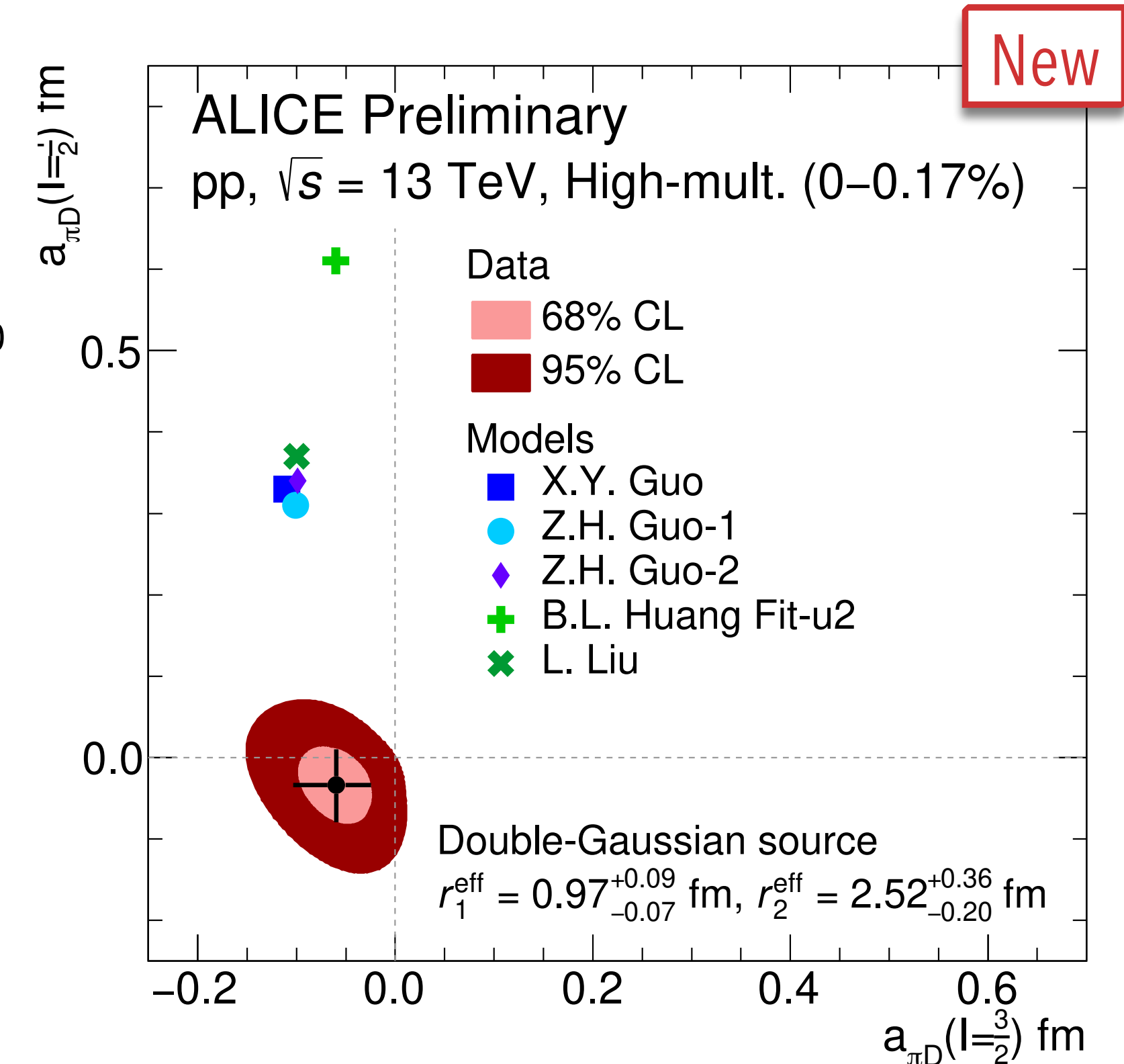


ALI-PREL-506581

L. Liu et al, Phys. Rev. D87 (2013) 014508  
 X.-Y. Guo et al, Phys. Rev. D 98 (2018) 014510  
 B.-L. Huang et al, Phys. Rev. D 105 (2022) 036016  
 Z.-H. Guo et al Eur. Phys. J. C 79 (2019) 13

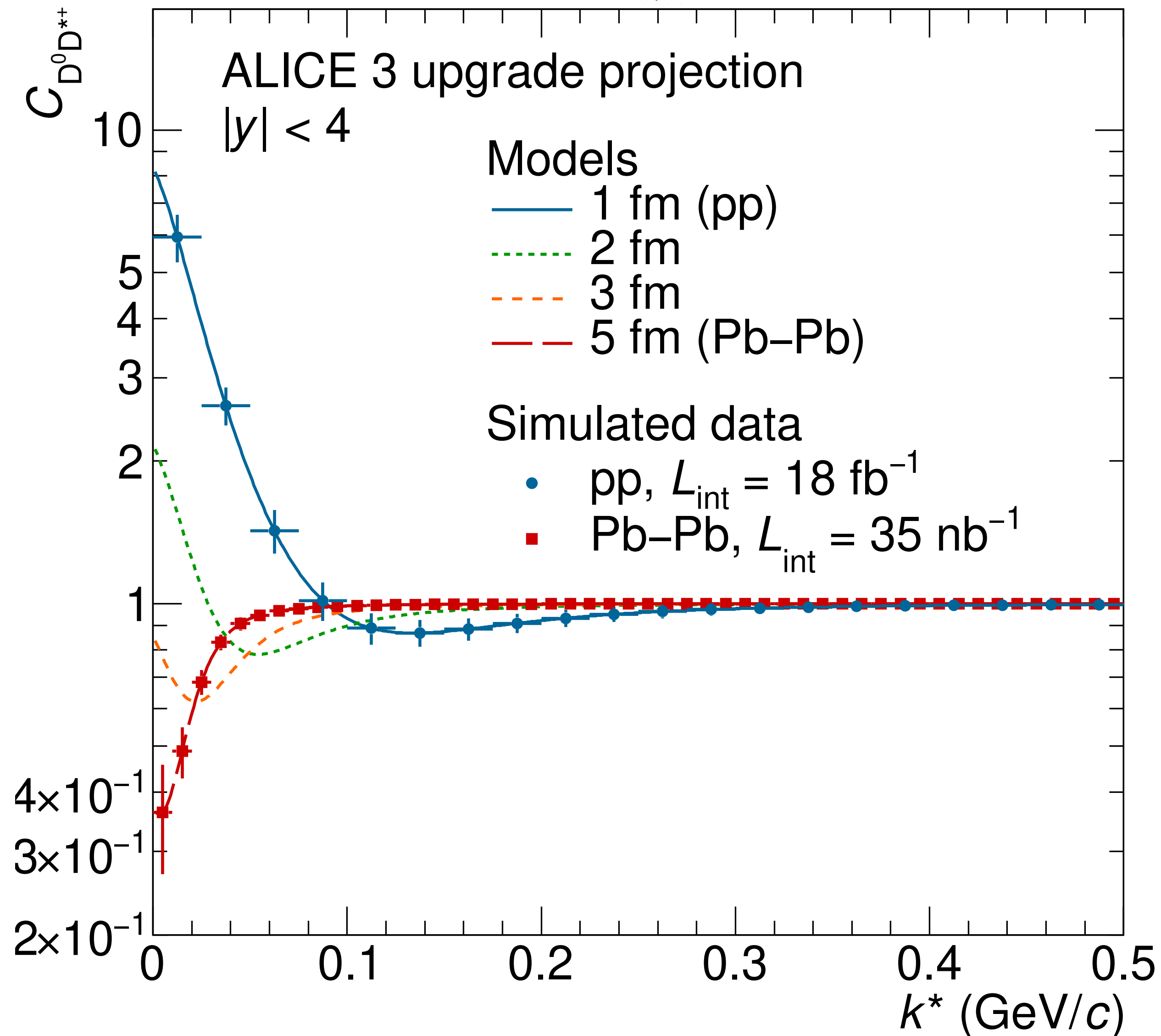


- $\pi^+D^+$
- $l=3/2$  channel only
- $\pi^+D^-$
- $l=3/2$  (33%)  
 $l=1/2$  (66%)



- The values found indicate a small rescattering of D mesons in the hadronic phase of heavy-ion collisions
- Scattering length for  $l=3/2$  in agreement with models
- Scattering length for  $l=1/2$  significantly smaller than models

ALICE3 LOI: CERN-LHCC-2022-009



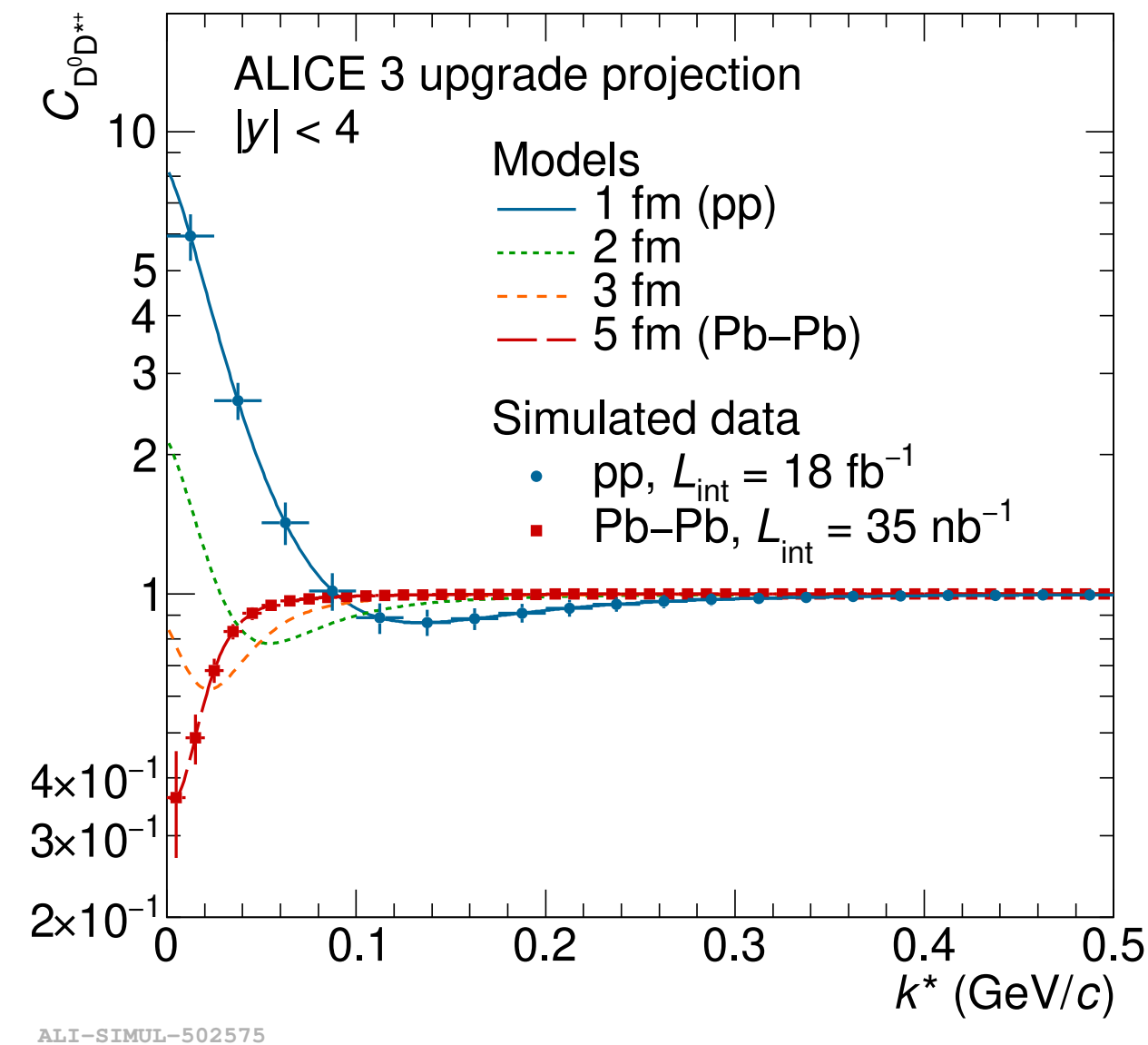
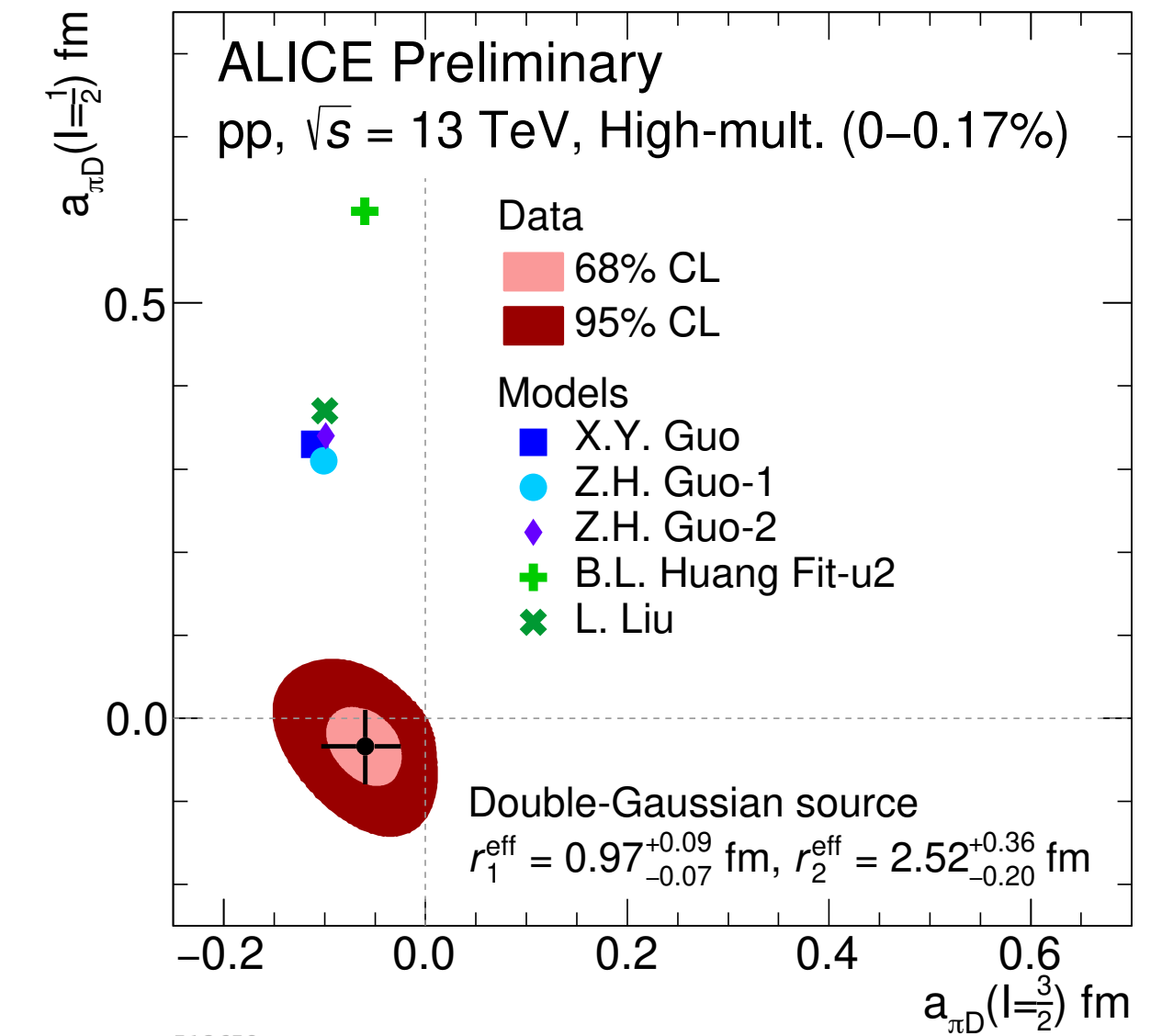
- ALICE 3: large acceptance, high luminosity, excellent spatial resolution
  - ➔ Run 5: ideal laboratory for the measurement of charm-hadron momentum correlations in different colliding systems
- Interplay between system size and scattering length
  - ➔ size-dependent modification of the correlation function in presence of a bound state

Yuki Kamyia et al, arXiv:2203.13814

ALICE upgrades, H. S. Scheid talk 07/04 2022, 16:00

- First studies of residual strong interaction between charm and light hadrons performed with Run 2 data
  - ➔ Shallow interaction between charm mesons and light hadrons suggests no important hadronic phase for heavy-flavour hadrons in heavy-ion collisions
  - ➔ Possible formation of bound state in ND interaction not excluded
  - ➔ Significant improvement foreseen with Run 3 data

- ALICE 3 will provide an ideal laboratory for the study of residual strong interaction among charm hadrons and for the search of charm molecular states



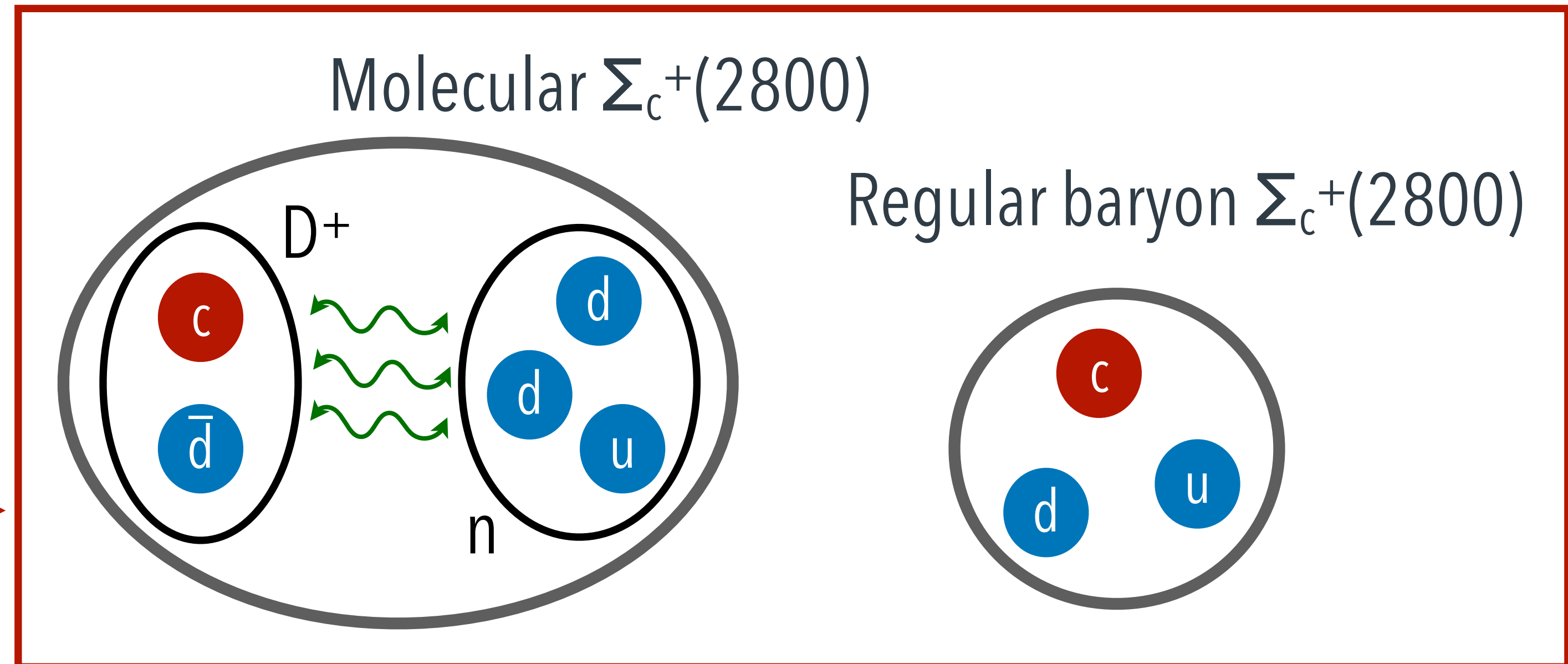
A stylized sun with a white center and red rays on a dark grey background. The sun is composed of a white circle with a grey outline, and the rays are made of numerous thin, red lines radiating outwards. The background is a dark grey color.

**ADDITIONAL SLIDES**

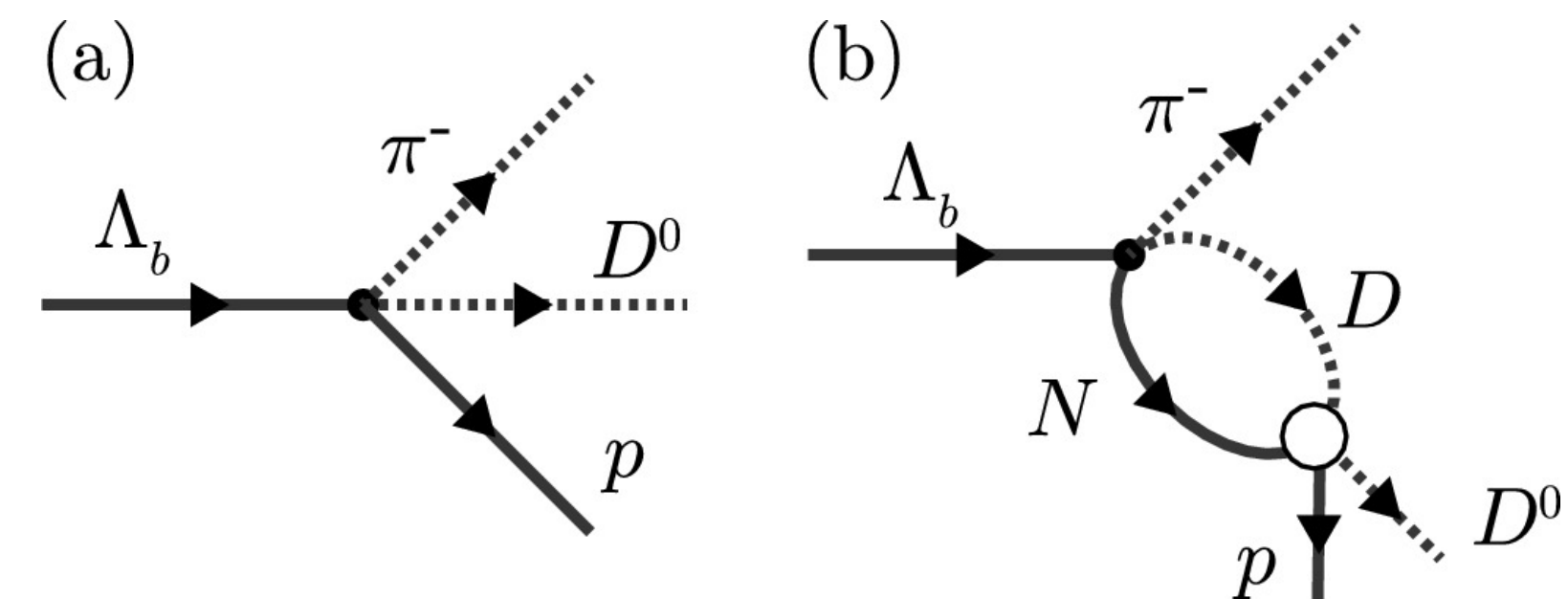


- Charm molecules?

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$\Sigma_c\bar{D}^*$	$1/2 (3/2^-)$	$P_c(4440)$



- Proposed as molecular state in J. Haidenbauer et al, Eur. Phys. J. A 47, 18 (2011)
- S. Sakai et al, Phys. Lett. B 808 (2020) 135623



Fang-Zheng Peng et al, Phys. Rev. D 105, 034028 (2022)

- Molecular states also relevant to explain some beauty-hadron decays

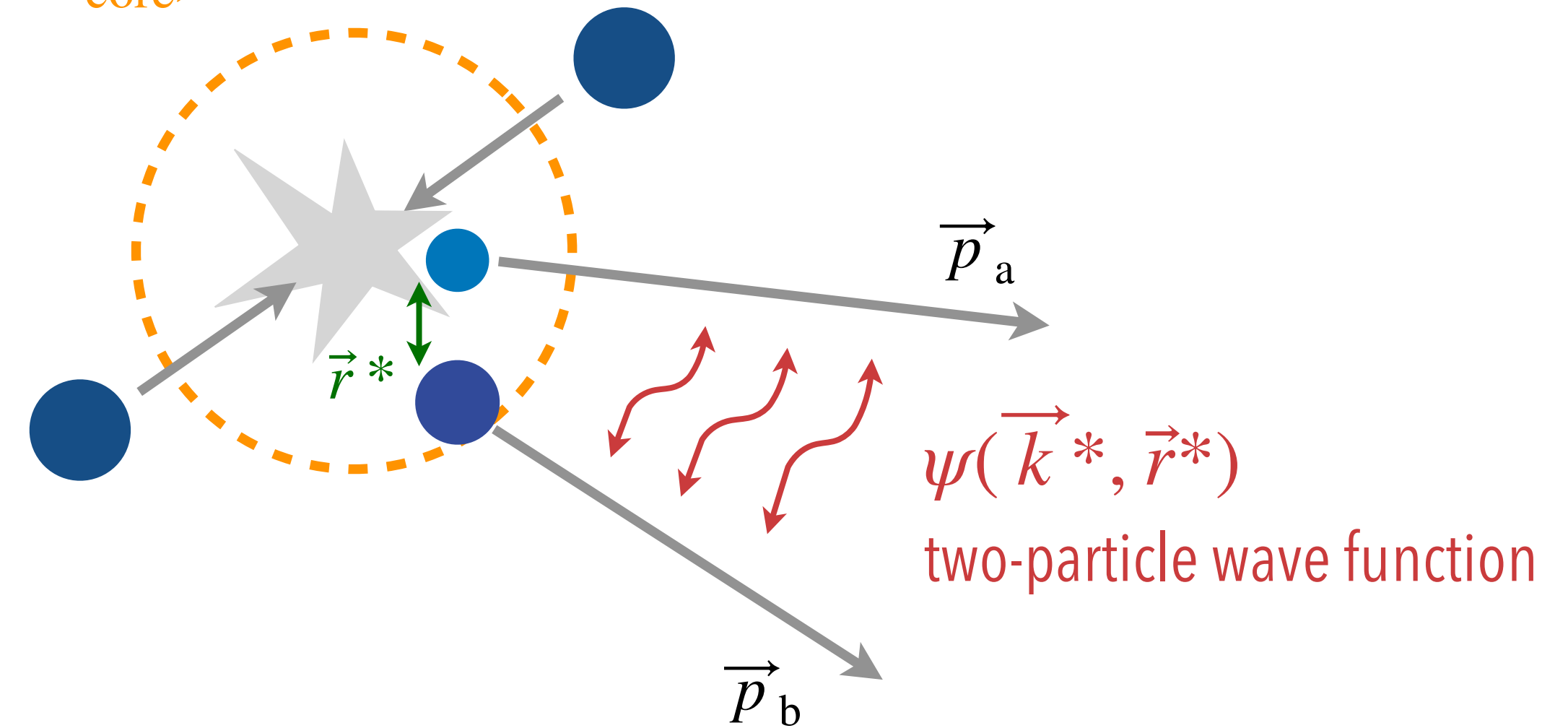
$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*$$

→ **Emitting source**: hypersurface at kinematic freezeout of final-state particles

- Described with a Gaussian core

$$G(r^*, r_{\text{core}}(m_T)) = \frac{1}{(4\pi r_{\text{core}}^2(m_T))^{3/2}} \cdot \exp\left(-\frac{r^{*2}}{4r_{\text{core}}^2(m_T)}\right)$$

$G(\vec{r}^*, r_{\text{core}})$  Gaussian core



$$C(\vec{k}^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3 r^*$$

$S(\vec{r}^*)$  source function  
 $G(\vec{r}^*, r_{\text{core}})$  Gaussian core

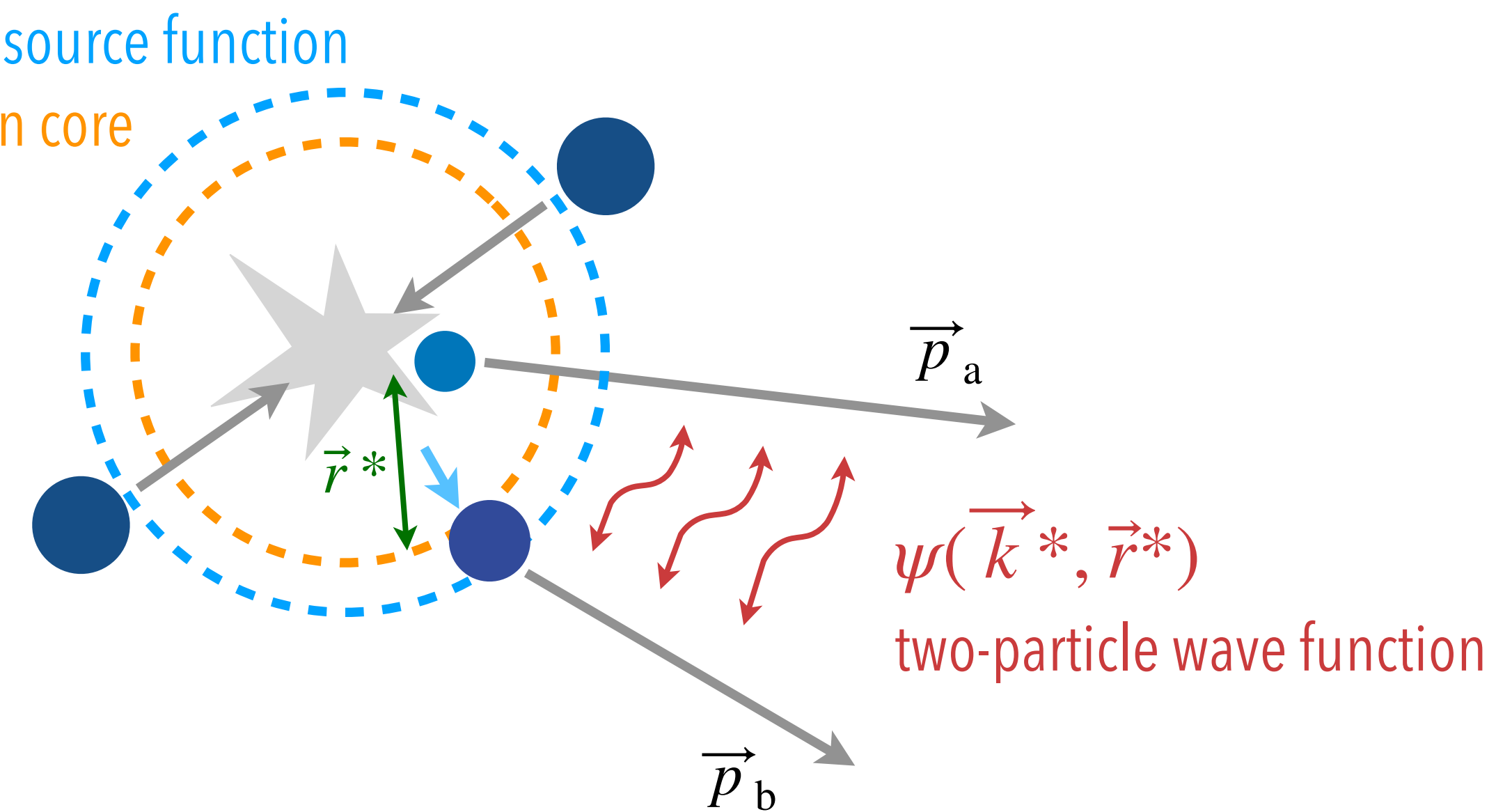
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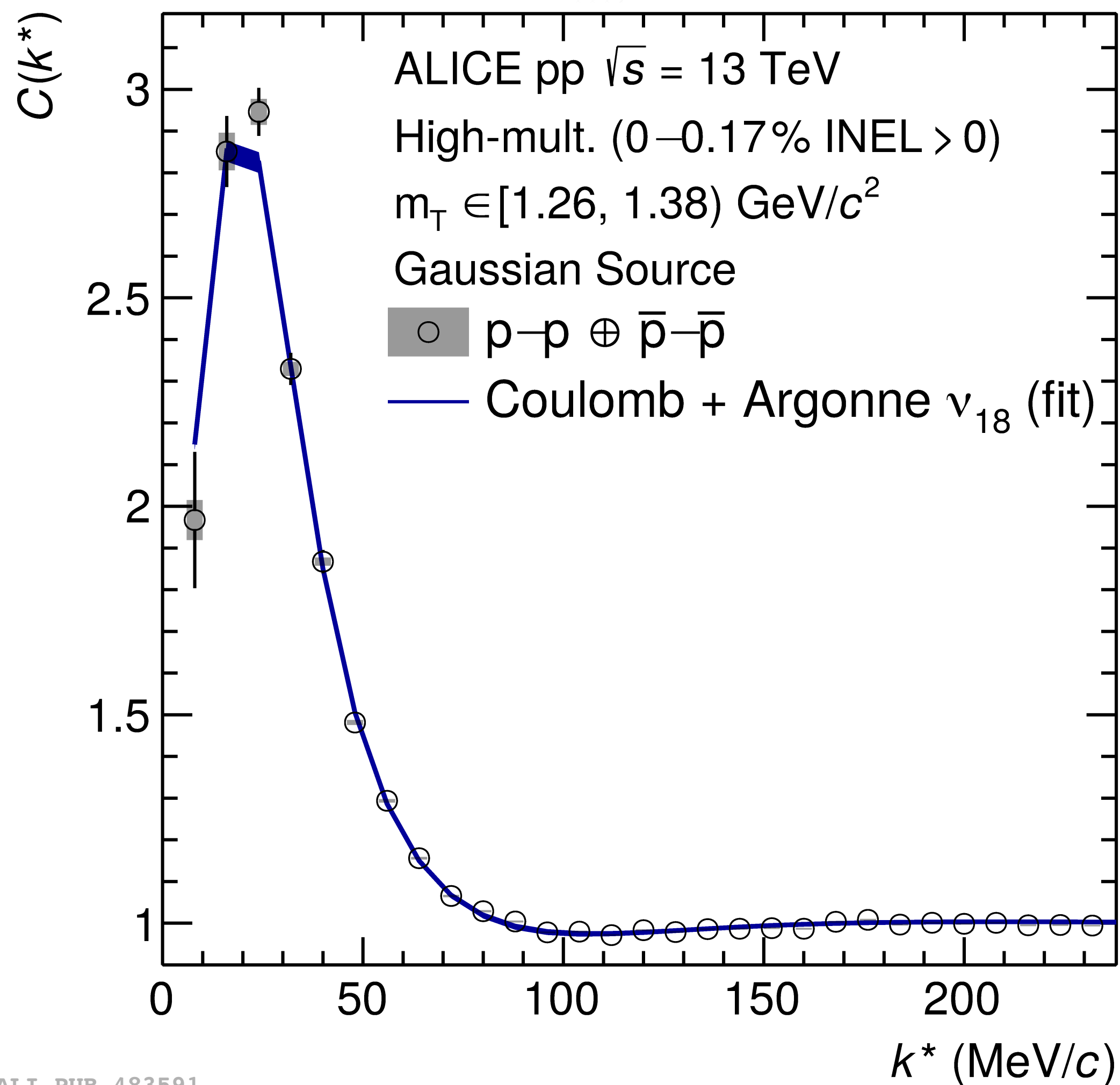
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- Short-lived strongly decaying resonances effectively enlarge it

$$E(r^*, M_{\text{res}}, \tau_{\text{res}}, p_{\text{res}}) = \frac{1}{s} \exp\left(-\frac{r^*}{s}\right) \quad \text{with} \quad s = \beta\gamma\tau_{\text{res}} = \frac{p_{\text{res}}}{M_{\text{res}}}\tau_{\text{res}}$$



Phys. Lett. B 811 (2020) 135849

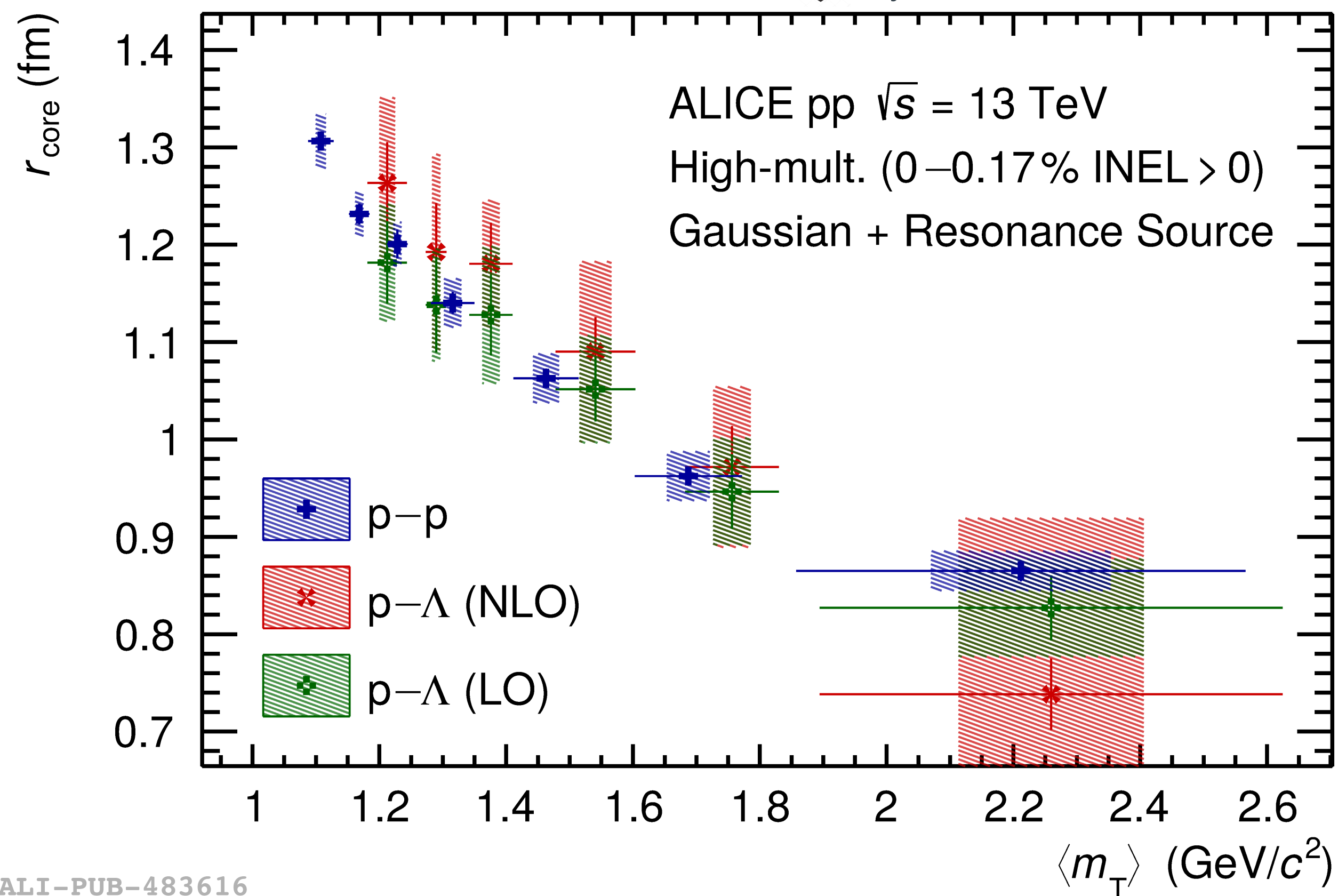


ALI-PUB-483591

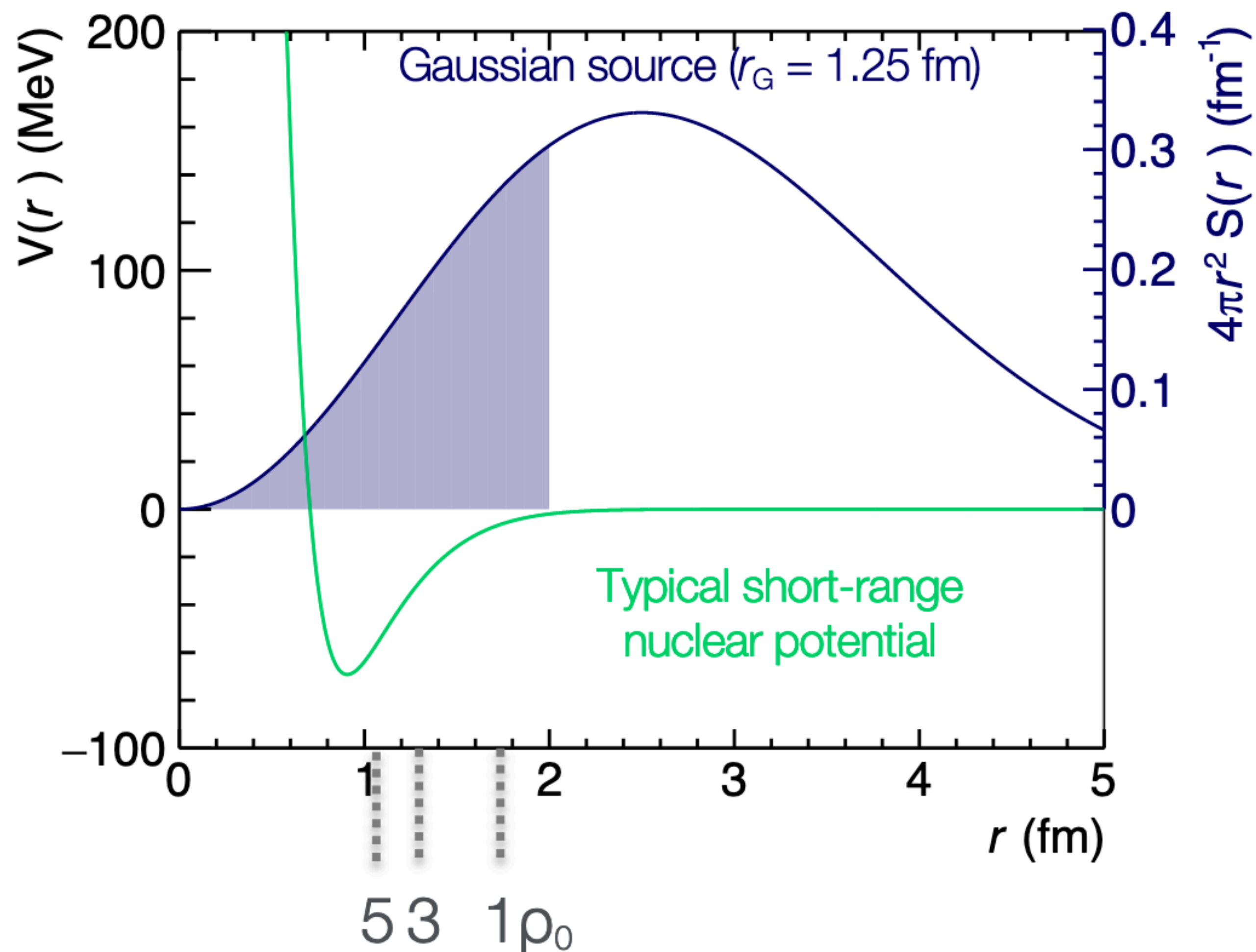
- Source size  $\sim 1$  fm makes the high-multiplicity pp system suitable for the study of hadron–hadron interactions

- Fit correlation functions of p–p and p– $\Lambda$  pairs
  - ➔ Interaction precisely described
  - ➔ Gaussian source with radius as free parameter

Phys. Lett. B 811 (2020) 135849

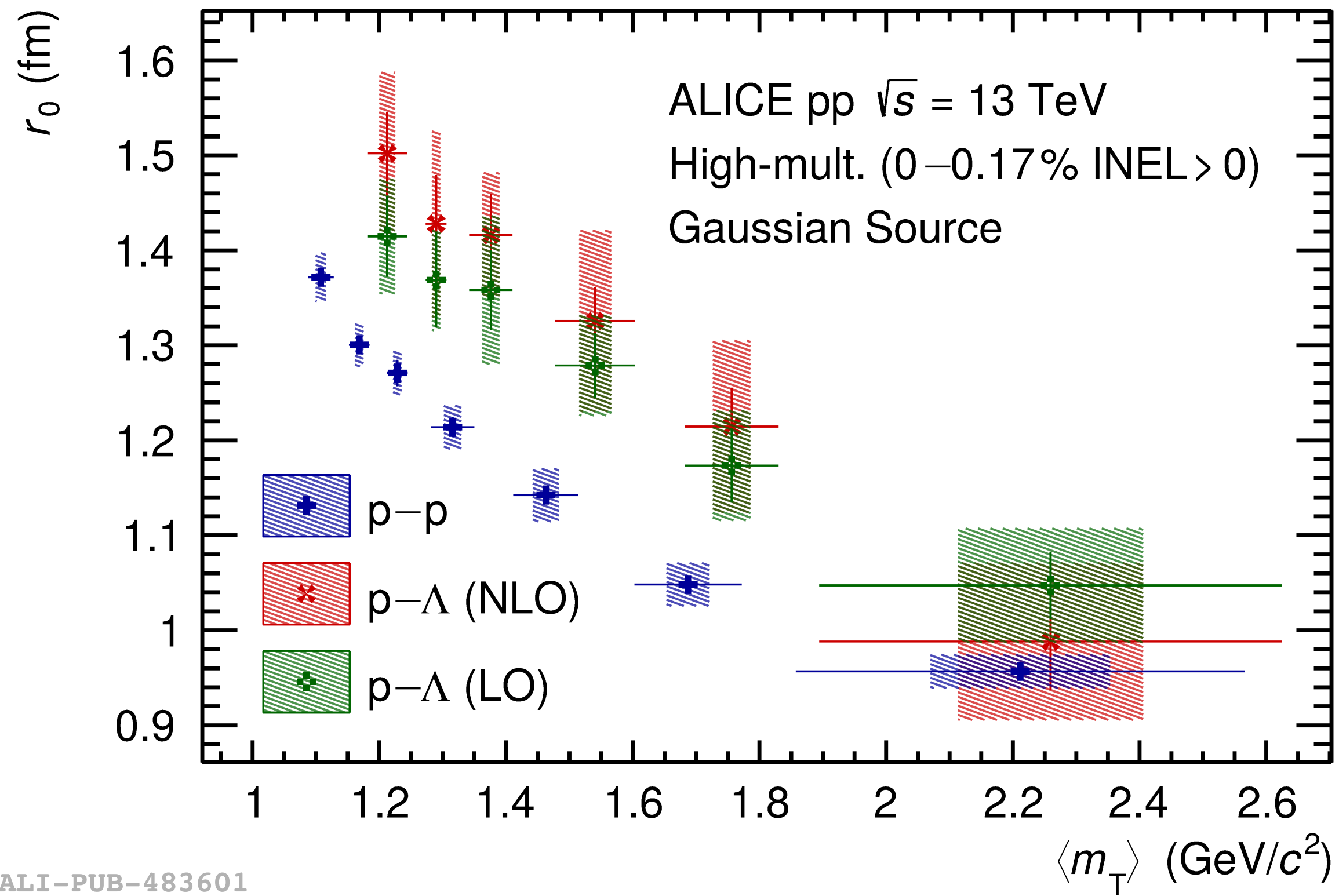


ALI-PUB-483616



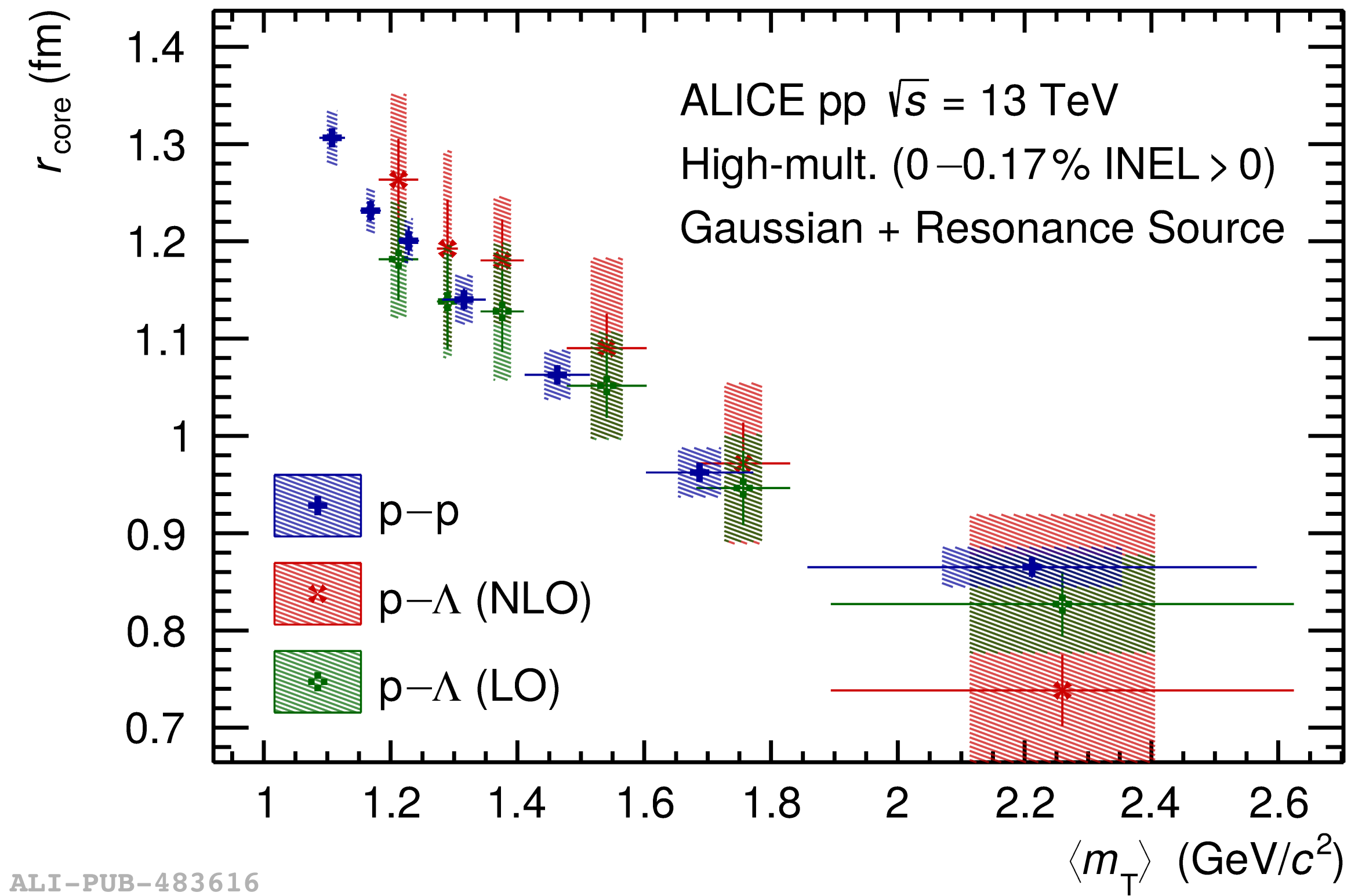
- Typical range of nuclear potential around 1-2 fm
  - ➔ study of strong interaction among hadrons not possible with larger sources
  - ➔ proton–proton and proton–nucleus collisions are the ideal laboratory to study the strong interaction

● Without considering resonances

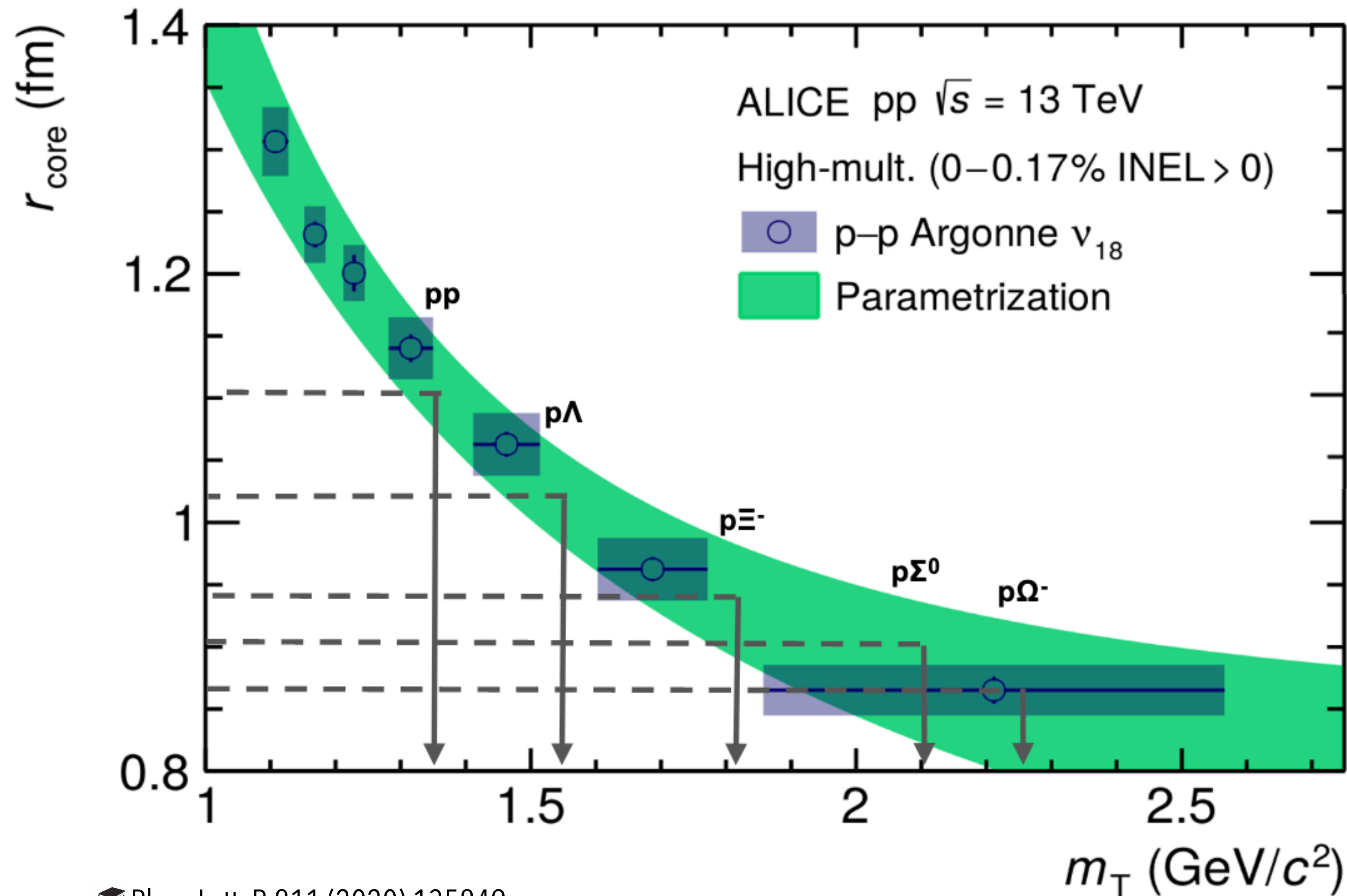


ALI-PUB-483601

● Considering resonances



ALI-PUB-483616



- Measurement of source radius obtained from p–p correlation used to obtain the values for other baryon species

- Decay channel:

- D<sup>±</sup> → K<sup>∓</sup>π<sup>±</sup>π<sup>±</sup>

- BR = (9.38 ± 0.16)%

PDG, Prog. Theor. Exp. Phys. 2020 083C01

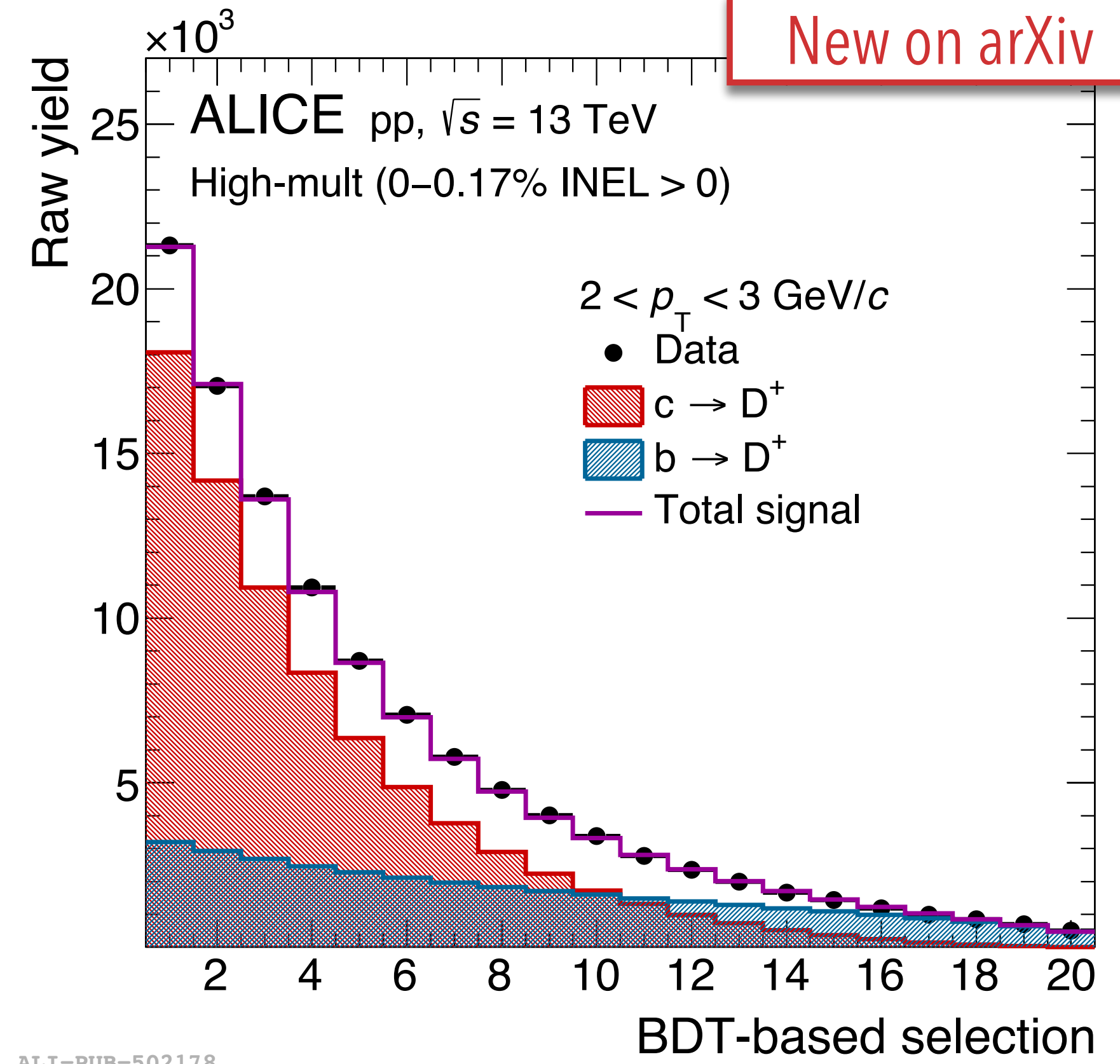
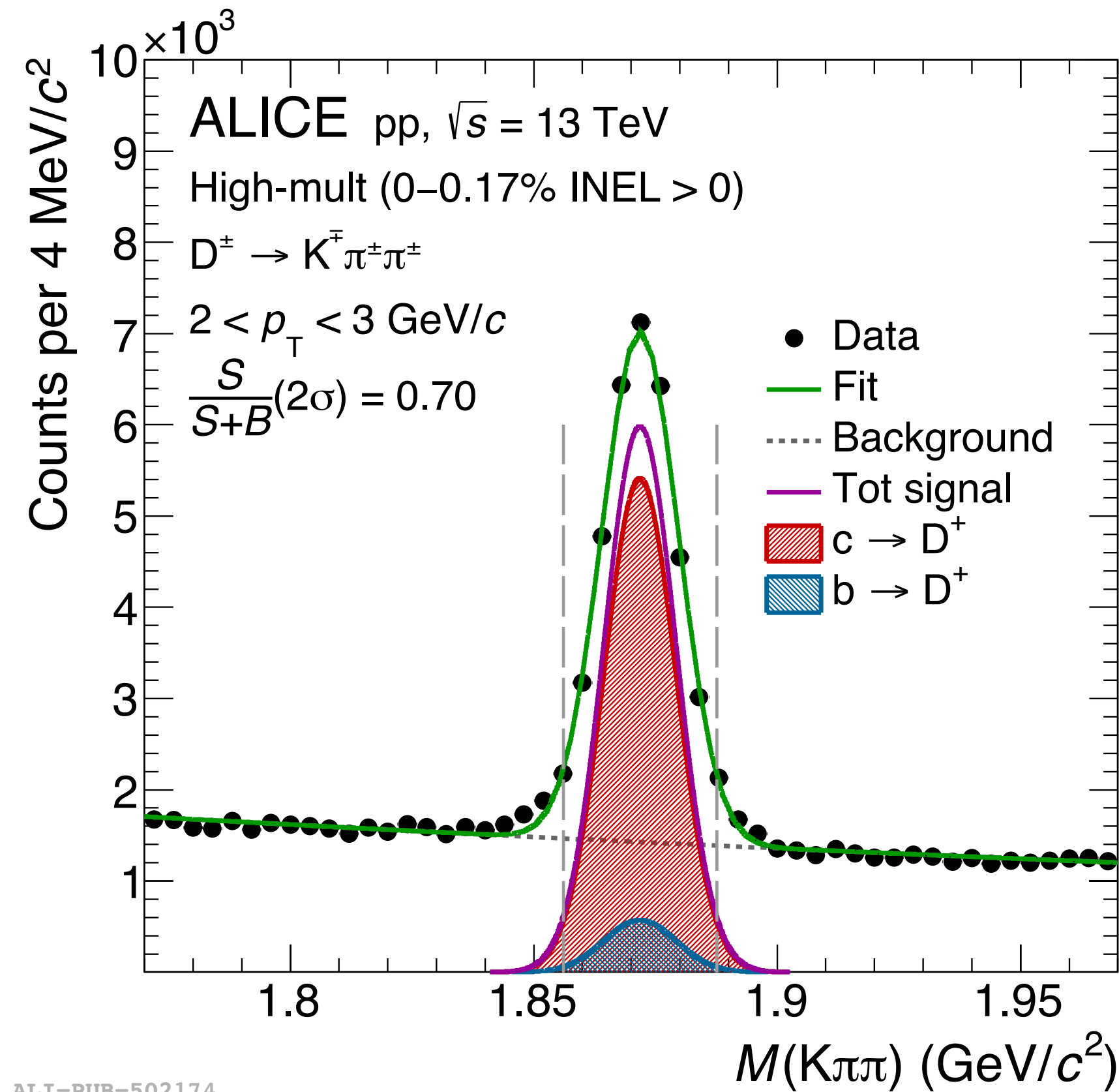
- Contributions:

- **Prompt:**

- ▶ from c hadronisation

- ▶ from D\* decays

- **Non-prompt:** from b decays



New on arXiv

- Fully reconstruct decay topologies ( $c\tau \approx 312 \mu\text{m}$ )

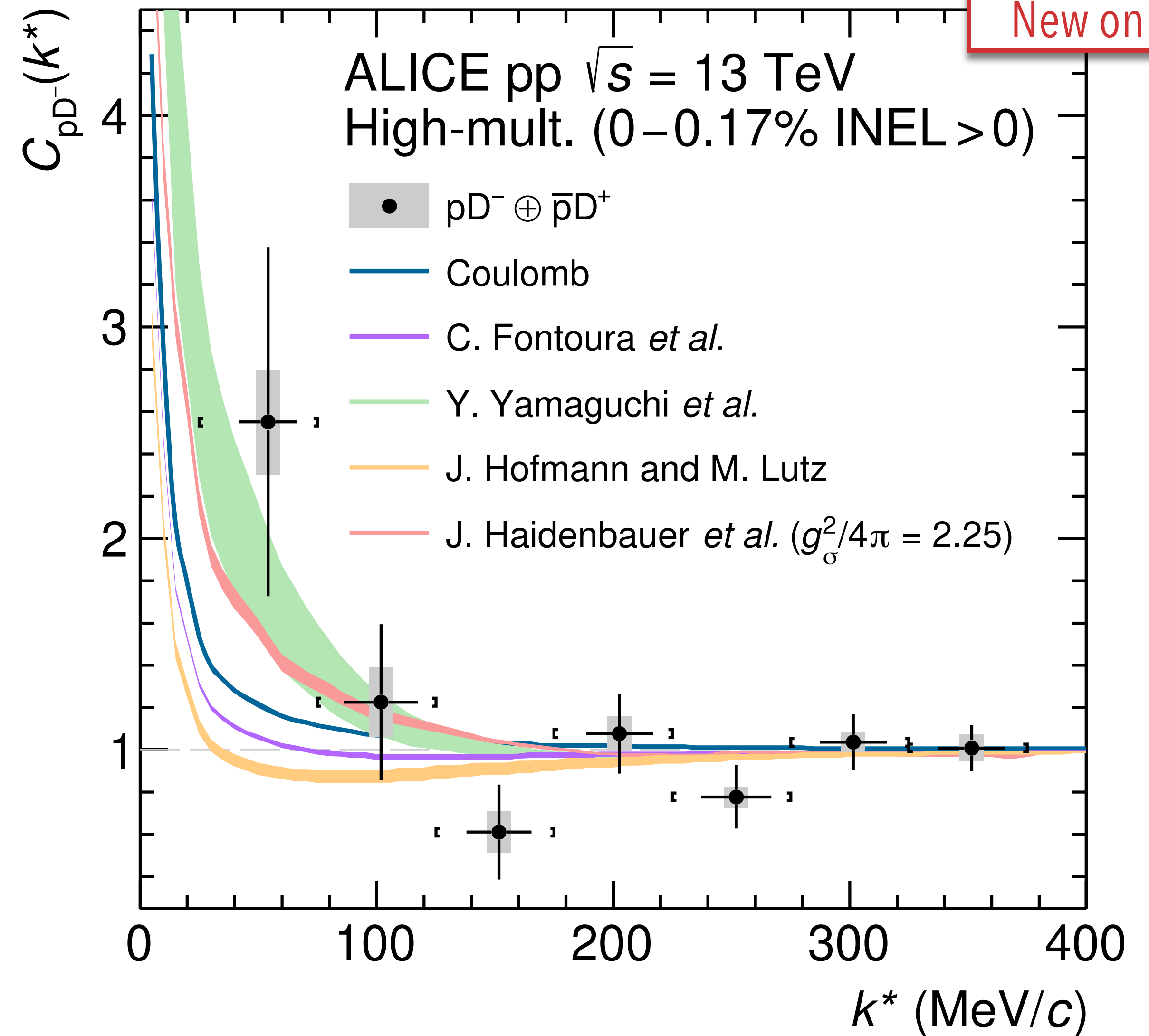
- Multiclass Boosted Decision Trees feed with geometrical, kinematic and PID variables used to reduce combinatorial background and non-prompt contribution

- Purity of about 70%, non-prompt contribution about 7%

arXiv: 2201.05352



Model	$f_0(l=0)$ [fm]	$f_0(l=1)$ [fm]
Haidenbauer $g_\sigma^2/4\pi = 1$ Meson-exchange model	0,14	-0,28
Haidenbauer $g_\sigma^2/4\pi = 2.25$ Meson-exchange model	0,67	0,04
Hofmann and Lutz SU(4) contact interaction	-0,16	-0,26
Yamaguchi meson-exchange on HQ symmetry	-4,38	-0,07
Fontoura Chiral-quark model	0,16	-0,25







J. Haidenbauer *et al.*, Eur. Phys. J. A33 (2007) 107–117

J. Hofmann and M. Lutz, Nucl. Phys. A 763 (2005) 90–139

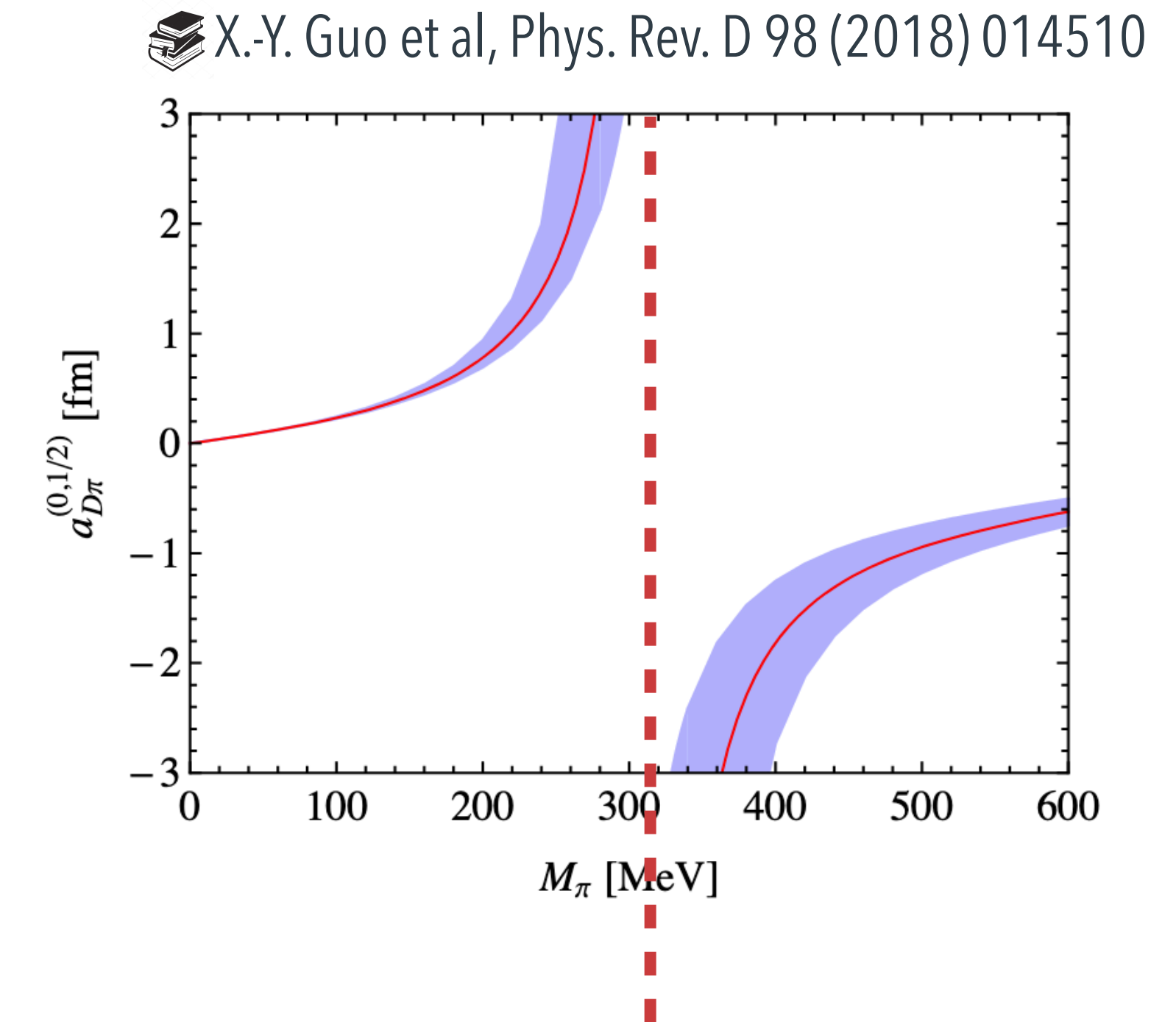
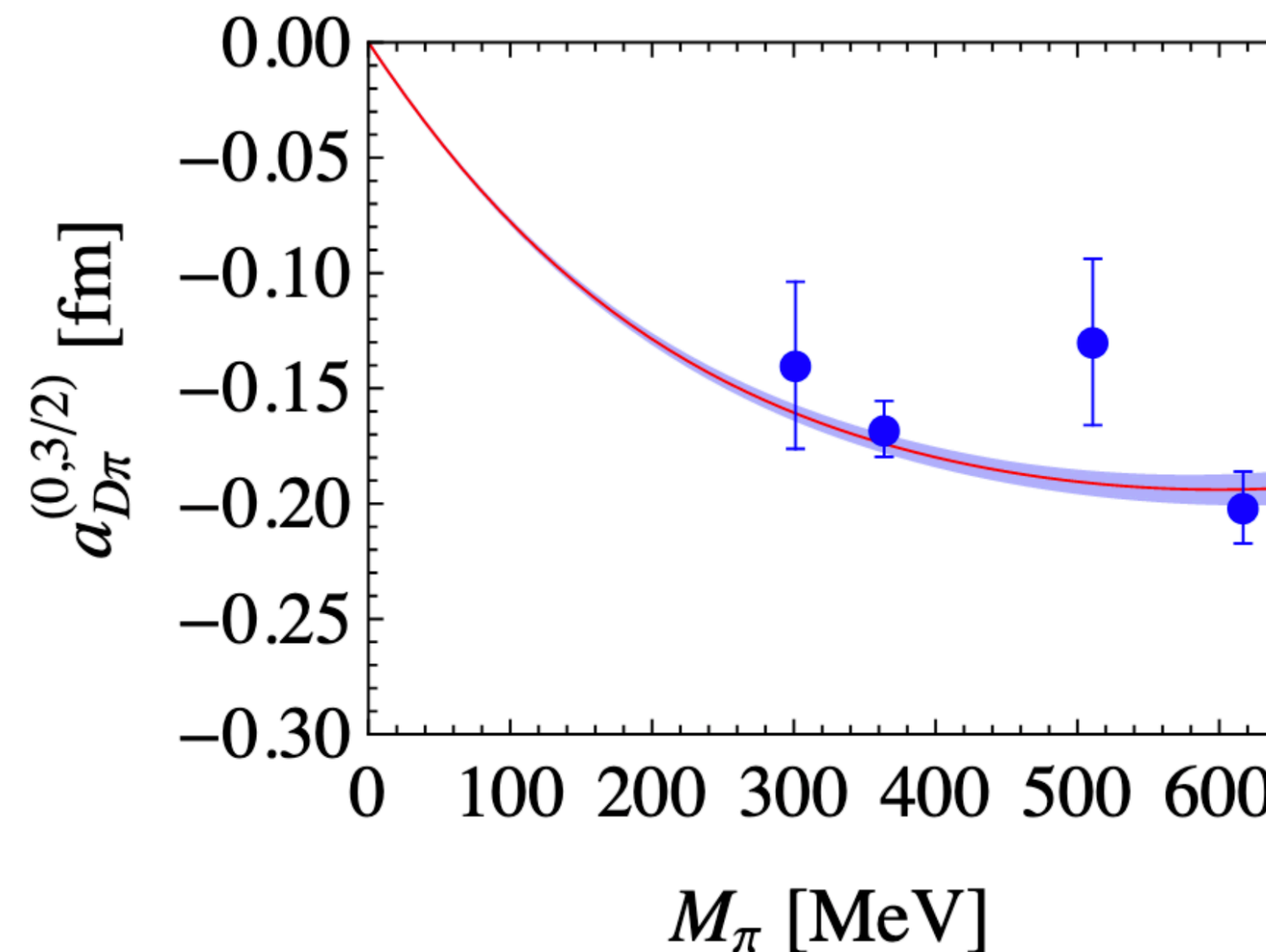
Fontoura *et al.*, Phys. Rev. C 87 (2013) 025206

Yamaguchi *et al.*, Phys. Rev. D84 (2011) 014032

Channel	L. Liu	X.-Y. Guo	Z.-H. Guo-1	Z.-H. Guo-2	B.-L. Huang
$D\pi(l=3/2)$ [fm]	-0,10	-0,11	-0,101	-0,099	-0,06
$D\pi(l=1/2)$ [fm]	0,37	0,33	0,31	0,34	0,61
$DK(l=1)$ [fm]	$0,07+i0,17$	-0,05	$0,06+i0,30$	$0,05+i0,17$	-0,01
$D\bar{K}(l=0)$ [fm]	0,84	0,46	0,96	0,68	1,81
$D\bar{K}(l=1)$ [fm]	-0,20	-0,22	-0,18	-0,19	-0,24

-  L. Liu et al, Phys. Rev. D87 (2013) 014508
-  X.-Y. Guo et al, Phys. Rev. D 98 (2018) 014510
-  B.-L. Huang et al, Phys. Rev. D 105 (2022) 036016
-  Z.-H. Guo et al Eur. Phys. J. C 79 (2019) 13

- Predictions of scattering lengths derived from lattice QCD calculations
  - ➔ Typically very small compared to other interactions (light-light  $\sim$  7-8 fm, light-strange  $\sim$  1.5 fm)
  - ➔ No constraints from data
  - ➔ For pions  $l=3/2$  channel more constrained than  $l=1/2$  channel

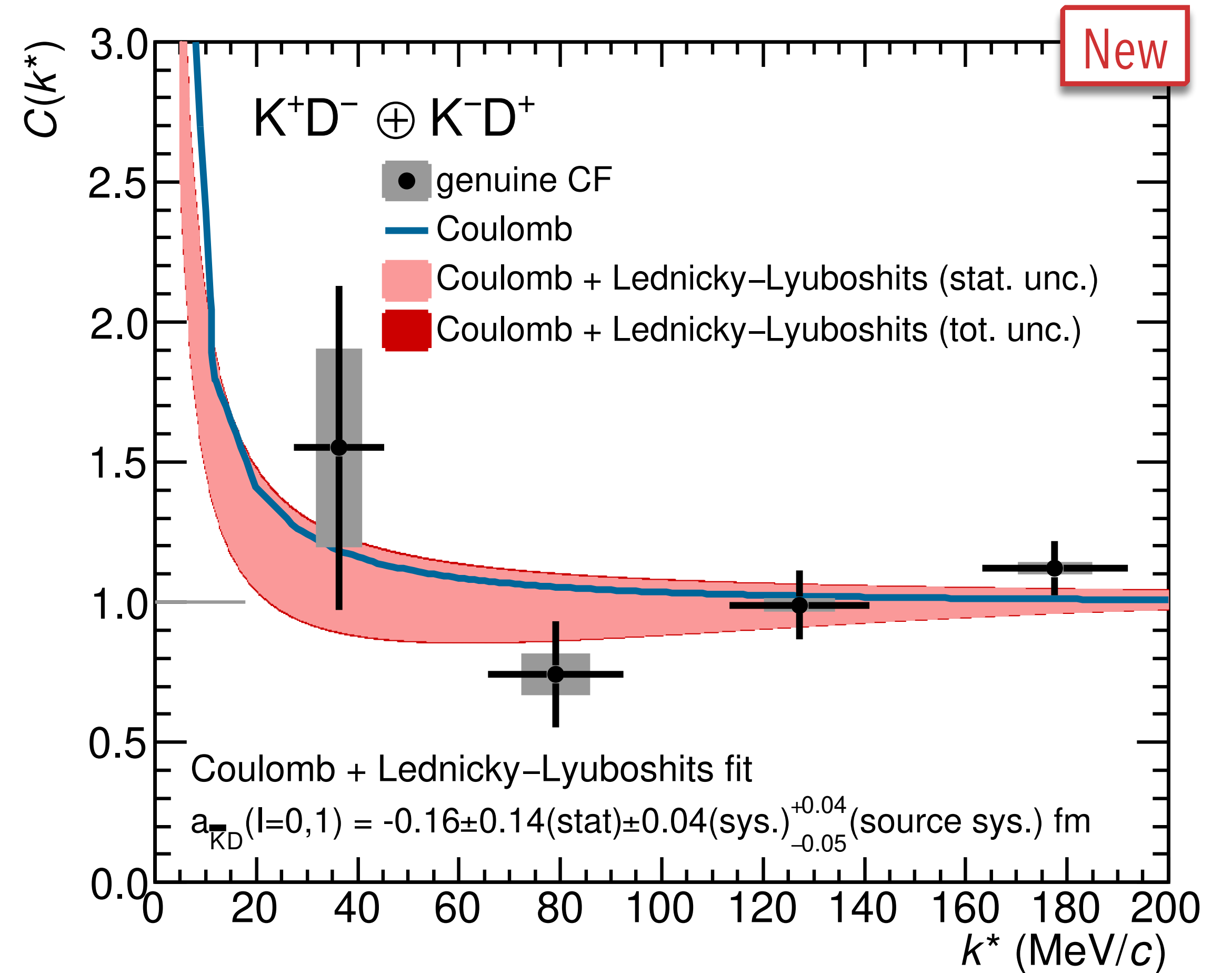
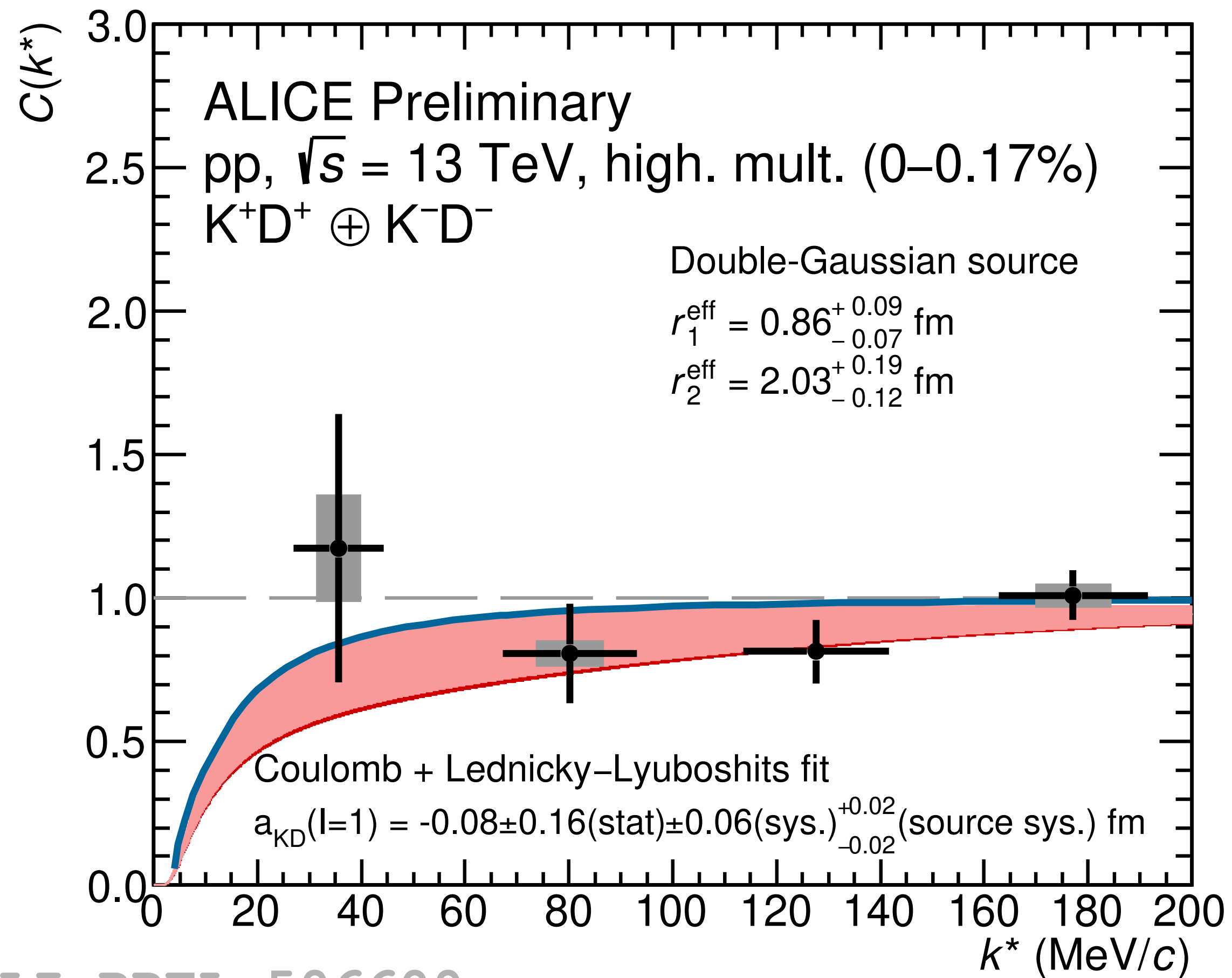


Bound-state pole formation corresponding to  $D_{s0}^*(2317)$

$$\begin{aligned}
 C'(k^*) = A_C(k^*) & \left\{ 2 \left[ \frac{1}{4} \left( \frac{|f_C(k^*)|}{r} \right)^2 \left[ 1 - \frac{d_0}{2\sqrt{\pi r}} + \frac{1}{2} (A_C(k^*) - 1)^2 (1 - e^{-4(rk^*)^2}) \right] + \right. \right. \\
 & + \mathcal{R}(f_C(k^*)) \frac{F_1(2k^*r)}{\sqrt{\pi r}} + \\
 & \left. \left. + \mathcal{I}(f_C(k^*)) \left[ \frac{F_2(2k^*r)}{2r} + (A_C(k^*) - 1)k^* \cos(rk^*) e^{-(rk^*)^2} \right] + 1 \right\}
 \end{aligned}$$

Where

$$f_C(k^*) = \left[ \frac{1}{a_0} + \frac{1}{2} d_0 k^{*2} - \frac{2}{a_C} h(k^* a_C) - ik^* A_C(k^*) \right]^{-1}$$



ALI-PREL-506600

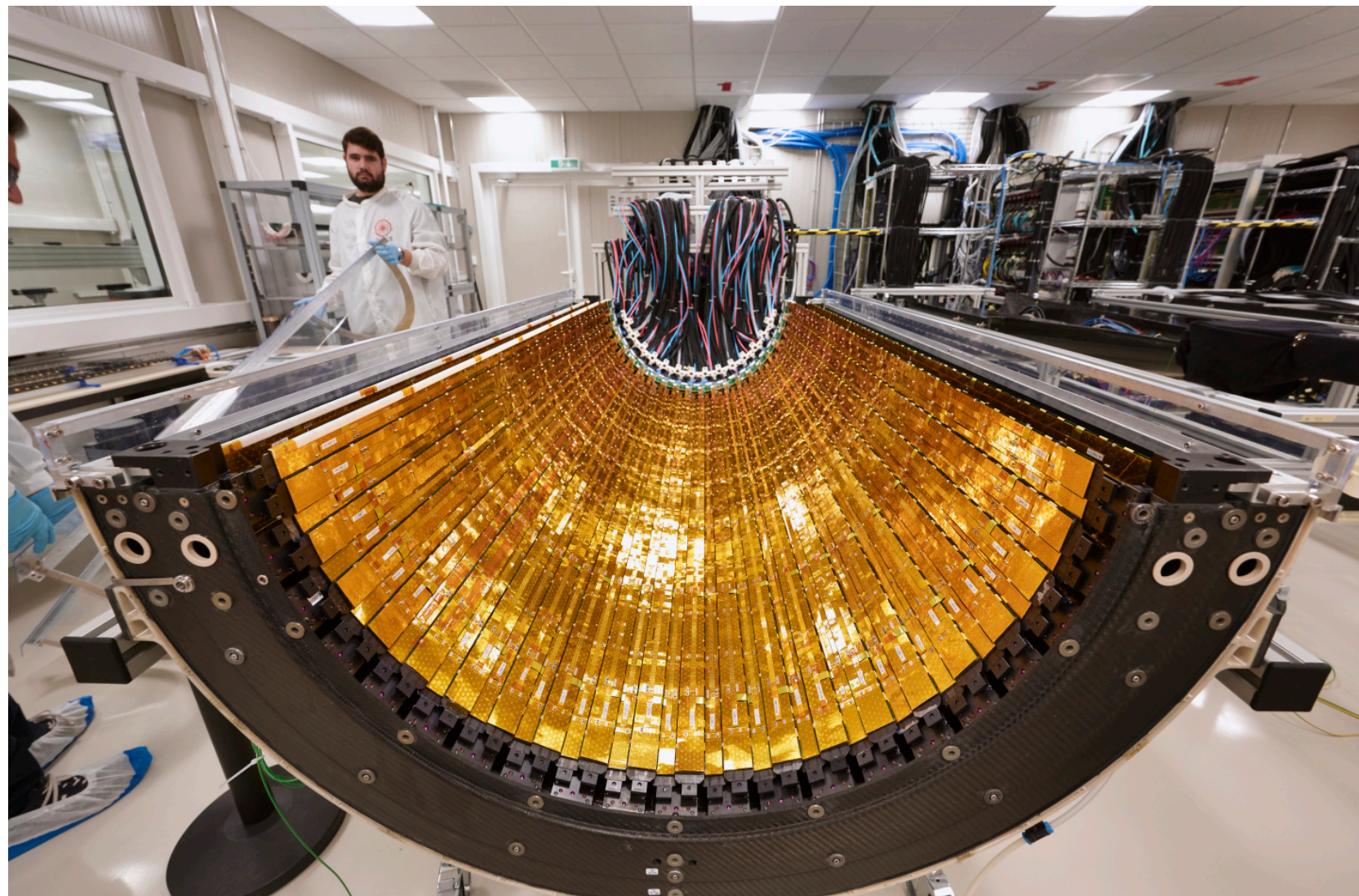
●  $K^+D^+$

●  $K^+D^-$

→  $l=1$  channel only

→  $l=0$  (50%)

→  $l=1$  (50%)



- Wide ALICE upgrade program for LHC Run 3 and 4 crucial for HF
  - increase collected Pb-Pb luminosity by more than one order of magnitude
  - new silicon Inner Tracking System

← Run 3: ITS2  TDR: CERN-LHCC-2013-024

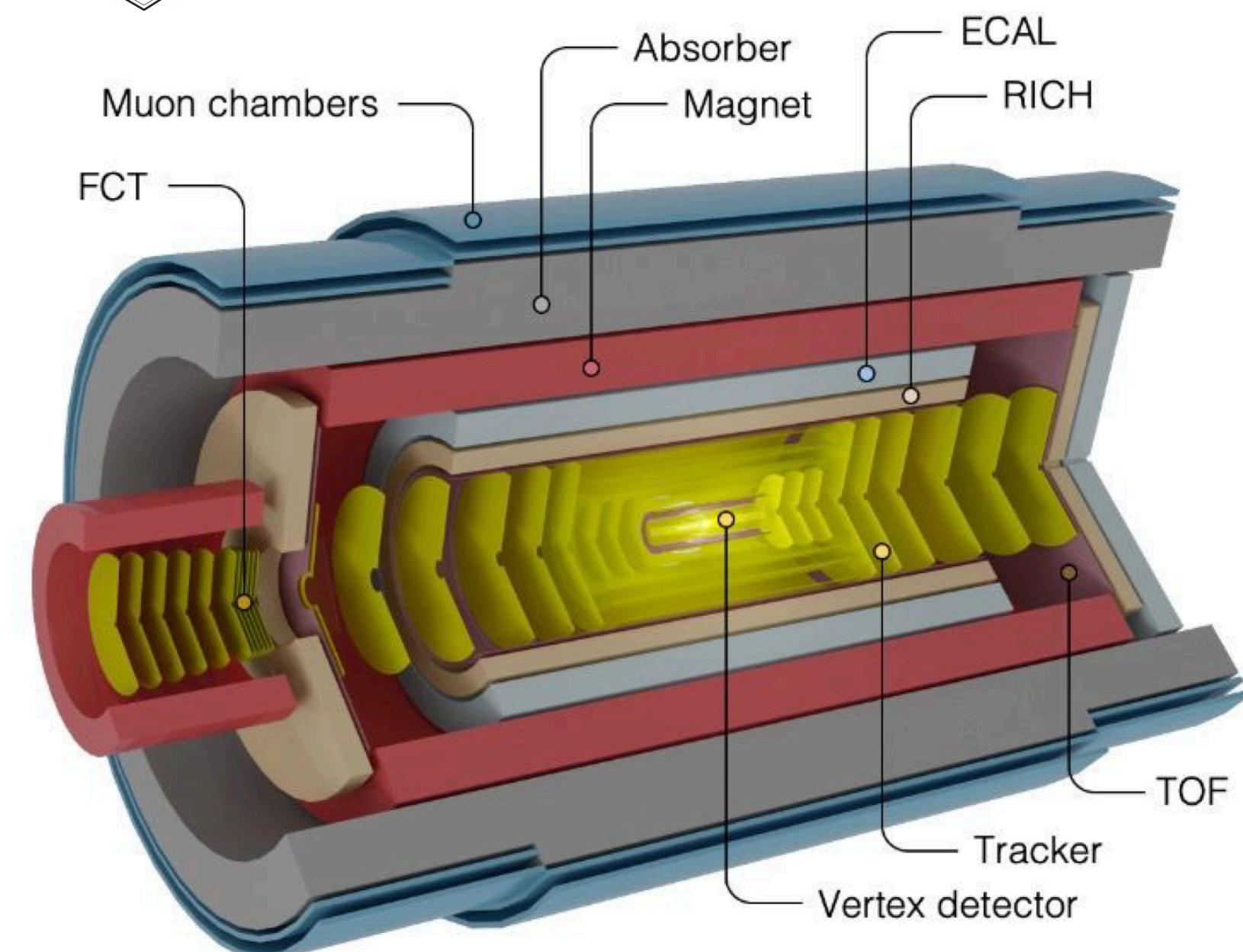
Run 4: ITS3  LOI: CERN-LHCC-2019-018 →

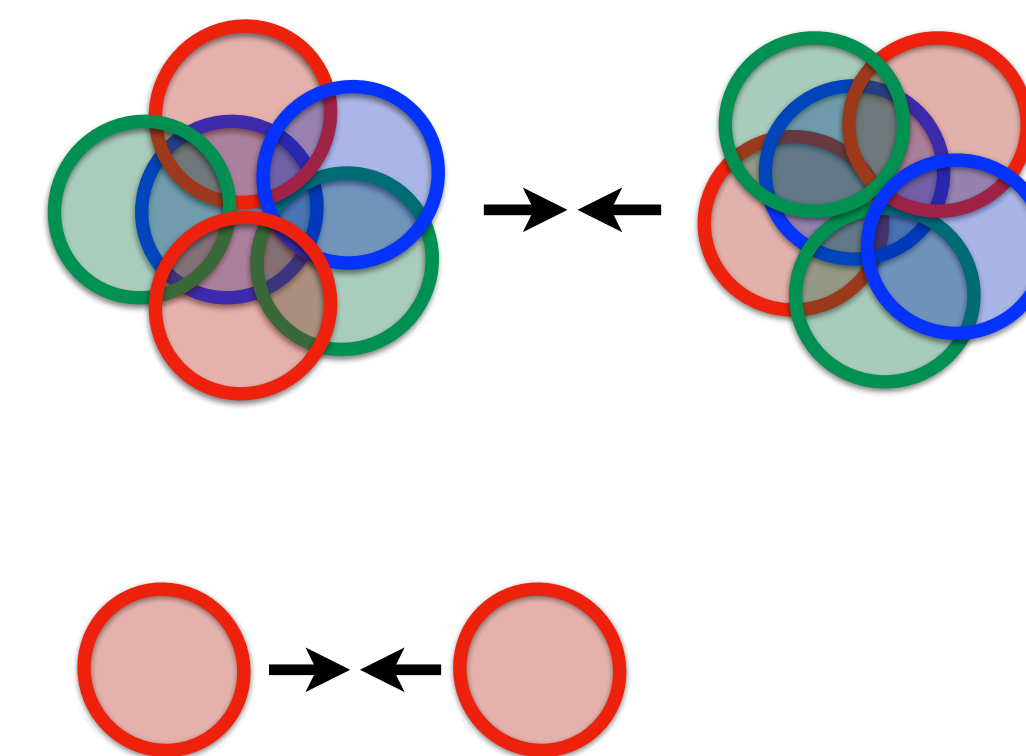
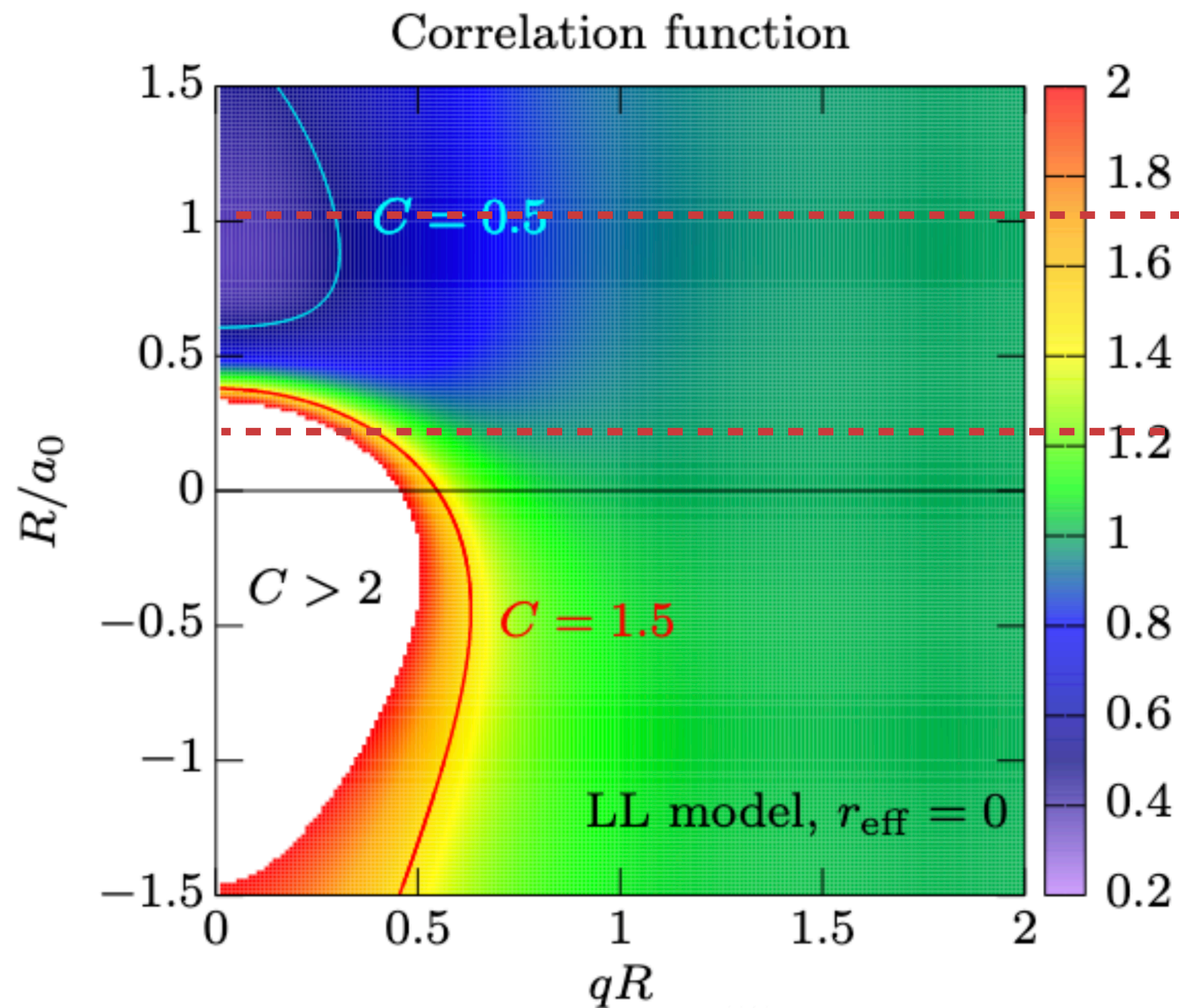
→ Run 5: all silicon ultra-light detector



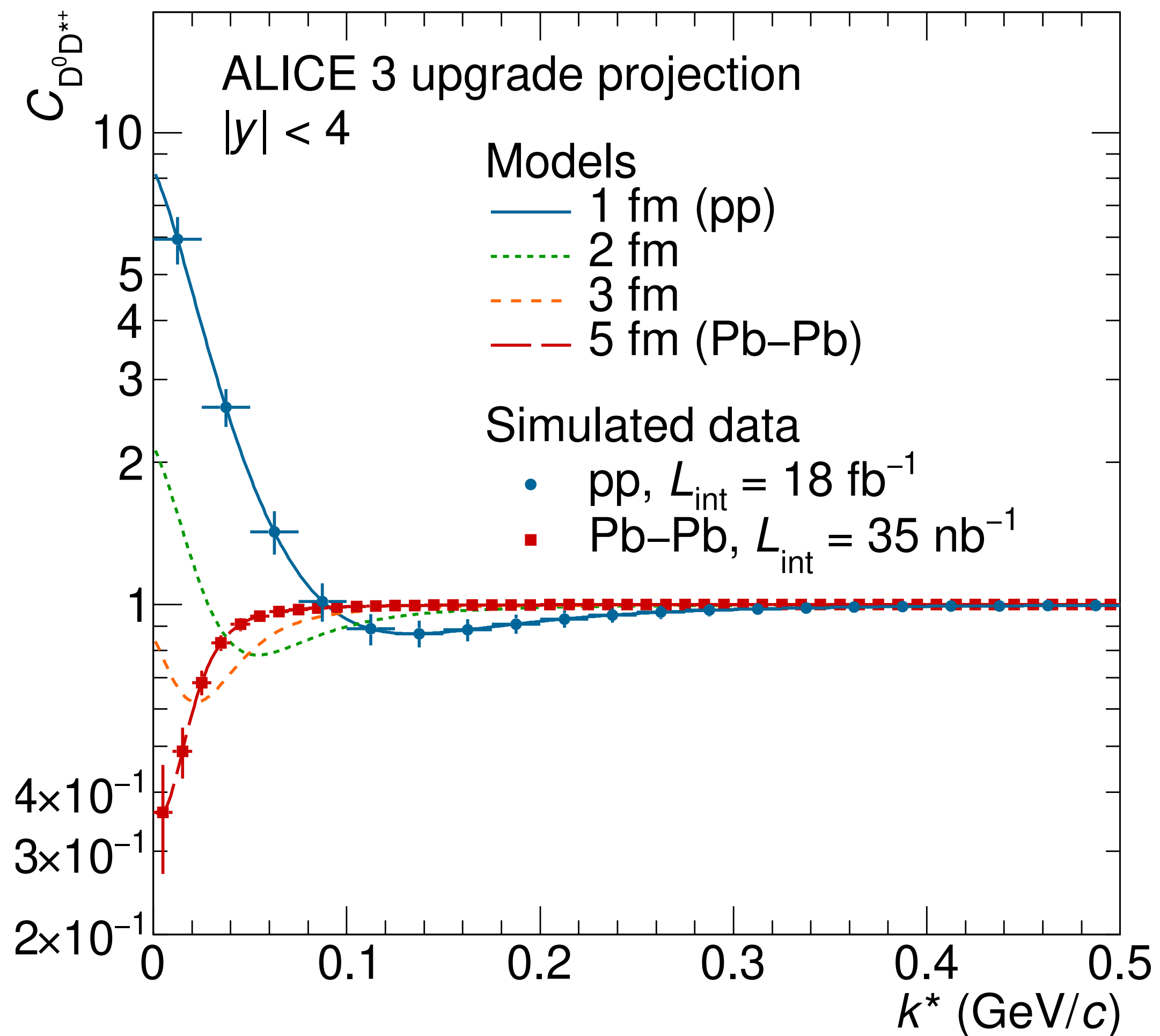
→  LOI: CERN-LHCC-2022-009

- New / more precise HF measurements down to low  $p_T$ 
  - Precise measurements of charm mesons and baryons
  - Access to measurements of beauty-strange meson and beauty-baryon production and azimuthal anisotropy
  - Run 5: multi-charm baryon production in heavy-ion collisions and charm-charm momentum correlations

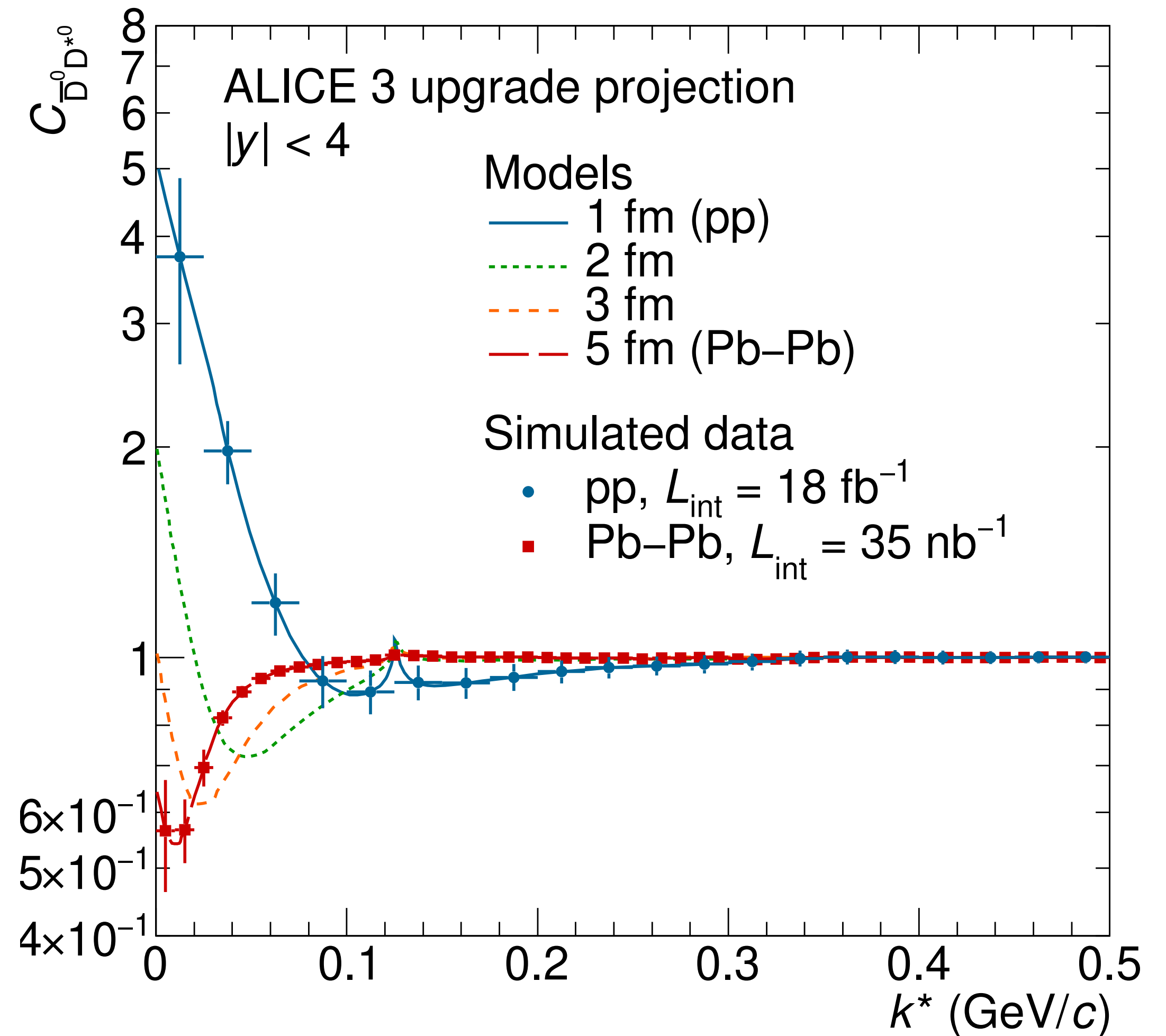




- Interplay between system size and scattering length
- ➔ size-dependent modification of the correlation function in presence of a bound state



ALI-SIMUL-502575



ALI-SIMUL-502579

- Expected precision enough to observe different behaviour of the CF in pp and Pb-Pb collisions in case  $T_{\text{cc}^+}$  and  $X(3872)$  are bound states of D mesons